

Relativistic 3D Viscous Hydrodynamics

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Outline

- ✦ Hydrodynamic Theory
- ✦ Computational Solution and Its Verification
- ✦ Is 3D Necessary at Midrapidity?

The Hydrodynamic Problem

- ✦ Given:
 - ✦ An Initial Condition
 - ✦ Equation of State
- ✦ Produce:
 - ✦ Evolution
 - ✦ Observables

Problem Description

- Conservation of Stress-Energy:

$$\partial^\mu T_{\mu\nu} = 0$$
$$T_{\mu\nu} = (\epsilon + P)u_\mu u_\nu - P g_{\mu\nu} + \Pi_{\mu\nu}$$

- Effective Pressure Altered by Velocity Gradients:

$$\Pi_{ij}^{(NS)} = -\eta[\partial_i n_j + \partial_j n_i - \delta_{ij} \partial_k n_k] - \zeta \partial_k n_k$$

- Rapid Longitudinal Expansion yields
Reduced Effective Longitudinal Pressure.

Israel-Stewart Theory

- ✦ Israel-Stewart Allows Dynamical Effective Pressure.
- ✦ Deviations from NS decay exponentially

$$\partial_t(a_i/\alpha) = \frac{-(a_i - a_i^{(NS)})}{\tau_\eta \alpha} + \Omega_{ij}(a_j/\alpha)$$

$$\alpha^2 = \eta T s$$

Problem Description

- Projected Pressure Corrections:

$$a_1 = \Pi_{xx} - \Pi_{yy}$$

$$\sqrt{3}a_2 = \Pi_{xx} + \Pi_{yy} - 2\Pi_{zz}$$

$$a_3 = \Pi_{xy}, \quad a_4 = \Pi_{xz}, \quad a_5 = \Pi_{yz}$$

$$b = (1/3)(\Pi_{xx} + \Pi_{yy} + \Pi_{zz})$$

Features

- ✦ 3d Reflection Symmetry (up to 8x speed up).
- ✦ Multiple Time Integration Methods.
- ✦ Optional Bjorken Expansion.
- ✦ 40^3 grid integrates in 2 hours.

Basic Testing

- ✦ Total Entropy Conservation
 - ✦ $\sim 10^{-5}$
 - ✦ Scales with Grid Density/Time Step.
- ✦ Test Problems
 - ✦ Exponential Energy Density
 - ✦ TECHQM Standard Solutions (soon)

Exponential Test

✦ For initial condition: $\epsilon(x) = \epsilon_0 \cdot \exp[-x/R]; v(x) = 0$

✦ Energy Conservation: $\partial_\mu T^{\mu 0} = 0$

$$2\epsilon(1 + c_s^2)u\dot{u} + [u_0^2 + c_s^2(u_0^2 - 1)]\dot{\epsilon} = (1 + c_s^2)u_0u\frac{\epsilon}{R}$$

✦ Momentum Conservation:

$$\partial_\mu T^{\mu x} = 0$$

$$\frac{(1 + c_s^2)\epsilon}{u_0} [u^2 + u_0^2]\dot{u} + (1 + c_s^2)u_0u\dot{\epsilon} = [u^2(1 + c_s^2) - c_s^2] \frac{\epsilon}{R}$$

Exponential Test (u)

- Separable for u:

$$\dot{u} = \frac{u_0 c_s^2}{R(1 + c_s^2)(u_0^2 - c_s^2 u^2)}$$

- Analytically Integrable:

$$\frac{c_s^2}{R(1 + c_s^2)} t = \frac{1}{2} \left[(1 + c_s^2) \sinh^{-1}(u) - (c_s^2 - 1) u \sqrt{u^2 + 1} \right]$$

$$\frac{c_s^2}{R(1 + c_s^2)} t \approx u$$

- But not invertible.

