

Investigation of jet quenching and elliptic flow within a pQCD-based partonic transport model

Oliver Fochler

Z. Xu C. Greiner

Institut für Theoretische Physik
Goethe-Universität Frankfurt

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Introduction

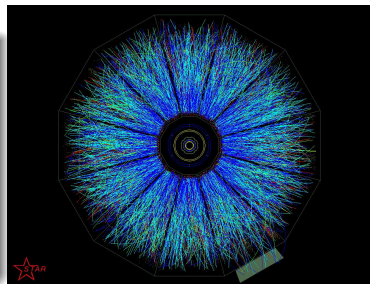
Heavy ion collisions are complicated!

Need models and theories for

- Initial state
- Medium evolution
- High- p_T physics
- Phase transition
- Freeze-out

Some tools:

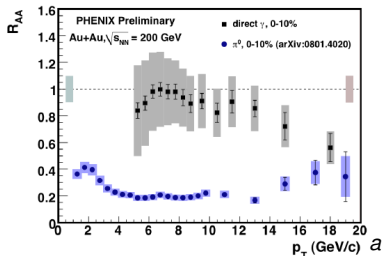
- Parameterizations (e.g. Bjorken)
- Hydrodynamics
- Transport models
- Lattice QCD
- pQCD (BDMPS, ASW, AMY, ...)



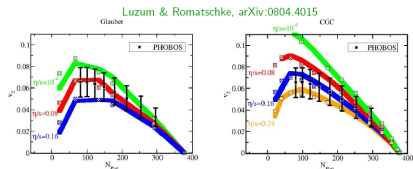
Problem

No model can describe all (most) aspects of the QGP evolution.

Introduction



^aHP 2008, C. Vale



Key observables

- Jet quenching, R_{AA} (high- p_t physics)
- Elliptic flow, v_2 (bulk physics)

- Difficult to describe within one model
- (Partonic) Transport with **only 2** \rightarrow **2** processes:
Need **unphysical cross sections** for v_2 . ($\Rightarrow R_{AA}$ too small)

BAMPS results

2 \leftrightarrow 3 processes + pQCD cross sections \Rightarrow enough elliptic flow

What about R_{AA} ?

- 1 The Model - BAMPS
- 2 Static Medium - Brick Scenario
- 3 Central and Non-Central Au+Au Collisions
- 4 Sensitivity on the LPM Cut-Off

BAMPS = Boltzmann Approach to Multiple Particle Scattering ¹

Microscopic transport simulations with full dynamics.

Attack various problems within *one* model.
(elliptic flow, R_{AA} , thermalization, ...)

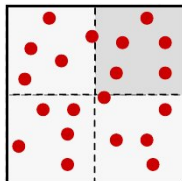
Solve Boltzmann equation for $2 \rightarrow 2$ and $2 \leftrightarrow 3$ processes based on LO pQCD matrix elements.

$$p^\mu \partial_\mu f(x, p) = C_{2 \rightarrow 2}(x, p) + C_{2 \leftrightarrow 3}(x, p)$$

¹Z. Xu, C. Greiner, Phys. Rev. C71 (2005)

Implemented processes

$$gg \rightarrow gg \text{ and } gg \leftrightarrow ggg$$



- Massless Boltzmann particles
- Sample transition **probabilities** for (test)particles in a given
 - Spatial cell, ΔV
 - Time step, Δt
- Use small angle cross section for $gg \rightarrow gg$ and **Gunion-Bertsch** matrix element for $gg \leftrightarrow ggg$

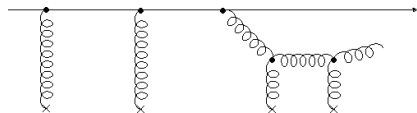
$$\frac{d\sigma_{gg \rightarrow gg}}{dq_{\perp}^2} \simeq \frac{9\pi\alpha_s^2}{(\mathbf{q}_{\perp}^2 + m_D^2)^2} \quad |\mathcal{M}_{gg \rightarrow ggg}|^2 = \frac{72\pi^2\alpha_s^2 s^2}{(\mathbf{q}_{\perp}^2 + m_D^2)^2} \frac{48\pi\alpha_s \mathbf{q}_{\perp}^2}{\mathbf{k}_{\perp}^2 [(\mathbf{k}_{\perp} - \mathbf{q}_{\perp})^2 + m_D^2]}$$

