

Physics for Scientists & Engineers 2

Spring Semester 2005 Lecture 29

Review

- Consider a circuit consisting of an inductor L and a capacitor C
- The charge on the capacitor as a function of time is given by

 $q = q_{\max} \cos(\omega_0 t + \phi)$

- The current in the inductor as a function of time is given by

 $i = -i_{\max}\sin(\omega_0 t + \phi)$

 $\omega_0 = \sqrt{2}$

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• where ϕ is the phase and ω_0 is the angular frequency



• The energy stored in the magnetic field of the inductor L as a function of time is

$$U_B = \frac{L}{2} i_{\max}^2 \sin^2(\omega_0 t + \phi)$$

The total energy stored in the circuit is given by

$$U = U_E + U_B = \frac{q_{\text{max}}^2}{2C}$$

RLC Circuit

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 Now let's consider a single loop circuit that has a capacitor C and an inductance L with an added resistance R



 We observed that the energy of a circuit with a capacitor and and an inductor

remains constant and that the energy translated from electric to magnetic and back gain with no losses

- If there is a resistance in the circuit, the current flow in the circuit will produce ohmic losses to heat
- Thus the energy of the circuit will decrease because of these losses



RLC Circuit (2)



The rate of energy loss is given by

$$\frac{dU}{dt} = -i^2 R$$

• We can rewrite the change in energy of the circuit as a function time as

$$\frac{dU}{dt} = \frac{d}{dt} \left(U_E + U_B \right) = \frac{q}{C} \frac{dq}{dt} + Li \frac{di}{dt} = -i^2 R$$

Remembering that i = dq/dt and di/dt = d²q/dt² we can write

$$\frac{q}{C}\frac{dq}{dt} + Li\frac{di}{dt} + i^2R = \frac{q}{C}\frac{dq}{dt} + L\frac{dq}{dt}\frac{d^2q}{dt^2} + \left(\frac{dq}{dt}\right)^2R = 0$$

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RLC Circuit (3)



- We can then write the differential equation $L\frac{d^2q}{dt^2} + \frac{dq}{dt}R + \frac{q}{C} = 0$
- The solution of this differential equation is $q = q_{\max} e^{-\frac{Rt}{2L}} \cos(\omega t + \phi)$
- where

$$\omega = \sqrt{\omega_0^2 - \left(\frac{R}{2L}\right)^2} \qquad \qquad \omega_0 = \frac{1}{\sqrt{LC}}$$

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 The charge varies sinusoidally with but the amplitude is damped out with time

RLC Circuit (5)

- After some time, no charge remains in the circuit
- We can study the energy in the circuit as a function of time by calculating the energy stored in the electric field of the capacitor

$$U_{E} = \frac{1}{2} \frac{q^{2}}{C} = \frac{1}{2} \frac{\left(q_{\max}e^{-\frac{Rt}{2L}}\cos(\omega t + \phi)\right)^{2}}{C} = \frac{q_{\max}^{2}}{2C}e^{-\frac{Rt}{L}}\cos^{2}(\omega t + \phi)$$

 We can see that the energy stored in the capacitor decreases exponentially and oscillates in time

RLC Circuit (4)



- Now consider a single loop circuit that contains a capacitor, an inductor and a resistor
- If we charge the capacitor then hook it up to the circuit, we will observe a charge in the circuit that varies sinusoidally with time and while at the same time decreasing in amplitude
- This behavior with time is illustrated below



Alternating Current

 Now we consider a single loop circuit containing a capacitor, an inductor, a resistor, and a source of emf



- This source of emf is capable producing
 a time varying voltage as opposed the
 sources of emf we have studied in previous chapters
- We will assume that this source of emf provides a sinusoidal voltage as a function of time given by
 - $V_{emf} = V_{\max} \sin \omega t$
- where ω is that angular frequency of the emf and V_{\max} is the amplitude or maximum value of the emf

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Circuit with Resistor

- To begin our analysis of *RLC* circuits, let's start with a circuit containing only a resistor and a source of time-varying emf as shown to the right

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- Applying Kirchhof's loop rule to this circuit we get $V_{\rm emf} v_{\rm R} = 0 \implies V_{\rm emf} = v_{\rm R}$
- where v_R is the voltage drop across the resistor
- Substituting into our expression for the emf as a function of time we get

