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# PHY481: Electromagnetism

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Curl and divergence of  $\mathbf{E}$   
& a bit more vector mathematics

## Problem 2.27

$$\mathbf{F}(\mathbf{x}) = \mathbf{x} / r^3 = \hat{\mathbf{r}} / r^2$$

Integral over sphere radius a

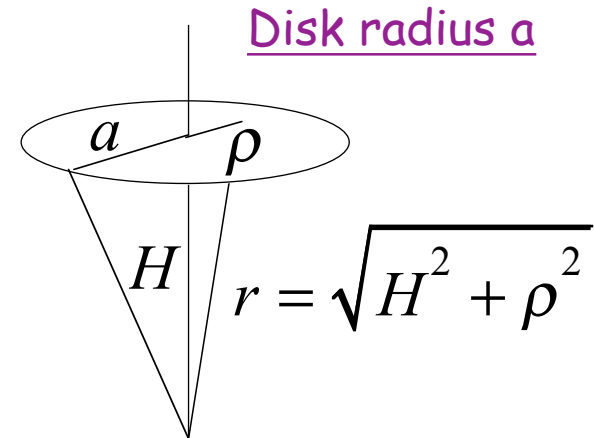
$$d\mathbf{A} = \hat{\mathbf{r}} r^2 \sin\theta d\theta d\phi$$

$$\oint_S \mathbf{F} \cdot d\mathbf{A} = \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi = 2\pi(-\cos\theta) \Big|_0^\pi = 4\pi$$

Prep for integral over Disk

$$d\mathbf{A} = \hat{\mathbf{k}} \rho d\rho d\phi = (\hat{\mathbf{r}} \cos\theta - \hat{\boldsymbol{\theta}} \sin\theta) \rho d\rho d\phi$$

$$r = \sqrt{H^2 + \rho^2}; \quad \mathbf{F} = \frac{\hat{\mathbf{r}}}{r^2} = \frac{\hat{\mathbf{r}}}{H^2 + \rho^2}; \quad \cos\theta = \frac{H}{\sqrt{H^2 + \rho^2}}$$



Integral over Disk radius a

$$\begin{aligned} \oint_S \mathbf{F} \cdot d\mathbf{A} &= H \int_0^a \frac{\rho}{(H^2 + \rho^2)^{3/2}} d\rho \int_0^{2\pi} d\phi \\ &= 2\pi \left( 1 - \frac{H}{(H^2 + a^2)^{1/2}} \right) \end{aligned}$$

Integral when H = a

$$\oint_S \mathbf{F} \cdot d\mathbf{A} = 2\pi \left( 1 - \frac{1}{\sqrt{2}} \right)$$

# Electric field of charge distributions

- Charge distributions

Distribution: Volume

Charge density:  $\rho(\mathbf{x}')$

Total Charge:  $Q = \int \rho(\mathbf{x}') d^3x'$

Surface

$\sigma(\mathbf{x}')$

$Q = \int \sigma(\mathbf{x}') d^2x'$

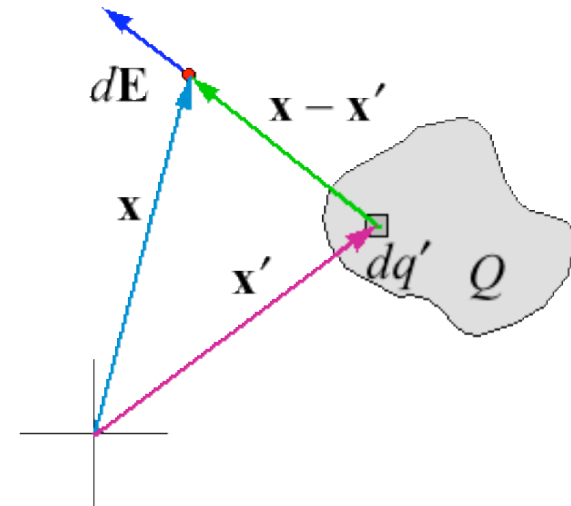
Line

$\lambda(\mathbf{x}')$

$Q = \int \lambda(\mathbf{x}') dx'$

- Integration over charge distribution

$$\mathbf{E}(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\mathbf{x} - \mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|^3} \rho(\mathbf{x}') d^3x'$$



