

Lecture 16:
Probability Current II and 1D
Scattering

Phy851 Fall 2009

Continuity Equation

$$j(x - \varepsilon) - j(x + \varepsilon) = \frac{dP(x, t)}{dt}$$

$$j(x - \varepsilon) - j(x + \varepsilon) = \frac{d\rho(x, t)2\varepsilon}{dt}$$

$$\frac{j(x - \varepsilon) - j(x + \varepsilon)}{2\varepsilon} = \frac{d\rho(x, t)}{dt}$$

$$-\frac{d}{dx} j(x, t) = \frac{d}{dt} \rho(x, t)$$

- This is the standard continuity equation, valid for any kind of fluid

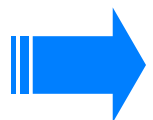
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- For energy eigenstates (stationary states), we need:

$$\frac{d}{dt} \rho(x, t) = 0 \Rightarrow \rho(x, t) = \rho(x, 0)$$

$$\frac{d}{dt} j(x, t) = 0 \Rightarrow j(x, t) = j(x, 0)$$

- This gives: $\frac{d}{dx} j(x, t) = 0 \Rightarrow j(x, t) = j_0$

- Must have **spatially uniform** current in steady state (of course j can be zero)



Derivation of the probability current:

Starting points:

- probability density: $\rho(x,t) = \psi^*(x,t) \psi(x,t)$
 $\rho = \psi^* \psi$

- continuity equation: $j' = -\dot{\rho}$ ← given ρ defines j

$$\begin{aligned} \dot{j}' &= -\dot{\rho} \\ &= -\dot{\psi}^* \psi - \psi^* \dot{\psi} \end{aligned}$$

$$\dot{\psi} = \frac{i\hbar}{2m} \psi'' - \frac{i}{\hbar} V \psi \xrightarrow{\text{c.c.}} \dot{\psi}^* = -\frac{i\hbar}{2m} \psi^{*''} + \frac{i}{\hbar} V \psi^*$$

therefore: $j' = \frac{i\hbar}{2m} \psi^{*''} \psi - \cancel{\frac{i}{\hbar} V \psi^* \psi} - \frac{i\hbar}{2m} \psi^* \psi'' + \cancel{\frac{i}{\hbar} V \psi^* \psi}$

$$j' = \frac{i\hbar}{2m} (\psi^{*''} \psi - \psi^* \psi'')$$

Trick: $\frac{d}{dx} (\psi^{*'} \psi - \psi^* \psi') = \psi^{*''} \psi + \cancel{\psi^{*'} \psi'} - \cancel{\psi^* \psi''} - \psi^* \psi''$
 $= \psi^{*''} \psi - \psi^* \psi''$

$$\frac{d}{dx} j(x,t) = \frac{i\hbar}{2m} \frac{d}{dx} (\psi^{*'} \psi - \psi^* \psi')$$

up to const.

$$j(x,t) = -\frac{i\hbar}{2m} \left(\psi^*(x,t) \frac{d}{dx} \psi(x,t) - \psi(x,t) \frac{d}{dx} \psi^*(x,t) \right)$$

Current of a plane wave

- For a plane wave we have:

$$\psi(x) = ae^{ikx}$$

- The corresponding probability current is:

$$j = -i \frac{\hbar}{2m} (\psi^* \psi' - \psi'^* \psi)$$

$$= -i \frac{\hbar |a|^2}{2m} (ik - (-ik)) = |a|^2 \frac{\hbar k}{m}$$

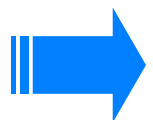
density

velocity

- So for a plane wave, we find:

$$\vec{j}(x, t) = \rho_0 \vec{v}_0$$

This result is fairly intuitive



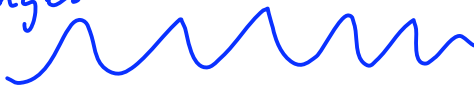
Quantum Interference terms

- Consider a superposition of plane waves:

$$\psi(x) = a_1 e^{ik_1 x} + a_2 e^{ik_2 x}$$

- The probability density is:

