Lecture 10: Coordinate and Momentum Representations

- We will start by considering the quantum description of the motion of a particle in one dimension.
- In classical mechanics, the state of the particle is given by its position and momentum coordinates, *x* and *p*.
- In quantum mechanics, we will consider position and momentum as observables and therefore represent them by Hermitian operators, *X* and *P*, respectively.
- Based on experimental evidence, we have deduced that:

$$[X,P] = i\hbar I$$

Incompatible Observables

 If two operators do not commute, then an eigenstate of one cannot be an eigenstate of the other.
 Proof:

-Strategy: assume opposite and show contradiction - Assume 10, by such that $A|a,b\rangle = a|a,b\rangle$ $B|a,b\rangle = b|a,b\rangle$ - Operate on labs with M:= [A, B] M(a,b) = AB(a,b) - BA(c,b)= bAla, b> - a RIG, b> $(ba-ab)(q,b) \implies M(a,b) = 0$ - either det [M] = 0 or [a,b) = 0 -if det [M] = o then (a, b) does not exist -Then A and B are compatible _iF [A,13]=0, then M=0, Ila, by do exist -Then A and B are incompatible

Coordinate Eigenstates

- Clearly *X* and *P* are *incompatible*, thus a particle cannot simultaneously have a well-defined position *and* momentum
- Since X is a Hermitian operator, it follows that its eigenstates form a complete set of unit vectors:



as cn = cn 14> => 4(x) = < x 14>

Momentum Representation

• The same logic applies also to momentum eigenstates:

$$P|p\rangle = p|p\rangle \quad \forall p \in \mathbb{R}$$
$$$$\int dp|p\rangle
$$\wedge (c_p) :=$$$$$$

 In order to move fluidly between coordinate and momentum basis, we need to know the transformation coefficients (x|p) and (p|x):

$$Proof: \Lambda(p) = \langle p|\Lambda \rangle \qquad \Lambda'(\gamma) = \langle \chi|\Lambda \rangle \\ = \int d\chi \langle p|\chi \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle p|\Lambda \rangle \\ = \int d\chi \langle p|\chi \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle p|\Lambda \rangle \\ = \int d\chi \langle p|\chi \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle P|\Lambda \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle P|\Lambda \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle P|\Lambda \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle P|\Lambda \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle P|\Lambda \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\Lambda \rangle \qquad = \int dp \langle \chi|p \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\Lambda \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\Lambda \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|p \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \\ = \int d\mu \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle \langle \chi|\rho \rangle$$

Deriving $\langle x|p \rangle$

- This is the direct derivation
 - Most textbooks use a round-about approach to avoid mathematical subtleties
 - We will just tackle them head on

Derivation of $\langle x|P|x' \rangle$

 $\langle x \rangle P | x' \rangle = i \frac{i \pi}{(x - x')}$

$$\langle x | P | x' \rangle = i\hbar \frac{\delta(x - x')}{(x - x')}$$

- At first glance this looks like a monstrosity, it is zero for *x*≠*x*', but at *x* = *x*', it is infinity divided by zero
 - By treating the delta-distribution correctly, we will see that we can easily understand the meaning of this result

Distribution Theory

- Q: is $\delta(x)$ a function?
 - A: no, technically it is a 'distribution'
 - A 'function' is a mapping from one space onto another:

$$y = f(x)$$

- A 'distribution' is more general than a function
- A distribution is defined only under integration

$$y = \int dx F(x) f(x)$$

- Here, *F*(*x*) is the distribution and *f*(*x*) is an ordinary function
- For example, the delta-distribution is defined by:

$$\int dx \,\delta(x - x_0) f(x) = f(x_0)$$

- A distribution can also be defined as the limit of a sequence of functions. All properties of the distribution are by definition, the limiting properties of the sequence
 - Example:

$$\delta(x) = \lim_{\sigma \to 0} \frac{1}{\sqrt{\pi\sigma^2}} e^{-x^2/\sigma^2}$$

$\langle x|P|x' \rangle$ is a Distribution

Our previous result can be written as:
 -(x-x')²/c⁻²

- This is well-behaved for any finite σ
- This proves that it has a clear meaning.
- Insert result into original equation:

$$(F_{0r} < x | p >)$$

$$\int dx' \text{ it } \frac{\delta(x - x')}{(x - x')} \langle x' | p > = p \langle x | p \rangle$$

• Expand $\langle x_{\perp}|p \rangle$ around $x_{\perp}=x$:

 $\langle x' | p \rangle = \langle x | p \rangle + (x' - x) \frac{d}{dx} \langle x | p \rangle + (x - x)^{2} \frac{d^{2}}{dx^{2}} \langle x | p \rangle + \dots$

Continued

$$\int dx' \frac{S(x-x')}{(x-x')} \left[(x|p) + (x'-x) \frac{d}{dx} (x|p) + \frac{(x'-x)^2 d^2}{2} \frac{d^2}{dx^2} (x|p) + \dots \right]$$

$$= -\frac{c}{\pi} p < x |p)$$

$$(x|p) \int dx' \frac{S(x-x')}{(x-x')} - \frac{d}{dx} (x|p) \int dx' \delta(x-x')$$

$$+ \frac{d^2}{dx^2} (x|p) \int dx' (x-x') \frac{S(x-x')}{x} + \dots = -\frac{c}{\pi} p < x |p)$$

$$\int dx' \frac{S(x-x')}{(x-x')} = \lim_{n \to \infty} \frac{1}{\sqrt{n}} \int dx' \frac{e}{(x-x')}$$

$$= \lim_{n \to \infty} \frac{1}{\sqrt{n}} \int dx' \frac{e}{(x-x')}$$

$$= \lim_{n \to \infty} \frac{1}{\sqrt{n}} \int dx' \frac{e}{\sqrt{n}} \int dx' \frac{e}{\sqrt{n}}$$

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