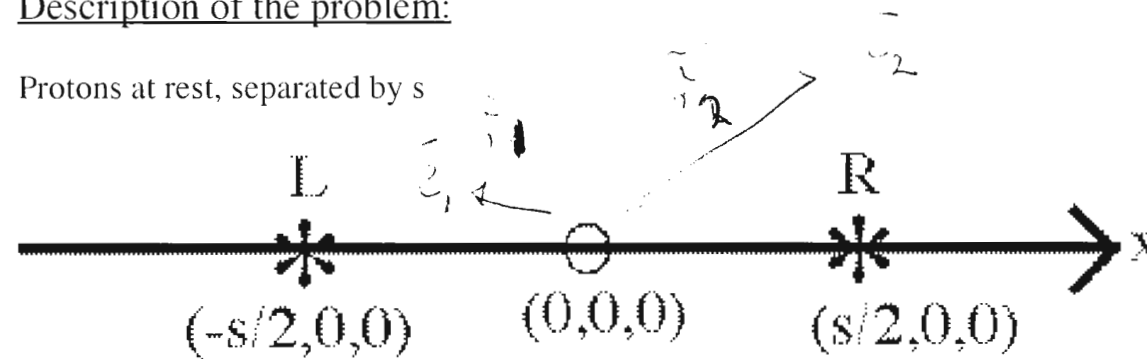


# HYDROGEN MOLECULE - VARIATIONAL CALCULATION

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## Description of the problem:

Protons at rest, separated by  $s$



\* represents a proton and o is the origin. Note that the protons are on the x-axis.  
L&R represent protons; 1 and 2 stand for electrons.

Initial conditions

- Electrons in spin antisymmetric state
- Spatial part of the wave function is symmetric
- Without loss of generality, one can take the spatial part of the wave function to be positive everywhere

Recall that

$$H\Psi(\vec{r}_1, \vec{r}_2, s) = E_o(s)\Psi(\vec{r}_1, \vec{r}_2, s) \quad (1)$$

The Hamiltonian for the system is

$$H = -\frac{\nabla_1^2}{2m} - \frac{\nabla_2^2}{2m} - \frac{e^2}{r_{1L}} - \frac{e^2}{r_{1R}} - \frac{e^2}{r_{2L}} - \frac{e^2}{r_{2R}} + \frac{e^2}{r_{12}} \quad (2)$$

where

$$\bullet \vec{r}_{1L} = \left| \vec{r}_1 + \frac{s}{2} \hat{x} \right| \quad (3a)$$

$$\bullet \vec{r}_{2L} = \left| \vec{r}_2 + \frac{s}{2} \hat{x} \right| \quad (3b)$$

$$\bullet \vec{r}_{1R} = \left| \vec{r}_1 - \frac{s}{2} \hat{x} \right| \quad (3c)$$

- $\bar{r}_{2R} = \left| \bar{r}_2 - \frac{s}{2} \hat{x} \right|$  (3d)

- $\bar{r}_{12} = |\bar{r}_1 - \bar{r}_2|$  (3e)

Multiply eq. (1) by  $\Psi(\bar{r}_1, \bar{r}_2, s)$  and integrate over  $\bar{r}_1$  and  $\bar{r}_2$

$$\int d\bar{r}_1 \int d\bar{r}_2 \Psi(\bar{r}_1, \bar{r}_2, s) H \Psi(\bar{r}_1, \bar{r}_2, s) = E_o(s) \int d\bar{r}_1 \int d\bar{r}_2 \Psi^2(\bar{r}_1, \bar{r}_2, s)$$

or

$$E_o(s) = \frac{\int d\bar{r}_1 \int d\bar{r}_2 \Psi(\bar{r}_1, \bar{r}_2, s) H \Psi(\bar{r}_1, \bar{r}_2, s)}{\int d\bar{r}_1 \int d\bar{r}_2 \Psi^2(\bar{r}_1, \bar{r}_2, s)} = \frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