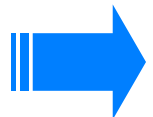
$$\rho(x) = |a_1|^2 + |a_2|^2 + \underbrace{\left(a_1^* a_2 e^{i(k_2 - k_1)x} + c.c. \right)}_{\text{Interference Term}}$$

'Fringes' 

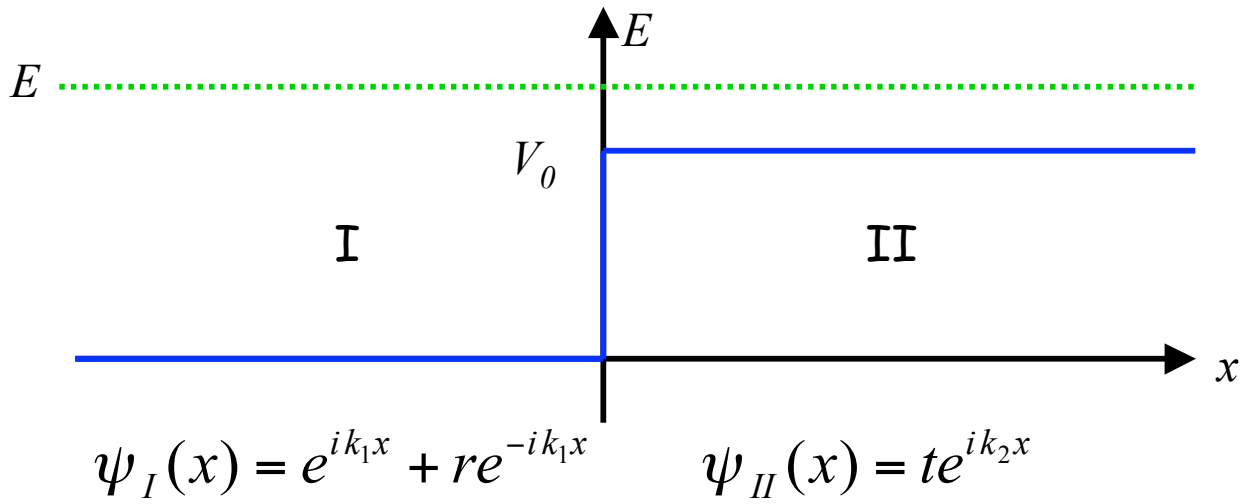
- The probability current density is:

$$j(x) = |a_1|^2 \frac{\hbar k_1}{m} + |a_2|^2 \frac{\hbar k_2}{m} + \underbrace{\left(a_1^* a_2 e^{i(k_2 - k_1)x} + c.c. \right)}_{\text{Interference Term}} \frac{\hbar(k_1 + k_2)}{2m}$$

- Note that the interference term in $j(x,t)$ vanishes for $k_2 = -k_1$
 - **This is always the case for Energy Eigenstates** → Currents are then purely additive
 - There is still interference in the *probability density* due to the presence of left and right currents, just not in the *probability current*.



Return to the Step Potential



$$r = \frac{k_1 - k_2}{k_1 + k_2}$$

$$t = \frac{2k_1}{k_1 + k_2}$$

- The probability current density is:

$$j_I(x) = \frac{\hbar k_1}{m} - |r|^2 \frac{\hbar k_1}{m} = (1 - |r|^2) \frac{\hbar k_1}{m}$$

$$j_{II}(x) = |t|^2 \frac{\hbar k_2}{m}$$

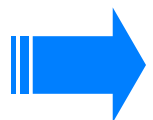
- Spatially uniform current requires: $j_I(x) = j_{II}(x)$
 - So the probability conservation law is:

$$(1 - |r|^2) \frac{\hbar k_1}{m} = |t|^2 \frac{\hbar k_2}{m}$$

Rearrange terms to
get more intuitive
result

$$|r|^2 + |t|^2 \frac{k_2}{k_1} = 1$$

$$j_{out} = j_{in}$$



Continued

$$|r|^2 + |t|^2 \frac{k_2}{k_1} = 1$$

$$t = \frac{2k_1}{k_1 + k_2}$$

$$r = \frac{k_1 - k_2}{k_1 + k_2}$$

- Transmission and reflection probabilities are derived from conservation law: $j_{in} = j_{out}$

$$T := \left| \frac{j_{out}(x > 0)}{j_{in}} \right|$$

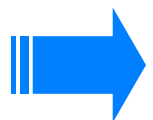
$$R := \left| \frac{j_{out}(x < 0)}{j_{in}} \right|$$