- Debye screening: $m_D^2 = d_G \pi \alpha_s \int \frac{d^3 p}{(2\pi)^3} \frac{1}{p} N_c f$

Landau-Pomeranchuk-Migdal (LPM) Effect

LPM effect

Multiple emission \Rightarrow Interference



- Not realizable in transport model with classical particles.
- Discard possible interference effects (Bethe-Heitler regime)

Parent must not scatter during formation time of emitted gluon.

$$|M_{gg \rightarrow ggg}|^2 \longrightarrow |M_{gg \rightarrow ggg}|^2 \Theta(\lambda - \tau)$$

Affects:

- Total cross section

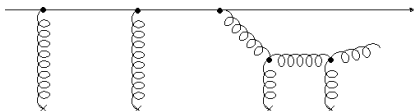
$$\sigma_{gg \rightarrow ggg} \sim \int dq_{\perp}^2 \int dk_{\perp}^2 \int dy \int d\phi \cdots |M_{gg \rightarrow ggg}|^2 \Theta(\lambda - \tau)$$

- Sampling of outgoing momenta

Landau-Pomeranchuk-Migdal (LPM) Effect

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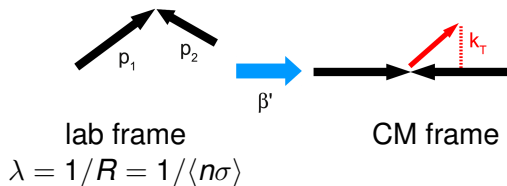
Reference Frames for the LPM Cut-Off



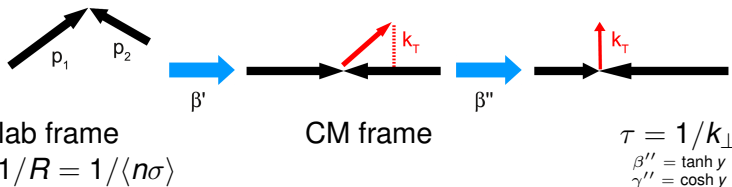
lab frame

$$\lambda = 1/R = 1/\langle n\sigma \rangle$$

Reference Frames for the LPM Cut-Off



Reference Frames for the LPM Cut-Off



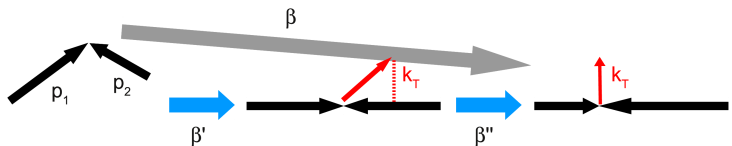
- Total boost: $\gamma = \gamma' \gamma'' (1 + \vec{\beta}' \vec{\beta}'') = \frac{\cosh y}{\sqrt{1 - \beta'^2}} (1 + \beta' \tanh y \cos \theta)$

$$\Theta(\lambda - \tau) \rightarrow \Theta\left(k_{\perp} - \frac{\gamma}{\lambda}\right) \rightarrow \Theta\left(k_{\perp} - \frac{1}{\lambda}\right) \Theta(B - (\cosh y + A \sinh y))$$

$$A = \beta' \cos \theta, B = k_{\perp} \Lambda_g \sqrt{1 - \beta'^2}$$

- Integral cuts: $\sigma_{gg \rightarrow ggg} \sim \int dq_{\perp}^2 \int_{1/\lambda^2}^{s/4} dk_{\perp}^2 \int_{y_{\min}}^{y_{\max}} dy \int d\phi$
- $\beta' \ll 1$ (thermal particles): $\Theta(\lambda - \tau) \simeq \Theta\left(k_{\perp} - \frac{\cosh y}{\lambda}\right)$