Exponential Test (u)

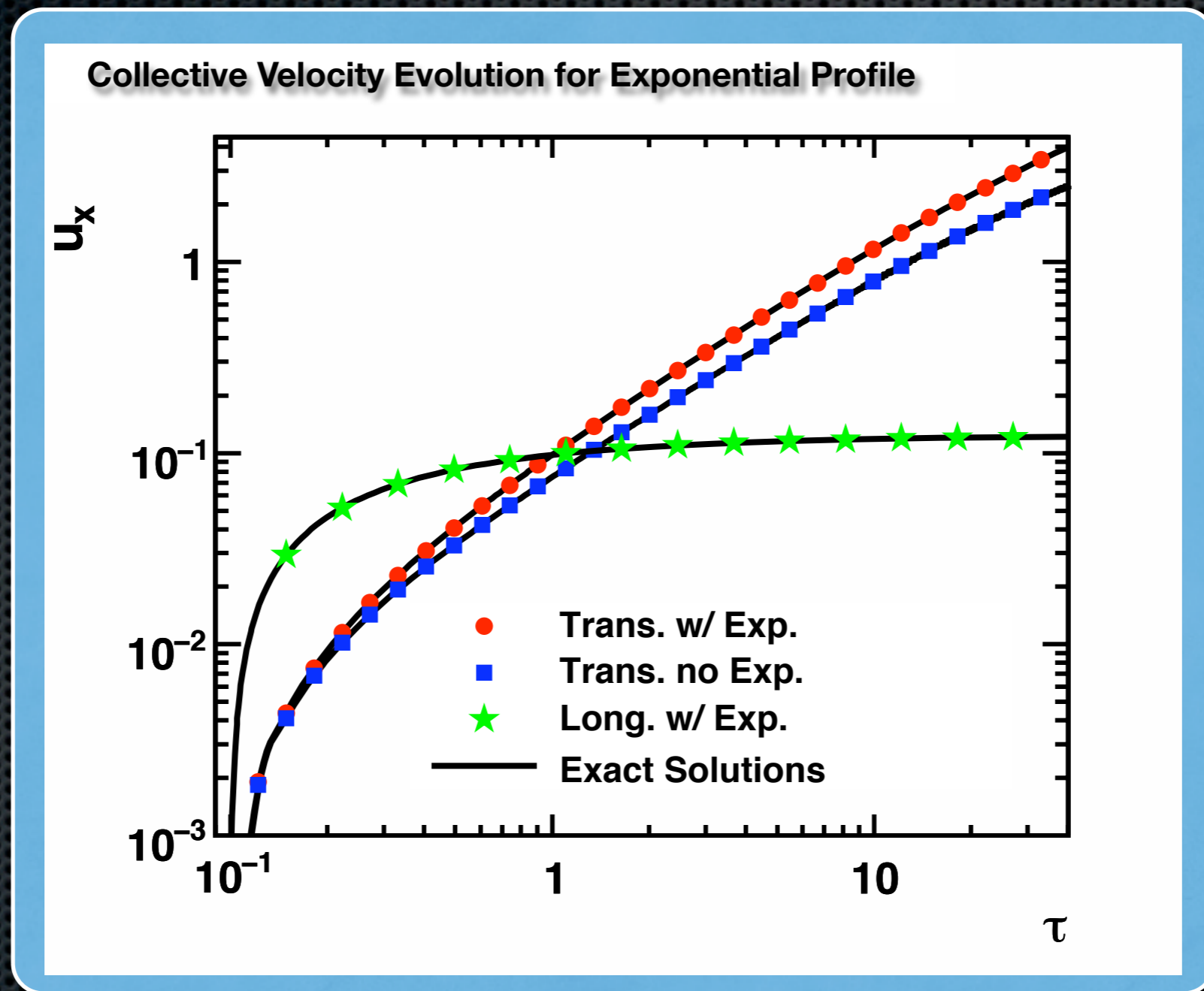
- Standard Collective Acceleration:

$$\dot{u} = \frac{u_0 c_s^2}{R(1 + c_s^2)(u_0^2 - c_s^2 u^2)}$$

- Modified by longitudinal expansion:

$$\dot{u} = \frac{u_0 c_s^2}{u_0^2 - c_s^2 u^2} \cdot \left[\frac{u_0 u}{\tau} + \frac{1}{R(1 + c_s^2)} \right]$$

Exponential Test (u)



Exponential Test (ϵ)

- We could have instead solved conservation equations for energy.