## Old way

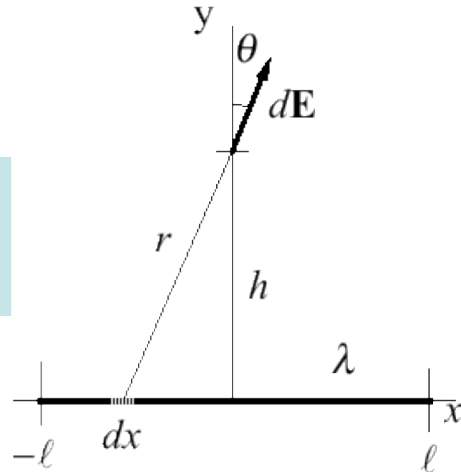
## Line of charge

## New way

$$r = (h^2 + x^2)^{1/2}$$

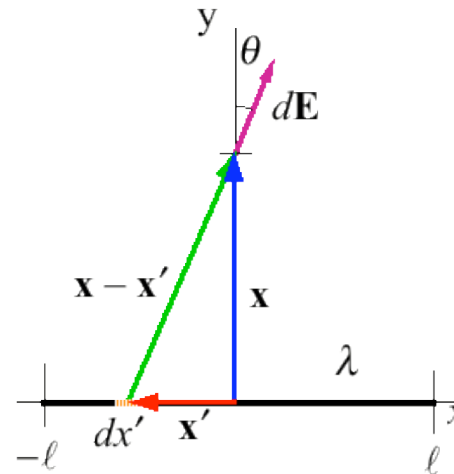
$$\cos \theta = \frac{h}{(h^2 + x^2)^{1/2}}$$

$$dq = \lambda dx$$



$$E_y = \int dE_y = \int dE \cos \theta = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r^2} \cos \theta$$

$$\begin{aligned} &= \frac{1}{4\pi\epsilon_0} \int_{-\ell}^{\ell} \frac{\lambda dx}{(h^2 + x^2)^{3/2}} \cos \theta \\ &= \frac{1}{4\pi\epsilon_0} \int_{-\ell}^{\ell} \frac{\lambda dx}{(h^2 + x^2)^{3/2}} \left( \frac{h}{(h^2 + x^2)^{1/2}} \right) \\ E_y &= \frac{\lambda h}{4\pi\epsilon_0} \int_{-\ell}^{\ell} (h^2 + x^2)^{-3/2} dx \end{aligned}$$



$$\begin{aligned} \mathbf{x} &= h \hat{\mathbf{j}} \\ \mathbf{x}' &= x' \hat{\mathbf{i}} \\ \mathbf{x} - \mathbf{x}' &= h \hat{\mathbf{j}} - x' \hat{\mathbf{i}} \end{aligned}$$

$$|\mathbf{x} - \mathbf{x}'| = (h^2 + x'^2)^{1/2}$$

$$\rho(\mathbf{x}') d^3 x' = \lambda dx'$$

$$\mathbf{E}(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\mathbf{x} - \mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|^3} \rho(\mathbf{x}') d^3 x'$$

$$\begin{aligned} &= \frac{1}{4\pi\epsilon_0} \int_{-\ell}^{\ell} \frac{h \hat{\mathbf{j}} - x' \hat{\mathbf{i}}}{(h^2 + x'^2)^{3/2}} \lambda dx' \\ \mathbf{E} &= \hat{\mathbf{j}} \frac{h\lambda}{4\pi\epsilon_0} \int_{-\ell}^{\ell} (h^2 + x'^2)^{-3/2} dx' \end{aligned}$$

Angular dependence  
automatically included !

## Curl and divergence of E

- **Prove** differential form of Gauss's Law

Gauss's Law:  $\nabla \cdot \mathbf{E}(\mathbf{x}) = \frac{\rho(\mathbf{x})}{\epsilon_0}$  True even for  $\mathbf{E}(\mathbf{x}, t)$  with  $\nabla \times \mathbf{E}(\mathbf{x}) = 0$  True only for static charge

- Will need definition of the **scalar potential** function  $V(\mathbf{x})$

$$V(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{x}') d^3x'}{|\mathbf{x} - \mathbf{x}'|}$$

$\mathbf{x}$  = evaluation coordinates

$\mathbf{x}'$  = source positions

$|\mathbf{x} - \mathbf{x}'|$  (not  $|\mathbf{x}|^2$  or  $|\mathbf{x}|^3$ )