Reference Frames for the LPM Cut-Off



lab frame

$$\lambda = 1/R = 1/\langle n\sigma \rangle$$

CM frame

$$\tau = 1/k_{\perp}$$

$$\beta'' = \tanh y$$

$$\gamma'' = \cosh y$$

- Total boost: $\gamma = \gamma' \gamma'' (1 + \vec{\beta}' \vec{\beta}'') = \frac{\cosh y}{\sqrt{1 - \beta'^2}} (1 + \beta' \tanh y \cos \theta)$

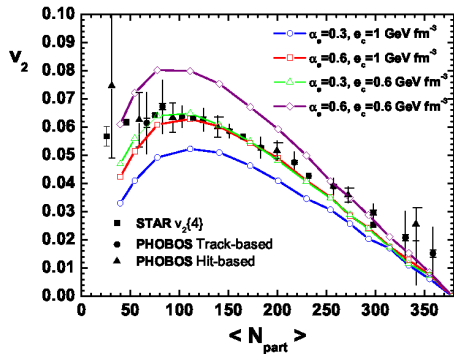
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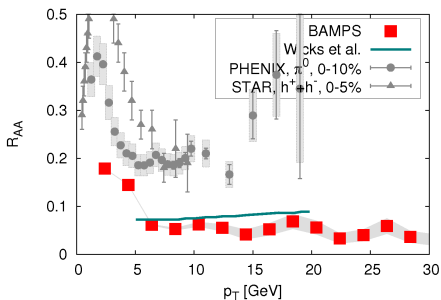
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v_2 and R_{AA} for Central Collisions

Integrated v_2 as function of N_{part}



Gluonic R_{AA} for $b = 0$ fm



OF, Z. Xu, C. Greiner, PRL 102 (2009)

We'll come back to this later..

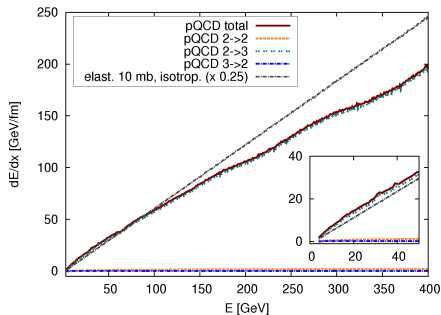
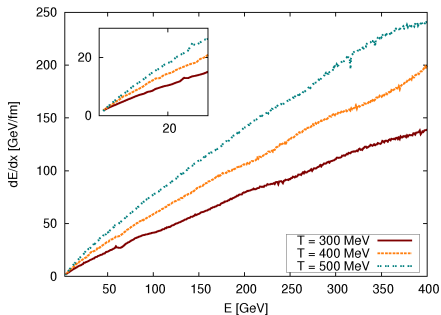
Energy Loss in a Static Medium

Brick

- Constant temperature T
- Only gluons, $N_f = 0$
- Static medium, no expansion

- $\alpha_s = 0.3$

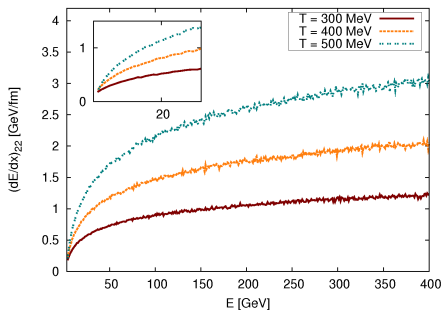
- $m_D^2 = \frac{8}{\pi} N_c \alpha_s T^2$



