### Lecture 25: Introduction to the Quantum Theory of Angular Momentum

Phy851 Fall 2009



# <u>Goals</u>

- 1. Understand how to use different coordinate systems in QM (cartesian, spherical,...)
- 2. Derive the quantum mechanical properties of Angular Momentum
  - Use an algebraic approach similar to what we did for the Harmonic Oscillator
- 3. Use the resulting theory to treat spherically symmetric problems in three dimensions
  - Calculating the Hydrogen atom energy levels will be our target goal







## Motion In 3Dimensions

- For a particle moving in three dimensions, there is a *distinct quantum state for every point* in space.
  - Thus each position state is now labeled by a vector  $\vec{D}$

$$\begin{array}{c} X \to R \\ \left| x \right\rangle \to \left| \vec{r} \right\rangle \end{array}$$

- Vector operators are really three operators

$$\vec{R} = X \vec{e}_x + Y \vec{e}_y + Z \vec{e}_z$$

 $\frac{\text{Scalar Operators}}{X,Y,Z} \qquad \frac{\text{Ordinary Vectors}}{\vec{e}_x,\vec{e}_y,\vec{e}_z}$ 

- Coordinate system not unique:
  - For example, we could use spherical coordinates

$$\vec{R} = R \, \vec{e}_r(\Theta, \Phi)$$

$$\frac{\text{Scalar Operators}}{R,\Theta,\Phi} \qquad \frac{\text{Vector Operator}}{\vec{e}_r(\Theta,\Phi)}$$

 $\vec{e}_r(\Theta, \Phi) = \sin \Theta \cos \Phi \vec{e}_x + \sin \Theta \sin \Phi \vec{e}_y + \cos \Theta \vec{e}_z$ 

#### You can never go wrong with Cartesian Coordinates

- In all other coordinate systems, the unit vectors are also operators
  - so must be treated carefully
- Eigenstates:

$$\vec{R} \big| \vec{r} \big\rangle = \vec{r} \big| \vec{r} \big\rangle$$

- For each point in space there is a position eigenstate
- How we want to label these points is up to us:

$$\begin{aligned} X \left| \vec{r} \right\rangle &= \vec{e}_x \cdot \vec{R} \left| \vec{r} \right\rangle = \vec{e}_x \cdot \vec{r} \left| \vec{r} \right\rangle = x \left| \vec{r} \right\rangle \\ Y \left| \vec{r} \right\rangle &= y \left| \vec{r} \right\rangle \qquad Z \left| \vec{r} \right\rangle = z \left| \vec{r} \right\rangle \end{aligned}$$

 $\begin{aligned} \left| \vec{r} \right\rangle &= \left| x, y, z \right\rangle \\ \left| \vec{r} \right\rangle &= \left| r, \theta, \phi \right\rangle \end{aligned} \qquad \begin{array}{c} \text{Different ways} \\ \text{to refer to} \\ \text{the same state} \end{aligned}$ 



### **Orthogonality**

• In three dimensions, the orthogonality condition becomes:

 $\langle \vec{r} | \vec{r}' \rangle = \delta^3 (\vec{r} - \vec{r}')$ 

- Cartesian coordinates:

$$\langle xyz | x'y'z' \rangle = \delta(x - x')\delta(y - y')\delta(z - z')$$

- Spherical coordinates:

$$\langle r\theta\phi|r'\theta'\phi'\rangle = \delta(r-r')\frac{\delta(\theta-\theta')}{r\sin\theta}\frac{\delta(\phi-\phi')}{r}$$

- Mixing coordinates:

 $\langle xyz | r\theta\phi \rangle = \delta(x - r\sin\theta\cos\phi)\delta(y - r\sin\theta\sin\phi)\delta(z - r\cos\theta)$ 



## Wavefunctions

Wavefunctions are defined in the usual way as:

$$\psi(r) = \langle r | \psi \rangle$$
$$\int d^3 r |\psi(\vec{r})|^2 = 1$$

$$\psi(x, y, z) = \langle xyz | \psi \rangle$$
$$\int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dz |\psi(x, y, z)|^{2} = 1$$

$$\psi(r,\theta,\phi) = \left\langle r\theta\phi \left|\psi\right.\right\rangle$$
$$\int_{0}^{\infty} dr \int_{0}^{\pi} r\sin\theta d\theta \int_{0}^{2\pi} rd\phi \left|\psi(r,\theta,\phi)\right|^{2} = 1$$



#### <u>Momentum</u>

• The three-dimensional momentum vector operator is:

$$\vec{P} = \vec{e}_x P_x + \vec{e}_y P_y + \vec{e}_z P_z$$

• The three-dimensional Hamiltonian is

$$H = \frac{\vec{P} \cdot \vec{P}}{2m} + V(\vec{R})$$

• In Cartesian components, this becomes:

$$H = \frac{P_x^2}{2m} + \frac{P_y^2}{2m} + \frac{P_z^2}{2m} + V(\vec{R})$$

$$\left\langle xyz \left| H \right| \psi \right\rangle = \left[ -\frac{\hbar^2}{2m} \left( \frac{d^2}{dx^2} + \frac{d^2}{dy^2} + \frac{d^2}{dz^2} \right) + V(x, y, z) \right] \left\langle xyz \left| \psi \right\rangle$$

or

 $\left\langle \vec{r} \left| H \right| \psi \right\rangle = \left[ -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \left\langle \vec{r} \left| \psi \right\rangle$ 



#### Angular Momentum

- We can decompose the momentum operator onto spherical components as:  $\vec{P} = \vec{e}_r(\Theta, \Phi)P_r + \vec{e}_\theta(\Theta, \Phi)P_\theta + \vec{e}_\phi(\Theta, \Phi)P_\phi$  $P_r = \vec{e}_r(\Theta, \Phi) \cdot \vec{P}$  $P_\theta = \vec{e}_\theta(\Theta, \Phi) \cdot \vec{P}$  $P_\theta = \vec{e}_\theta(\Theta, \Phi) \cdot \vec{P}$
- The unit-vector operators are:  $\vec{e}_r(\Theta, \Phi) = \sin \Theta \cos \Phi \vec{e}_x + \sin \Theta \sin \Phi \vec{e}_y + \cos \Theta \vec{e}_z$   $\vec{e}_{\theta}(\Theta, \Phi) = -\cos \Theta \cos \Phi \vec{e}_x - \cos \Theta \sin \Phi \vec{e}_y + \sin \Theta \vec{e}_z$   $\vec{e}_{\phi}(\Theta, \Phi) = -\sin \Phi \vec{e}_x + \cos \Phi \vec{e}_y$ 
  - From the above decomposition, we are led to define the Angular momentum as:

$$\vec{L} = \vec{e}_{\theta} (\Theta, \Phi) P_{\theta} + \vec{e}_{\phi} (\Theta, \Phi) P_{\phi}$$

 This is not the usual definition (L=RxP) but is equivalent.