- Will need **Gradient** of  $V$

- Derivatives are with respect to  $\mathbf{x}$ , **not**  $\mathbf{x}'$ :

$$\nabla = \hat{\mathbf{e}}_i \frac{\partial}{\partial x_i}$$

- Bring Del into the integral over  $d^3x'$

$$\nabla V(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \nabla \left( \int \frac{\rho(\mathbf{x}') d^3x'}{|\mathbf{x} - \mathbf{x}'|} \right) = \frac{1}{4\pi\epsilon_0} \int \rho(\mathbf{x}') \nabla \left( \frac{1}{|\mathbf{x} - \mathbf{x}'|} \right) d^3x'$$

## Gradient of $1/|\mathbf{x}-\mathbf{x}'|$

Since:

$$\begin{aligned}\nabla\left(\frac{1}{|\mathbf{x}-\mathbf{x}'|}\right) &= \frac{-1}{|\mathbf{x}-\mathbf{x}'|^2} \nabla|\mathbf{x}-\mathbf{x}'| \\ &= -\frac{\mathbf{x}-\mathbf{x}'}{|\mathbf{x}-\mathbf{x}'|^3}\end{aligned}$$

, then

$$\begin{aligned}\nabla V(\mathbf{x}) &= \frac{1}{4\pi\epsilon_0} \int \nabla\left(\frac{1}{|\mathbf{x}-\mathbf{x}'|}\right) \rho(\mathbf{x}') d^3x' \\ &= -\frac{1}{4\pi\epsilon_0} \int \frac{\mathbf{x}-\mathbf{x}'}{|\mathbf{x}-\mathbf{x}'|^3} \rho(\mathbf{x}') d^3x' \\ &= -\mathbf{E}(\mathbf{x})\end{aligned}$$

Using chain rule

$$\frac{\partial}{\partial x_i} f[g(x)] = f'[g(x)] \frac{\partial g(x)}{\partial x_i}$$

$$\begin{aligned}|\mathbf{x}-\mathbf{x}'| &= \sqrt{(x_i-x'_i)^2} \\ \nabla|\mathbf{x}-\mathbf{x}'| &= \frac{2(x_i-x'_i)\mathbf{e}_i}{2\sqrt{(x_i-x'_i)^2}} = \frac{\mathbf{x}-\mathbf{x}'}{|\mathbf{x}-\mathbf{x}'|}\end{aligned}$$

Presented as

$$\mathbf{E}(\mathbf{x}) = -\nabla V(\mathbf{x})$$

## Del squared of $V$

- Still trying to **prove** Gauss's Law !

$$\nabla \cdot \mathbf{E}(\mathbf{x}) = \frac{\rho(\mathbf{x})}{\epsilon_0}$$

- We now have available  $\mathbf{E}(\mathbf{x}) = -\nabla V(\mathbf{x})$

$$\nabla \cdot \mathbf{E}(\mathbf{x}) = \nabla \cdot (-\nabla V) = -\nabla^2 V \quad (-) \text{ Laplacian of } V$$

$$\nabla \cdot \mathbf{E}(\mathbf{x}) = -\nabla^2 V(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int -\nabla^2 \left[ \frac{1}{|\mathbf{x} - \mathbf{x}'|} \right] \rho(\mathbf{x}') d^3x'$$

- Two will be the same if we can show that

$$\int_{Vol.} -\nabla^2 \left[ \frac{1}{|\mathbf{x} - \mathbf{x}'|} \right] d^3x' = 4\pi \text{ if Volume includes } \mathbf{x}' = \mathbf{x}$$

Involves Dirac's delta function

## Laplacian of $1/|\mathbf{x}-\mathbf{x}'|$

- To **prove** Gauss's Law we need to show

$$\int_{Vol.} -\nabla^2 \left[ \frac{1}{|\mathbf{x} - \mathbf{x}'|} \right] d^3x' = 4\pi \text{ if Volume includes } \mathbf{x}' = \mathbf{x}$$

- Point source at  $\mathbf{x}' = 0$ , and use spherical coordinates  $\mathbf{x} = \mathbf{r}$

$$\begin{aligned} \nabla^2 \left[ \frac{1}{|\mathbf{x}|} \right] &= \nabla^2 \left[ \frac{1}{r} \right] = \frac{d^2}{dr^2} \left[ \frac{1}{r} \right] + \frac{2}{r} \frac{d}{dr} \left[ \frac{1}{r} \right] \\ &= \frac{2}{r^3} - \frac{2}{r^3} = 0 \text{ everywhere?} \end{aligned}$$

What about  
at  $r = 0$ ?