$T = 400$ MeV

Energy Loss in Binary Collisions

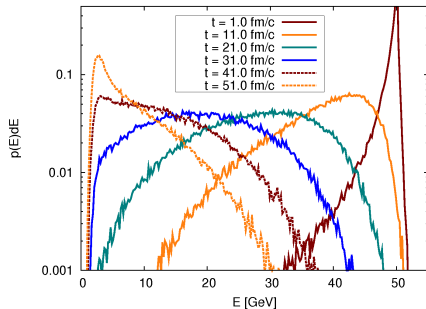
Elastic dE/dx for different T



- Elastic energy loss

$$\left. \frac{dE}{dx} \right|_{2 \rightarrow 2} \propto C_R \pi \alpha_s^2 T^2 \ln \left(\frac{4ET}{m_D^2} \right)$$

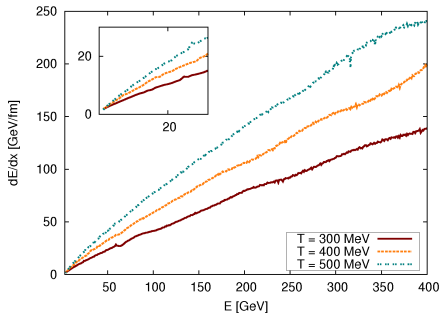
E-spectrum ($T = 400$ MeV, $E_0 = 50$ GeV)



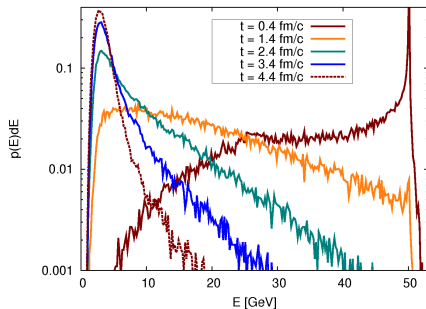
- Broad distribution
- Not: Mean energy loss + momentum diffusion (Gaussian)

Energy Loss in $gg \rightarrow ggg$ Processes

Total dE/dx for different T



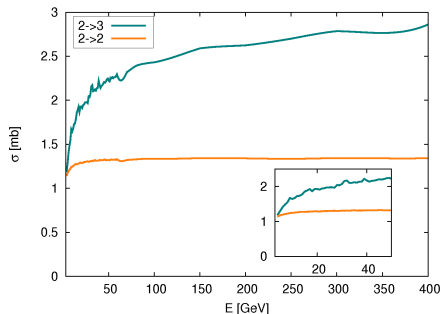
E-spectrum ($T = 400$ MeV, $E_0 = 50$ GeV)



- Large differential energy loss due to $gg \rightarrow ggg$
- Roughly $dE/dx \propto E$
- Rapid evolution of the energy spectrum

Energy Loss in $gg \rightarrow ggg$ Processes

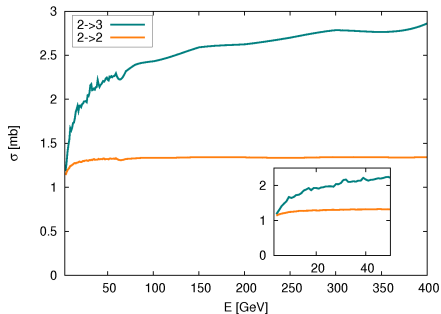
Cross sections ($T = 400$ MeV)



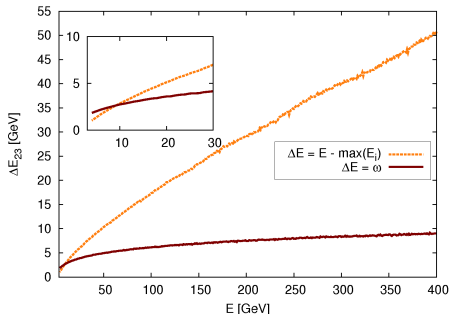
- Reasonable partonic cross sections over the whole energy range.
- Definition of the energy loss ΔE matters:
 - $\Delta E = E_{in} - \max(E_{out}^i)$
 - $\Delta E = \omega$

Energy Loss in $gg \rightarrow ggg$ Processes

Cross sections (T = 400 MeV)