- With no Bjorken expansion:
$$\frac{\dot{\epsilon}}{\epsilon} = \frac{uu_0(1 - c_s^2)}{R(u^2(1 - c_s^2) + 1)}$$

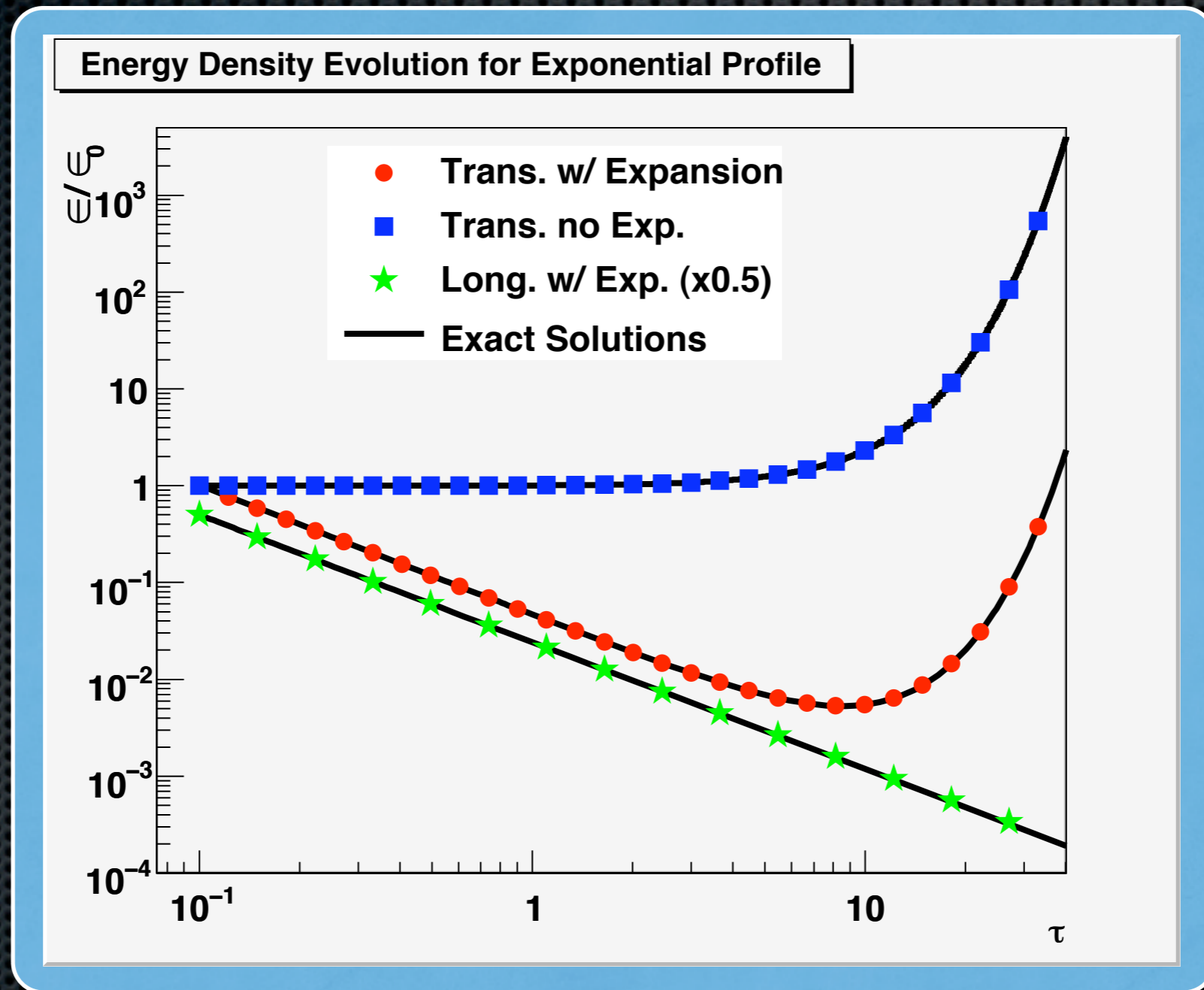
- Which may be integrated:

$$\epsilon = \epsilon_0 \cdot \exp \left[\frac{1 - c_s^4}{2c_s^2} u^2 \right]$$

- Or with Bjorken expansion:

$$\frac{\dot{\epsilon}}{\epsilon} = \frac{u_0}{u^2(1 - c_s^2) + 1} \cdot \left[\frac{(1 - c_s^2)u}{R} - \frac{(1 + c_s^2)u_0}{\tau} \right]$$