Now let

$$\omega(\bar{r}_1, \bar{r}_2, s) = \frac{\Psi^2(\bar{r}_1, \bar{r}_2, s)}{\int d\bar{r}_1 \int d\bar{r}_2 \Psi^2(\bar{r}_1, \bar{r}_2, s)} = \text{probability density function} \quad (4)$$

and

$$\epsilon(\bar{r}_1, \bar{r}_2, s) = \frac{H \Psi(\bar{r}_1, \bar{r}_2, s)}{\Psi(\bar{r}_1, \bar{r}_2, s)} \quad (5)$$

giving the following

$$E_o(s) = \int d\bar{r}_1 \int d\bar{r}_2 \omega(\bar{r}_1, \bar{r}_2, s) \epsilon(\bar{r}_1, \bar{r}_2, s) = \int \omega \epsilon \quad (6)$$

### Choice of the wave function:

Suppose the electron - electron interaction is zero. In that case, let us find out what the wave function will be. To the non-interacting wave function, we will add terms that bring in the effects of electron-electron interaction ( $e^- - e^-$ ) corrections.

### Non-interacting simulation:

For the non-interacting system, let

$$H_0 \tilde{\Psi}_0(\vec{r}_1, \vec{r}_2, s) = \tilde{E}_0(s) \tilde{\Psi}_0(\vec{r}_1, \vec{r}_2, s).$$

where the Hamiltonian is

$$H_0 = -\frac{\nabla_1^2}{2m} - \frac{\nabla_2^2}{2m} - \frac{e^2}{r_{1L}} - \frac{e^2}{r_{1R}} - \frac{e^2}{r_{2L}} - \frac{e^2}{r_{2R}} = H_{01} + H_{02}. \quad (7)$$

In equation (7)

$$H_{01} = -\frac{\nabla_1^2}{2m} - \frac{e^2}{r_{1L}} - \frac{e^2}{r_{1R}} \quad (8)$$

and

$$H_{02} = -\frac{\nabla_2^2}{2m} - \frac{e^2}{r_{2L}} - \frac{e^2}{r_{2R}}. \quad (9)$$

Substitute (7) in (1)

$$(H_{01} + H_{02}) \tilde{\Psi}_0(\vec{r}_1, \vec{r}_2) = \tilde{E}_0 \tilde{\Psi}_0(\vec{r}_1, \vec{r}_2) \quad (10)$$

Let

$$\tilde{\Psi}_0(\vec{r}_1, \vec{r}_2, s) = \phi_1(\vec{r}_1) \phi_2(\vec{r}_2). \quad (11)$$

Plugging (11) into (10), one finds

$$\tilde{E}_0(s) = \tilde{E}_1 + \tilde{E}_2 \quad (12)$$

where

$$H_{01} \phi(\vec{r}_1) = \tilde{E}_1 \phi(\vec{r}_1) \quad (13)$$

and

$$H_{02} \phi(\vec{r}_2) = \tilde{E}_2 \phi(\vec{r}_2) \quad (14)$$

$\phi(\vec{r}_i, s)$  is the independent particle wave function shared equally between the 2 protons,

$$\Psi(r_1, s) = e^{-r_{1L}/a} + e^{-r_{1R}/a} \quad (15)$$

$$\Psi(r_2, s) = e^{-r_{2L}/a} + e^{-r_{2R}/a} \quad (16)$$

In the presence of electron-electron interaction, let

$$\Psi(r_1, r_2, s) = \varphi(r_1, s) \cdot \varphi(r_2, s) \cdot f(r_{12}) \quad (17)$$

which is the correlated wave function. In the above equation,  $f(r_{12})$  is

$$f(r_{12}) = \exp\left(\frac{r_{12}}{\alpha(1 + \beta r_{12})}\right) \quad (18)$$