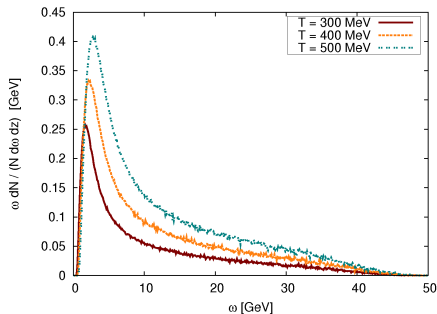
Energy loss in single $gg \rightarrow ggg$



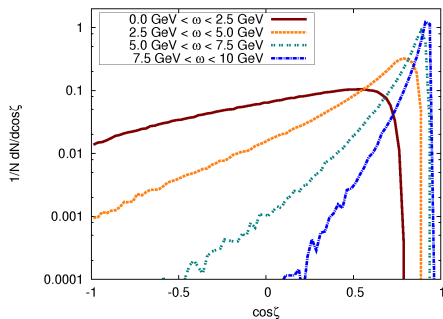
- Reasonable partonic cross sections over the whole energy range.
- Definition of the energy loss ΔE matters:
 - $\Delta E = E_{in} - \max(E_{out}^i)$
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Gluon Radiation from the Gunion-Bertsch Matrix Element

Energy spectrum

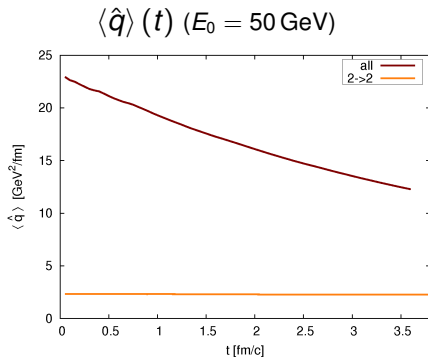
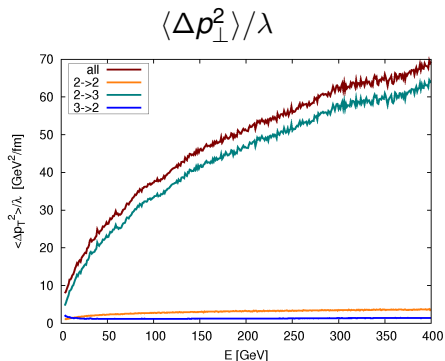


Angular distribution



- Emission at small angles not allowed due to LPM cut-off $\Theta(k_{\perp} - \frac{\gamma}{\lambda})$
- Effect more pronounced for small gluon energies.

Accumulated Transverse Momenta



- \hat{q} : Transverse momentum accumulated over path length L .

$$\hat{q}(L) = \frac{1}{L} \sum_i (\Delta p_{\perp}^2)_i \qquad \langle \hat{q} \rangle (t) = \frac{1}{t} \int_0^t \left. \frac{\langle \Delta p_{\perp}^2 \rangle}{\lambda} \right|_{E(\tilde{t})} d\tilde{t}$$

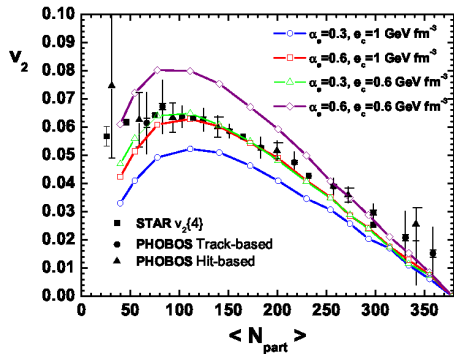
- We extend investigation to $2 \rightarrow 3$ processes.