Energy Density Test

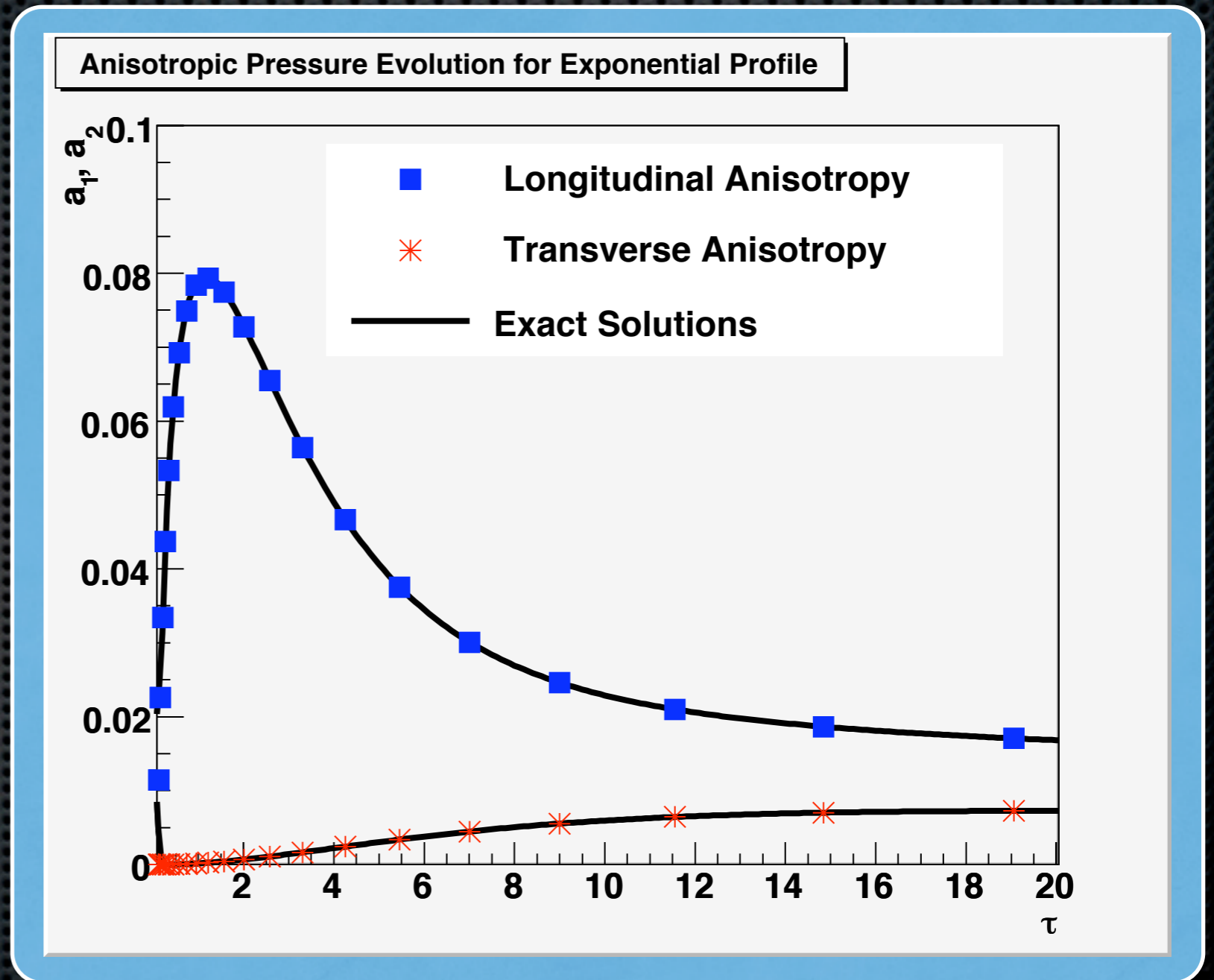


Exponential Test ($\pi_{\alpha\beta}$)

- ✦ Similar testing of the viscous deviations.