- Strategy to evaluate at the singular point,  $r = 0$

Use Gauss's theorem & a little sphere radius  $R$  around  $r = 0$



## Strange Laplacian of $1/|\mathbf{x}-\mathbf{x}'|$

- Value of  $\nabla^2 \left[ \frac{1}{r} \right] = 0$  But perhaps not at  $r = 0$  !

- Use Gauss's theorem to evaluate at  $r = 0$

$$\int_V \nabla \cdot \mathbf{F} d^3x = \int_S \mathbf{F} \cdot d\mathbf{A}$$

$$\text{with } \mathbf{F} = \nabla \left[ \frac{1}{r} \right]$$

$$\int_V \nabla^2 \left[ \frac{1}{r} \right] d^3x = \int_S \nabla \left[ \frac{1}{r} \right] \cdot \hat{\mathbf{r}} dA$$

$$\nabla \left[ \frac{1}{r} \right] = \frac{\partial}{\partial r} \left[ \frac{1}{r} \right] \hat{\mathbf{r}} = \frac{-1}{r^2} \hat{\mathbf{r}}$$
$$dA = R^2 \sin \theta d\theta d\phi$$

- Evaluate right side with sphere, radius R around origin

$$\int_V \nabla^2 \left[ \frac{1}{r} \right] d^3x = \frac{-1}{R^2} \int_0^\pi R^2 \sin \theta d\theta \int_0^{2\pi} d\phi = -4\pi \quad (\text{pretty amazing})$$

Does not depend on  
sphere radius R !

$$\int_{Vol.} -\nabla^2 \left[ \frac{1}{r} \right] d^3x' = 4\pi \text{ if Volume includes } r = 0$$

# Dirac delta function

- Strange behavior of the Laplacian of  $1/r$

$$\int_{Vol.} -\nabla^2 \left[ \frac{1}{|\mathbf{x}|} \right] d^3x' = \begin{cases} 4\pi & \text{if Volume includes } \mathbf{x} = 0 \\ 0 & \text{elsewhere} \end{cases}$$

- Define Dirac delta function (very useful in advanced physics)

$$\delta^3(\mathbf{x}) = -\frac{1}{4\pi} \nabla^2 \left[ \frac{1}{|\mathbf{x}|} \right] \quad \int_{Vol.} \delta^3(\mathbf{x}) d^3x' = \begin{cases} 1 & \text{if Volume includes } \mathbf{x} = 0 \\ 0 & \text{elsewhere} \end{cases}$$

- Finally, we prove Gauss's Law

delta function  
picks out value  
of  $\rho$  at point  
where  $\mathbf{x} = \mathbf{x}'$

$$\nabla \cdot \mathbf{E} = -\nabla^2 V(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int -\nabla^2 \left[ \frac{1}{|\mathbf{x} - \mathbf{x}'|} \right] \rho(\mathbf{x}') d^3x'$$

$$\nabla \cdot \mathbf{E} = \frac{1}{\epsilon_0} \int \delta^3(\mathbf{x} - \mathbf{x}') \rho(\mathbf{x}') d^3x' = \frac{\rho(\mathbf{x})}{\epsilon_0}$$

# Green's functions

- Poisson's equation and a general solution

$$-\nabla^2 V(\mathbf{x}) = \frac{\rho(\mathbf{x})}{\epsilon_0}$$

$$V(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{x}') d^3x'}{|\mathbf{x} - \mathbf{x}'|}$$

- Green's function of a differential operator
  - The Laplacian is just one of many important differential operators

- Consider a unit point source at  $\mathbf{x}'$   $\rho_0(\mathbf{x}) = \delta^3(\mathbf{x} - \mathbf{x}')$

- The Green's function  $G$  for the Laplacian (which we have just found) is the solution for the potential due a unit point source

$$-\nabla^2 G(\mathbf{x} - \mathbf{x}') = \delta^3(\mathbf{x} - \mathbf{x}')$$

$$G(\mathbf{x} - \mathbf{x}') = \frac{1}{4\pi |\mathbf{x} - \mathbf{x}'|}$$

- The potential  $V(\mathbf{x})$  for an arbitrary source distribution is then

$$V(\mathbf{x}) = \int G(\mathbf{x} - \mathbf{x}') \rho(\mathbf{x}') d^3x'$$