For large distances  $r_{12}$ ,  $\frac{e^2}{r_{12}} \rightarrow 0$

$$f(r_{12}) \rightarrow \text{constant}$$

and

$$\Psi(\vec{r}_1, \vec{r}_2, s) \rightarrow \varphi_1(\vec{r}_1, s) \varphi_2(\vec{r}_2, s) \quad (19)$$

### Coulomb-Cusp condition:

The Schrodinger equation for 1 hydrogen atom is

$$H\varphi = \left( -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{r} \right) \varphi = E\varphi \quad (20)$$

Let us examine the kinetic energy term as electron 1 approaches proton L. Then

$$\frac{1}{\varphi} \nabla_i^2 \varphi \rightarrow -\frac{2}{ar_{iL}} \left[ \frac{1}{1 + e^{-s/a}} \right] \quad (21)$$

Multiplying the above equation by  $-\frac{\hbar^2}{2m}$ , one obtains

$$-\frac{\hbar^2}{2m} \frac{1}{\phi} \nabla^2 \phi \rightarrow \frac{\hbar^2}{m} \frac{1}{ar_{1L}} \frac{1}{1+e^{-s/a}} \quad (22)$$

The singularity as  $r_{1L} \rightarrow 0$  is offset by the singularity from the coulomb attraction. This yields the condition

$$\frac{\hbar^2}{ma} \frac{1}{1+e^{-s/a}} = e^2 \quad (23)$$

Recall that  $\frac{\hbar^2}{me^2} = a_0$ , Bohr Radius. Plugging  $a_0$  into the above equation one finds

$$a = \frac{a_0}{1+e^{-s/a}} \quad (24)$$

Similarly, one can show that  $\alpha = 2a_0$ .

Start with

$$f(r_{12}) = \exp\left(\frac{r_{12}}{\alpha(1+\beta r_{12})}\right) \quad (25)$$

$$\nabla_1^2 f(r_{12}) = f \frac{1}{\alpha^2(1+\beta r_{12})^4} + \frac{3f}{\alpha r_{12}(1+\beta r_{12})^2} + \frac{f}{\alpha} \left[ -\frac{1}{r_{12}} \frac{1}{(1+\beta r_{12})^2} - \frac{2\beta}{(1+\beta r_{12})^3} \right] \quad (26)$$

For  $r_{12} \rightarrow 0$

$$\nabla_1^2 f(r_{12}) \rightarrow \frac{2f}{\alpha r_{12}(1+\beta r_{12})^2} \quad (27)$$

Similarly, we can arrive at  $\nabla_2^2 f(r_{12}) \rightarrow \frac{2f}{\alpha r_{21}(1+\beta r_{12})^2}$  as  $r_{12} \rightarrow 0$ .

For small  $r_{12}$ , we arrive at the condition

$$-\frac{\hbar^2}{2m} \left[ \frac{2f}{\alpha} + \frac{2f}{\alpha} \right] + e^2 = 0 \quad (28)$$

$$\text{or } \alpha = \frac{2\hbar^2}{me^2} = 2a_0 \quad (29)$$

Outline of program:

The Monte Carlo simulation will be performed on equation (6) with  $\Psi(\vec{r}_1, \vec{r}_2, s)$  from equation (16). Fix  $s$  and  $\beta$  and vary the positions of the electrons one at a time until the minimum energy is found. (Note:  $s$  can be looked up.) Every time that you move an electron, if  $\omega_{\text{new}} > \omega_{\text{old}}$  accept the new position. If  $\omega_{\text{new}} < \omega_{\text{old}}$  then generate a random number,  $\gamma$ , between 0 and 1. Find  $\frac{\omega_{\text{new}}}{\omega_{\text{old}}}$  if this is greater than  $\gamma$  then accept the new position and calculate  $\epsilon$ , but if it is less than  $\gamma$  retain the old position. After obtaining a reliable estimate of energy  $E_0(s)$ , choose another  $\beta$  and find the minimum value of the energy.