Setup for Au+Au Collisions in BAMPS

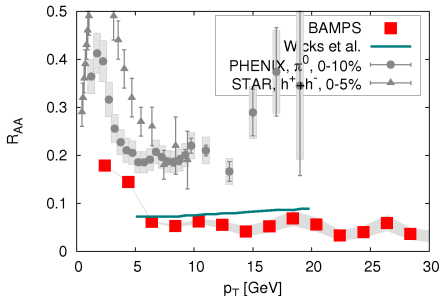
- Au+Au at RHIC energy 200 AGeV
- Currently limited to gluons
- $\alpha_s = 0.3$
- Mini-Jet initial conditions (lower cut-off $p_0 = 1.4$ GeV)
 - Glück-Reya-Vogt PDFs
 - Wood-Saxon density profile + Glauber model
- Formation time for each parton $\Delta t_f = \cosh(y)/p_T$
- Free streaming applied to cells where $\varepsilon < \varepsilon_c$
(choices: $\varepsilon_c = 0.6$ GeV/fm³ or $\varepsilon_c = 1$ GeV/fm³)

v_2 and R_{AA} for Central Collisions

Integrated v_2 as function of N_{part}



Gluonic R_{AA} for $b = 0 \text{ fm}$

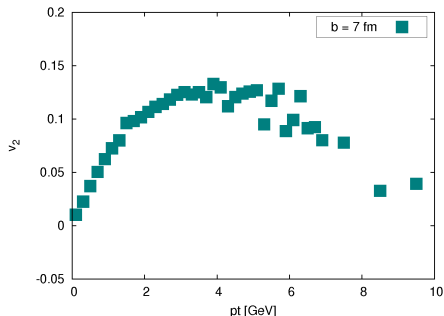


- Elliptic flow generated in BAMPS compatible with experiment.
- Suppression of gluon jets slightly stronger than found by Wicks et al. (Nucl.Phys.A784).

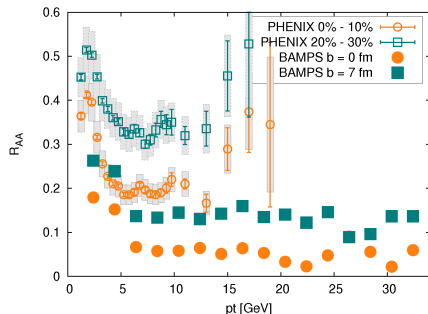
OF, Z. Xu, C. Greiner, PRL 102 (2009)

Non-Central R_{AA} and High- p_T Elliptic Flow

Differential v_2 for $b = 7$ fm

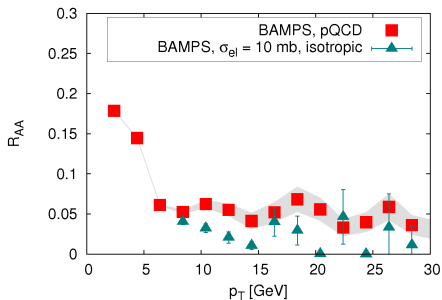


Gluon R_{AA} for $b = 0$ and $b = 7$ fm

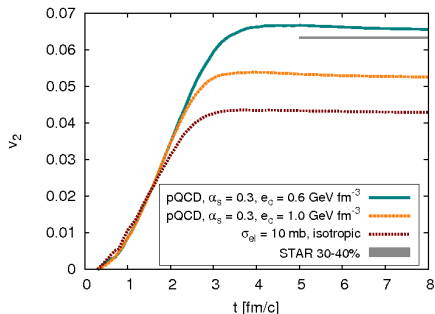


R_{AA} and v_2 from Elastic Collisions with Fixed σ

Gluon R_{AA} for fixed $\sigma = 10$ mb



Time evolution of v_2



Elastic collisions can't get R_{AA} and v_2 right simultaneously

For binary $\sigma = 10$ mb: R_{AA} too strong, v_2 still too weak

OF, Z. Xu, C. Greiner, PRL 102, 202301 (2009)

Parameters in BAMPS

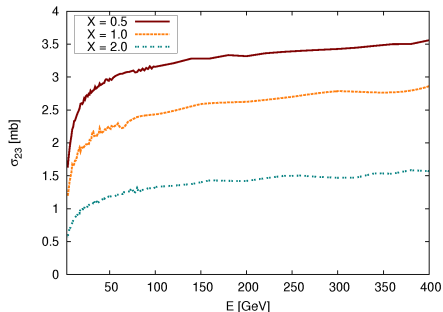
- Coupling strength α_s
- Critical freeze-out energy density ε_c
- LPM cut-off

The effective implementation of the LPM cut-off requires $\Lambda_g > \tau$.
Only qualitative argument, introduce factor X to test sensitivity.