Any constant we might have put in front of the incident wave would cancel out here

- For step potential, this gives:

$$T = \frac{|t|^2 k_2}{k_1} = \frac{4k_1 k_2}{|k_1 + k_2|^2} \quad R = \frac{|r|^2 k_1}{k_1} = |r|^2$$

$$R + T = \frac{(k_1 - k_2)^2 + 4k_1 k_2}{(k_1 + k_2)^2} = \frac{k_1^2 + 2k_1 k_2 + k_2^2}{(k_1 + k_2)^2} = 1$$

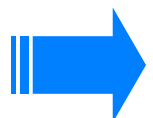


Probability current: Summary/conclusions

- The proper way to compute probability in scattering is via probability current.
- The probability for a particle to scatter into a certain channel is the ratio of the outgoing current in that channel to the total incoming current.
- In 1D scattering at fixed energy, we can treat the left-traveling and right-traveling components of the current as independent
 - because there are no interference terms in the current density for $+k$ and $-k$ currents.
 - Allows us to group components into 'incoming' and 'outgoing' currents
- For a plane-wave, the current is the amplitude squared times the velocity.

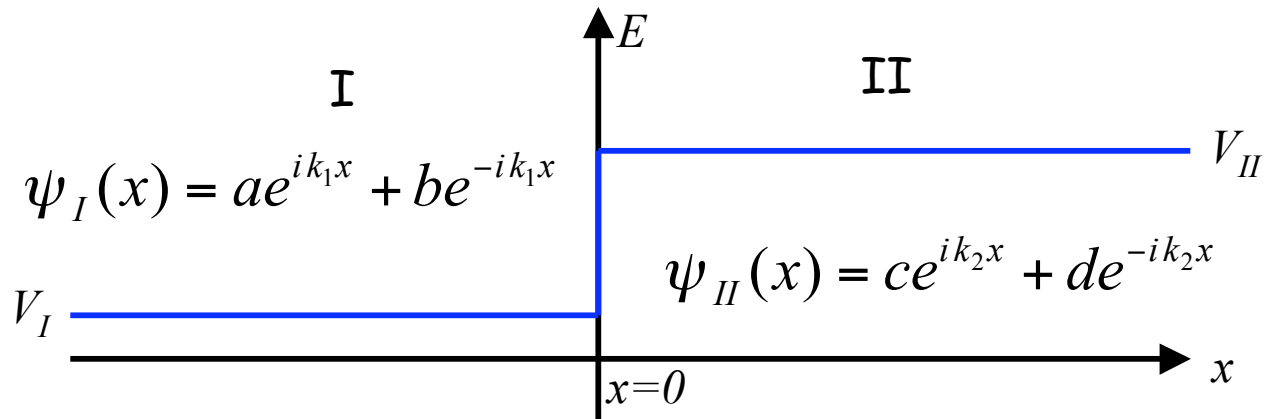
Important Shortcut:

- If you are asked to compute R and T , for a 1d scattering problem, you can:
 - Compute $R=|r|^2$
 - Then use $T=1-R$ (conservation law)
 - i.e. you don't need to compute t or worry about current unless specifically asked to do so.
 - **Wrong:** $T=|t|^2$ then $R = 1-T$ would not work



Scattering From a Potential Step revisited

- Lets solve the step potential again, but with the possibility for left *and* right travelling incoming waves:



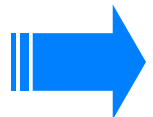
$$k_1 = \left(U(E - V_I) + iU(V_I - E) \right) \frac{\sqrt{2m|E - V_I|}}{\hbar}$$

$$k_2 = \left(U(E - V_{II}) + iU(V_{II} - E) \right) \frac{\sqrt{2m|E - V_{II}|}}{\hbar}$$

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- Boundary condition equations:

$$\psi_I(0) = \psi_{II}(0) \Rightarrow a + b = c + d$$

$$\psi'_I(0) = \psi'_{II}(0) \Rightarrow k_1(a - b) = k_2(c - d)$$



Scattering Matrix

- The scattering matrix gives the outgoing amplitudes in terms of the incoming amplitudes:

Definition of scattering matrix, S :
$$\begin{pmatrix} b \\ c \end{pmatrix} = S \begin{pmatrix} a \\ d \end{pmatrix}$$

$$a + b = c + d$$

$$b - c = -a + d$$

$$k_1(a - b) = k_2(c - d) \quad \longrightarrow \quad k_1b + k_2c = k_1a + k_2d$$

Boundary
condition
equations

Move outgoing amplitudes to
l.h.s. and incoming to r.h.s.

$$\begin{pmatrix} 1 & -1 \\ k_1 & k_2 \end{pmatrix} \begin{pmatrix} b \\ c \end{pmatrix} = \begin{pmatrix} -1 & 1 \\ k_1 & k_2 \end{pmatrix} \begin{pmatrix} a \\ d \end{pmatrix}$$

b.c. eqs in
matrix
form

Solve for the
outgoing
amplitudes

$$\begin{pmatrix} b \\ c \end{pmatrix} = \begin{pmatrix} 1 & -1 \\ k_1 & k_2 \end{pmatrix}^{-1} \begin{pmatrix} -1 & 1 \\ k_1 & k_2 \end{pmatrix} \begin{pmatrix} a \\ d \end{pmatrix}$$

Calculate the
inverse

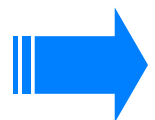
$$\begin{pmatrix} b \\ c \end{pmatrix} = \frac{1}{k_1 + k_2} \begin{pmatrix} k_2 & 1 \\ -k_1 & 1 \end{pmatrix} \begin{pmatrix} -1 & 1 \\ k_1 & k_2 \end{pmatrix} \begin{pmatrix} a \\ d \end{pmatrix}$$

Multiply
matrices to
get:

$$\begin{pmatrix} b \\ c \end{pmatrix} = \frac{1}{k_1 + k_2} \begin{pmatrix} k_1 - k_2 & 2k_2 \\ 2k_1 & k_2 - k_1 \end{pmatrix} \begin{pmatrix} a \\ d \end{pmatrix}$$

Result:

$$S = \frac{1}{k_1 + k_2} \begin{pmatrix} k_1 - k_2 & 2k_2 \\ 2k_1 & k_2 - k_1 \end{pmatrix}$$



Example: wave incident from left

- Consider case $a=1$, $b=r$, $c=t$, and $d=0$:

$$\begin{pmatrix} b \\ c \end{pmatrix} = S \begin{pmatrix} a \\ d \end{pmatrix}$$

$$S = \frac{1}{k_1 + k_2} \begin{pmatrix} k_1 - k_2 & 2k_2 \\ 2k_1 & k_2 - k_1 \end{pmatrix}$$

$$\begin{pmatrix} r \\ t \end{pmatrix} = \frac{1}{k_1 + k_2} \begin{pmatrix} k_1 - k_2 & 2k_2 \\ 2k_1 & k_2 - k_1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

- We find: $r = \frac{k_1 - k_2}{k_1 + k_2}$ $t = \frac{2k_1}{k_1 + k_2}$

- For left and right incoming waves:

$$\begin{pmatrix} r \\ t \end{pmatrix} = \frac{1}{k_1 + k_2} \begin{pmatrix} k_1 - k_2 & 2k_2 \\ 2k_1 & k_2 - k_1 \end{pmatrix} \begin{pmatrix} c_L \\ c_R \end{pmatrix}$$

$$r = \frac{(k_1 - k_2)c_L + 2k_2c_R}{k_1 + k_2} \quad t = \frac{2k_1c_L + (k_2 - k_1)c_R}{k_1 + k_2}$$

$$j_{in} = |c_L|^2 k_1 + |c_R|^2 k_2$$

$$R = \frac{|r|^2 k_1}{j_{in}} = \frac{|k_1 - k_2|^2 |c_L|^2 + 4|k_2|^2 |c_R|^2 + 4 \operatorname{Re} \{ k_2^* (k_1 - k_2) c_R^* c_L \}}{|k_1 + k_2|^2 (|c_L|^2 k_1 + |c_R|^2 k_2)}$$

$$T = 1 - R$$