- ✦ For simplicity:

$$\eta = \epsilon/4\pi$$



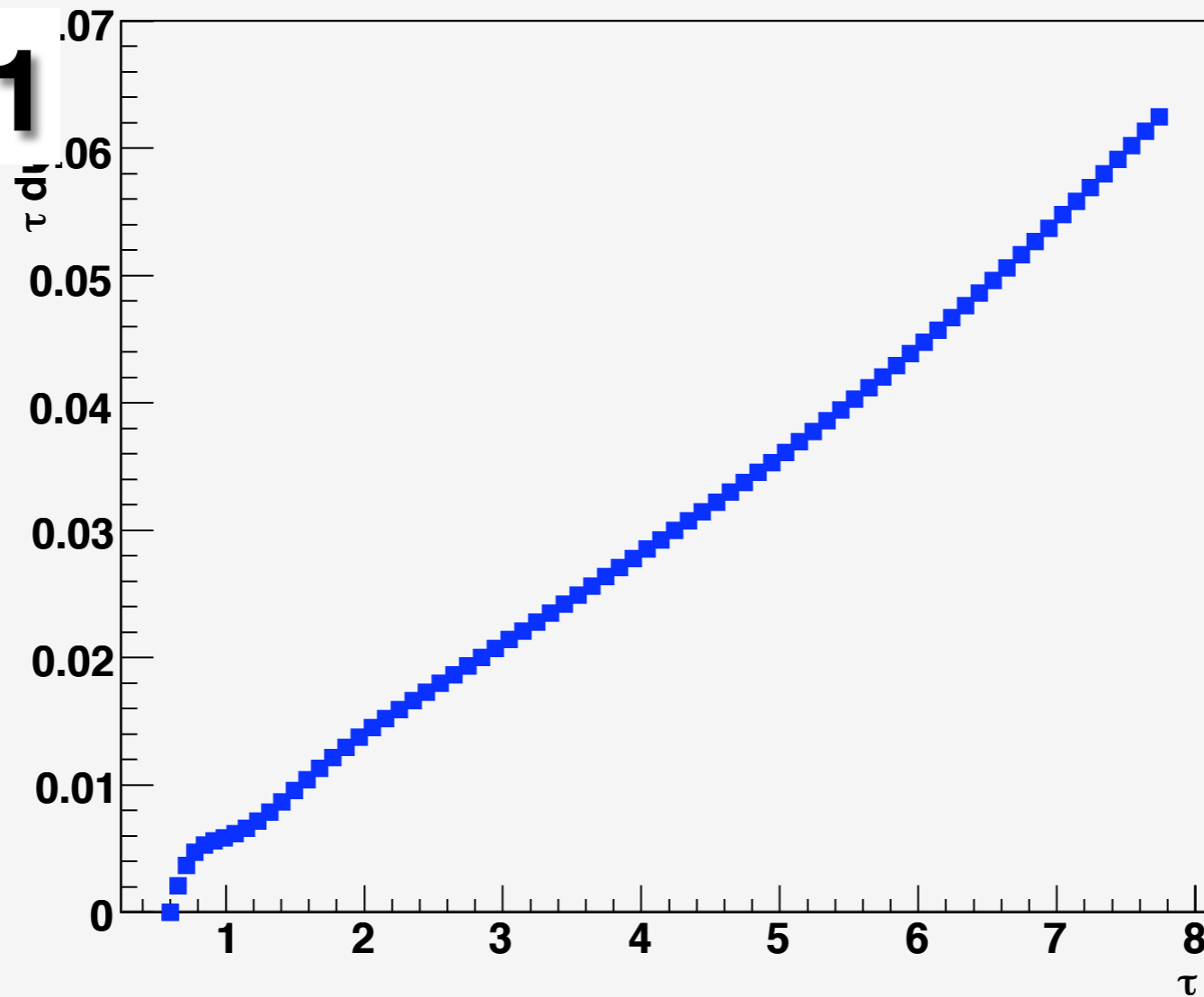
Is the 3rd dimension necessary?

