Chapter 13: Relativistic Quantum Mechanics

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The problem

Suppose that an electron of momentum k coming from the left strikes a one-dimensional potential barrier

$$e\Phi(x) = V, x > 0$$
$$e\Phi(x) = 0, x < 0.$$

Calculate the transmission and reflection coefficients for the cases where E < V < 2m and V > 2m and interpret the results.



First, starting with the Dirac Equation:

$$(\vec{\alpha} \cdot \vec{p} + \beta m + e\Phi)\Psi(\vec{r}, t) = 0$$

We are looking at an electron, so we'll start by assuming that we are using the positive energy, right helicity solution for an electron.



This general solution looks like:

$$\Psi_R^+ = e^{i(kx - Et)} \begin{pmatrix} 1\\ 0\\ k\\ \overline{m+E}\\ 0 \end{pmatrix}$$

How to define helicity:

Right-handed has spin projection parallel to momentum.

Left-handed has spin projection antiparallel to momentum.



We next need to write the solutions of the Dirac equation in both regions. Writing the incident, reflected, and transmitted wavefunctions separately.

$$\begin{split} \Psi_{i}(x) &= e^{ikx} \begin{pmatrix} 1\\0\\k\\\overline{E+m}\\0 \end{pmatrix} \\ \Psi_{r}(x) &= be^{-ikx} \begin{pmatrix} 1\\0\\-k\\\overline{E+m}\\0 \end{pmatrix} + b'e^{-ikx} \begin{pmatrix} 0\\1\\0\\k\\\overline{E+m}\\0 \end{pmatrix} \\ \Psi_{t}(x) &= de^{iqx} \begin{pmatrix} 1\\0\\q\\\overline{E-V+m}\\0 \end{pmatrix} + d'e^{-iqx} \begin{pmatrix} 0\\1\\0\\-k\\\overline{E-V+m}\\0 \end{pmatrix} \end{split}$$



Where $k = \sqrt{E^2 - m^2}$

and

$$q = \sqrt{(E - V)^2 - m^2}$$

Note that b' = d' = 0 because we are assuming that the barrier isn't causing a spin-flip.

After removing states (due to no spin-flip at the barrier).

Putting these together gives the following

$$\begin{split} \Psi_{I}(x) &= \Psi_{i} + \Psi_{r} = e^{ikx} \begin{pmatrix} 1\\0\\k\\\overline{E+m}\\0 \end{pmatrix} + be^{-ikx} \begin{pmatrix} 1\\0\\-k\\\overline{E+m}\\0 \end{pmatrix} \\ \Psi_{II}(x) &= de^{iqx} \begin{pmatrix} 1\\0\\\frac{q}{E-V+m}\\0 \end{pmatrix} \end{split}$$



Now, at the boundary, $\Psi_I(0) = \Psi_2(0)$

We match the respective spinor indices with each other. This will give two equations with 2 unknowns,

$$(1-b)\frac{k}{E+m} = d \frac{q}{E-V+m}$$

Now, divide the equations and solving for b and d (after some algebra/Mathematica simplification, whichever is preferred)

$$b = \frac{1-\xi}{1+\xi}$$
 and $d = \frac{2}{1+\xi}$ where $\xi = \frac{q}{k} \frac{(E+m)}{E-V+m}$

We want the transmission and reflection coefficients, which are defined in terms of currents as

$$R = -\frac{j_r}{j_i}$$
 and $T = \frac{j_t}{j_i}$ Using the definition of current from Sakurai to be

In 4-vector notation:

For the 3-vector component

 $j^{\mu} = \overline{\Psi} \gamma^{\mu} \Psi$ where $\overline{\Psi} = \Psi^{\dagger} \beta$

$$\vec{J} = \Psi^{\dagger} \vec{\alpha} \Psi$$

where
$$\vec{\alpha} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix}$$

We are working in 1-D, so we are going to use α_z , or

$$\alpha = \begin{pmatrix} 0 & \sigma_z \\ \sigma_z & 0 \end{pmatrix}$$

Solving for the incident, reflected, and transmitted current densities:

$$j_{i} = \Psi_{i}^{\dagger} \alpha \Psi_{i} = e^{-ikx} \begin{pmatrix} 1 \\ 0 \\ \frac{k}{E+m} \end{pmatrix}^{T} \begin{pmatrix} 0 & \sigma_{z} \\ \sigma_{z} & 0 \end{pmatrix} e^{ikx} \begin{pmatrix} 1 \\ 0 \\ \frac{k}{E+m} \end{pmatrix} = 2\frac{k}{E+m}$$
$$j_{r} = \Psi_{r}^{\dagger} \alpha \Psi_{r} = be^{ikx} \begin{pmatrix} 1 \\ 0 \\ \frac{-k}{E+m} \end{pmatrix}^{T} \begin{pmatrix} 0 & \sigma_{z} \\ \sigma_{z} & 0 \end{pmatrix} be^{-ikx} \begin{pmatrix} 1 \\ 0 \\ \frac{-k}{E+m} \end{pmatrix} = -2b^{2}\frac{k}{E+m}$$
$$j_{t} = \Psi_{t}^{\dagger} \alpha \Psi_{t} = de^{-iqx} \begin{pmatrix} 1 \\ 0 \\ \frac{q}{E-V+m} \end{pmatrix}^{T} \begin{pmatrix} 0 & \sigma_{z} \\ \sigma_{z} & 0 \end{pmatrix} de^{iqx} \begin{pmatrix} 1 \\ 0 \\ \frac{q}{E-V+m} \end{pmatrix} = 2d^{2}\frac{q}{E-V+m}$$

Solving for the incident, reflected, and transmitted current densities:

$$j_{i} = 2\frac{k}{E+m} \qquad j_{r} = -2b^{2}\frac{k}{E+m} \qquad j_{t} = 2d^{2}\frac{q}{E-V+m}$$
$$R = -\frac{j_{r}}{j_{i}} = |b|^{2} = \left|\frac{1-\xi}{1+\xi}\right|^{2} \qquad T = \frac{j_{t}}{j_{i}} = |d|^{2}\frac{q}{k}\frac{E+m}{E-V+m}$$

Now, the different cases:

1) V<2m
2) V>2m
$$\xi = \frac{q}{k} \frac{(E+m)}{E-V+m}$$

For a sanity check, can also show that at the boundary

$$j_i + j_r = j_t$$

so the current is conserved.



Now, the different cases:

1) For V<2m, everything is normal. $\xi > 0$, and like normal R+T = 1, and R,T will be between 0 and 1

- 2) For V>2m, weird stuff happens. $\xi < 0$, meaning that b > 1 which implies that R>1 and T<0. But how can this happen?
 - \rightarrow This is the Klein Paradox. Electron-holes are produced at the boundary

2) For V>2m, weird stuff happens. ξ < 0, meaning that b > 1 which implies that R>1 and T<0. But how can this happen?
→This is the Klein Paradox. Electron-holes are produced at the boundary



 $T = \frac{j_t}{j_i} = |d|^2 \frac{q}{k} \frac{E+m}{E-V+m}$