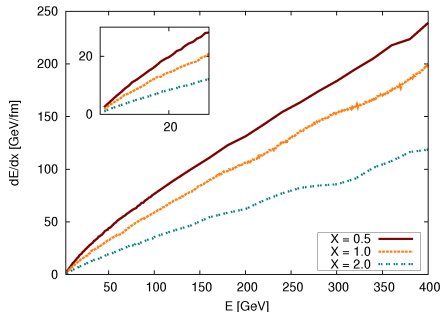
$$\Theta\left(k_{\perp} - \frac{\gamma}{\Lambda_g}\right) \rightarrow \Theta\left(k_{\perp} - X \frac{\gamma}{\Lambda_g}\right)$$

Sensitivity on the LPM Cut-Off

$\sigma_{gg \rightarrow ggg}$ for different X ($T = 400$ MeV)



dE/dx for different X ($T = 400$ MeV)

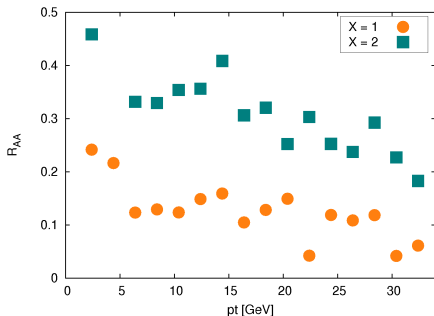


$$\sigma_{gg \rightarrow ggg} \sim \int dq_{\perp}^2 \int dk_{\perp}^2 \int y \dots \Theta \left(k_{\perp} - X \frac{\gamma}{\Lambda_g} \right)$$

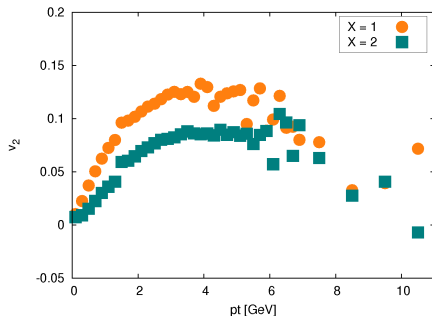
- Large X reduces total cross section
- Sampling of outgoing particles affected in non-trivial way
- Energy loss per collision only slightly affected, main contribution to the change in energy loss from change in σ .

Sensitivity on the LPM Cut-Off

R_{AA} for different X ($b = 7$ fm)



v_2 for different X ($b = 7$ fm)



- Partonic transport based on pQCD including $2 \rightarrow 3$ processes provides means of investigating medium and high- p_T physics within a common framework.
- Integrated gluon v_2 ok, gluon R_{AA} slightly below analytic reference

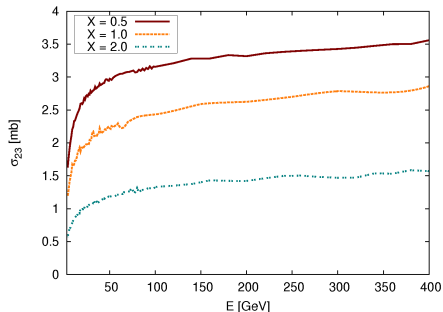
More quantitative and detailed studies needed

- More accurate identification $\langle b \rangle \leftrightarrow$ centrality class needed
- Need to include light quarks
- Need fragmentation (hadronization) scheme
- Need to investigate more (differential) observables

Thank you!

Sensitivity on the LPM Cut-Off

$\sigma_{gg \rightarrow ggg}$ for different X ($T = 400$ MeV)



$\Delta E_{gg \rightarrow ggg}$ for different X ($T = 400$ MeV)