- Used Ideal EOS - constant sound speed.
- $4\pi \cdot \frac{\eta}{s} = 2$.
- Initial conditions:
 - Woods-Saxon; quick fit to TECHQM tests.
 - Sharper Woods-Saxon in Eta.

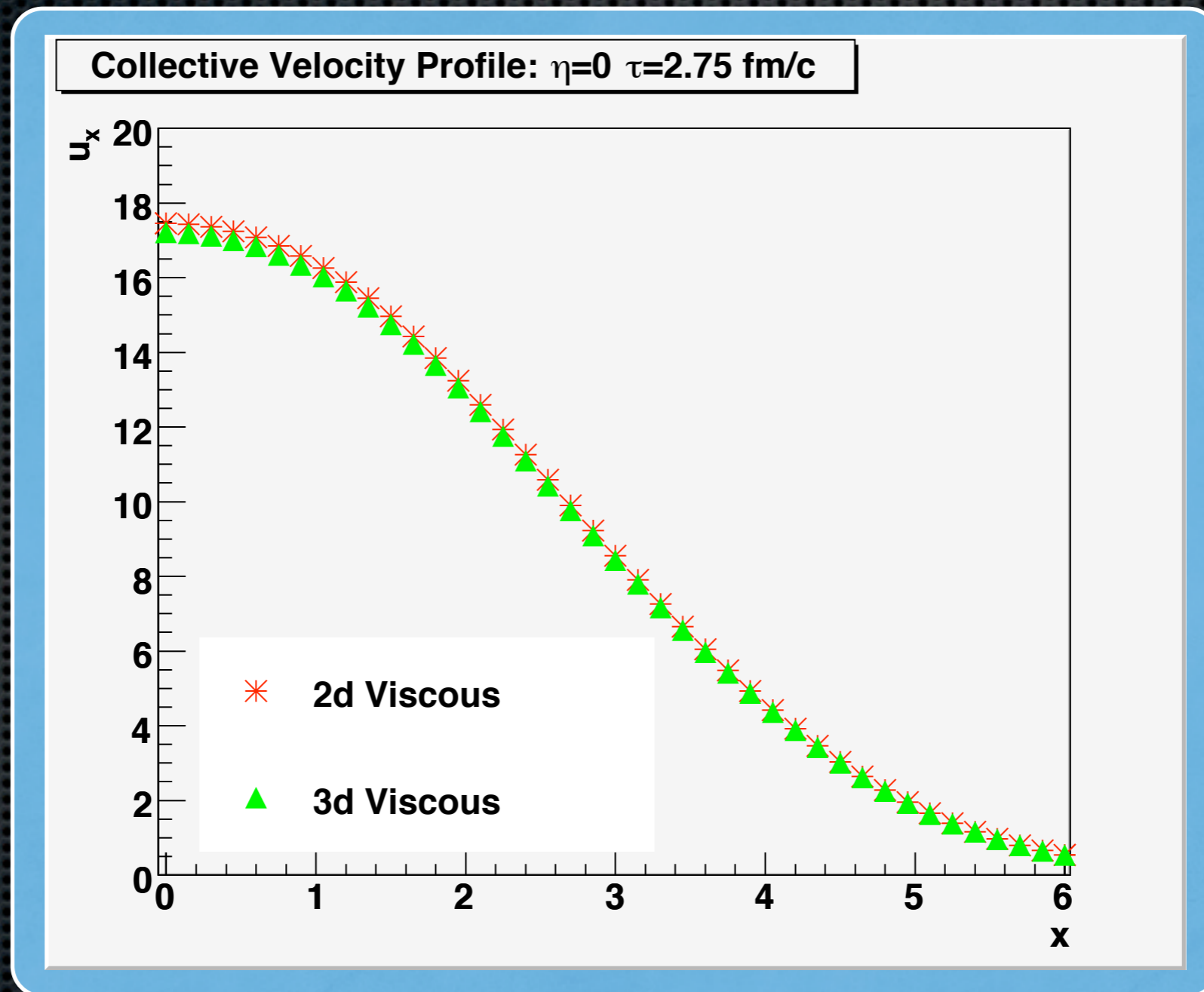
Effects of (3+1) - Long. Velocity Gradient