 $v_R = V_R \sin \omega t$

Remembering Ohm's Law, V = iR, we get

$$\bar{v}_R = \frac{v_R}{R} = I_R \sin \omega t$$

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Alternating Current (2)



- The current induced in the circuit will also vary sinusoidally with time
- This time-varying current is called alternating current
- However, this current may not always remain in phase with the time-varying emf
- We can express the induced current as $i = I \sin(\omega t \phi)$
- where the angular frequency of the time-varying current is the same as the driving emf but the phase ϕ is not zero
- Note that traditionally the phase enters here with a negative sign
- Thus the voltage and the current in the circuit are not necessarily in phase



Circuit with Resistor (2)



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- Thus we can relate the current amplitude and the voltage amplitude by $V_{\rm R} = I_{\rm R} R$
- We can represent the time varying current by a phasor I_R and the timevarying voltage by a phasor V_R as shown below



• The current flowing through the resistor and the voltage across the resistor are in phase, which means that the phase difference between the current and the voltage is zero

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Circuit with Capacitor



- Now let's address a circuit that contains a capacitor and a time varying emf as shown to the right
- The voltage across the capacitor is given by Kirchhof's loop rule $v_c = V_c \sin \omega t$
- Remembering that q = CV for a capacitor we can write $q = Cv_c = CV_c \sin \omega t$
- We would like to know the current as a function of time rather than the charge so we can write

$$i_C = \frac{dq}{dt} = \frac{d(CV_C \sin \omega t)}{dt} = \omega CV_C \cos \omega t$$

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Circuit with Capacitor (2)



 We can rewrite the last equation by defining a quantity that is similar to resistance and is called the capacitive reactance

$$X_C = \frac{1}{\omega C}$$

Which allows us to write

$$i_C = \frac{V_C}{X_C} \cos \omega t$$

• We can now express the current in the circuit as

$$i_C = I_C \cos \omega t = I_C \sin (\omega t + 90^\circ)$$

- We can see that the current and the time varying emf are out of phase by 90°

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Circuit with Capacitor (3)



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 We can represent the time varying current by a phasor I_c and the time-varying voltage by a phasor V_c as shown below



 The current flowing this circuit with only a capacitor is similar to the expression for the current flowing in a circuit with only a resistor except that the current is out of phase with the emf by 90°

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 We can also see that the amplitude of voltage across the capacitor and the amplitude of current in the capacitor are related by

 $V_C = I_C X_C$

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- This equation resembles Ohm's Law with the capacitive reactance replacing the resistance
- One major difference between the capacitive reactance and the resistance is that the capacitive reactance depends on the angular frequency of the time-varying emf

Circuit with Inductor



- Now let's consider a circuit with a source of time-varying emf and an inductor as shown to the right
- We can again apply Kirchhof's Loop Rule to this circuit to obtain the voltage across the inductor as

 $v_L = V_L \sin \omega t$

• A changing current in an inductor will induce an emf given by $v_L = L \frac{di_L}{dt}$

$$L\frac{di_L}{dt} = V_L \sin \omega t \implies \frac{di_L}{dt} = \frac{V_L}{L} \sin \omega t$$

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Circuit with Inductor (2)



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• We are interested in the current rather than its time derivative so we integrate

$$i_L = \int \frac{di_L}{dt} dt = \int \frac{V_L}{L} \sin \omega t dt = -\frac{V_L}{\omega L} \cos \omega t$$

• We define inductive reactance as

 $X_L = \omega L$

- which, like the capacitive reactance, is similar to a resistance
- We can then write
 - $v_L = i_L X_L$
- which again resembles Ohm's Law except that the inductive reactance depends on the angular frequency of the time-varying emf



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Circuit with Inductor (3)



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- The current in the inductor can then be written as $i_L = -\frac{v_L}{V} \cos \omega t = -I_L \cos \omega t = I_L \sin(\omega t - 90^\circ)$
- Thus the current flowing in a circuit with an inductor and a source timevarying emf will be -90° out of phase with the emf



• We can write the relationship between the amplitude of the current and the amplitude of the voltage as

 $V_L = I_L X_L$

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