Central Longitudinal Velocity Gradient

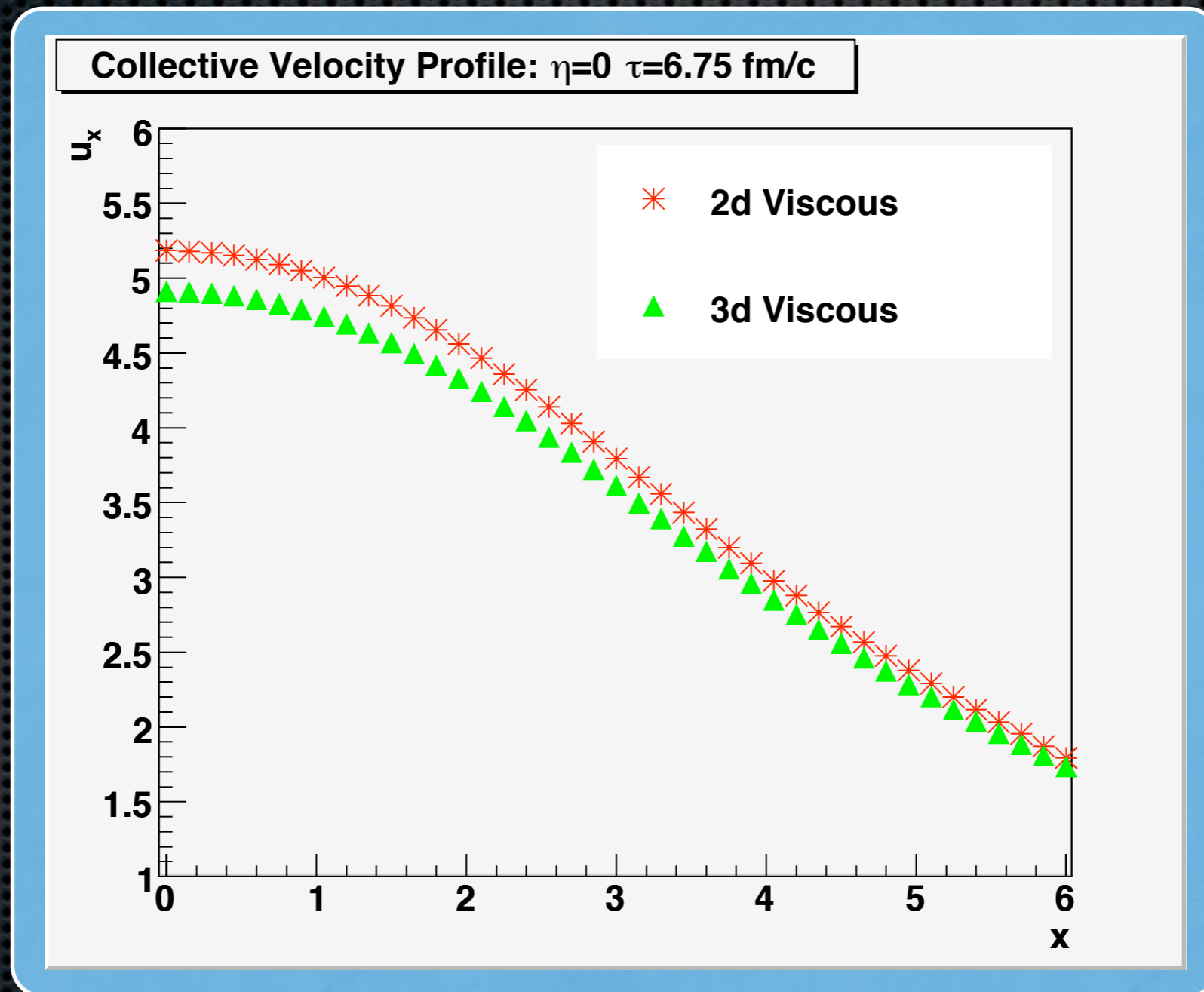
$\tau \frac{du_z}{dz} - 1$



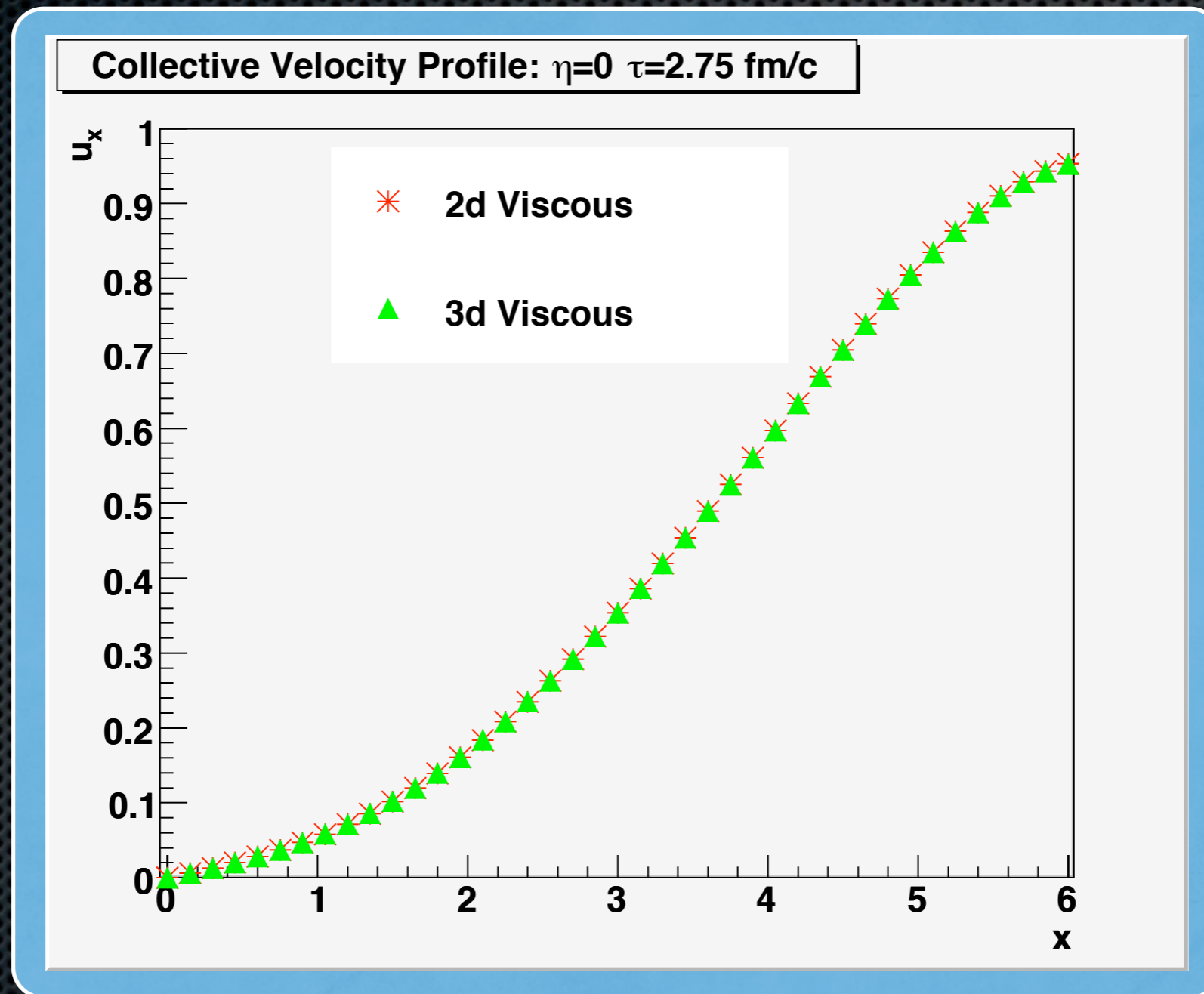
Effects of (3+1) - Energy Density Profiles



Effects of (3+1) - Energy Density Profiles



Effects of (3+1) - Velocity Profiles



Moving Forward

- ✦ Significant progress has been made toward a (3+1) dimensional Israel-Stewart code.
 - ✦ Additional testing required
 - ✦ TechQM tests
 - ✦ Shocks (Travel to Frankfurt in Spring)
 - ✦ Lattice EOS (R. Soltz, J. Newby and A. Glenn).