
PHY481: Electromagnetism

Exam 1 Solutions
Charge on a conductor

Exam 1

1) Stokes's theorem: $\oint_S (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} dA = \oint_C \mathbf{F} \cdot d\boldsymbol{\ell}$

$$\begin{aligned}\mathbf{F}(\mathbf{x}) &= \hat{\mathbf{i}} \times (x\hat{\mathbf{i}} + y\hat{\mathbf{j}} + z\hat{\mathbf{k}}) \\ &= y\hat{\mathbf{k}} - z\hat{\mathbf{j}}\end{aligned}$$

$$1) d\boldsymbol{\ell} = \hat{\mathbf{j}} dy, z = 0, \quad 2) d\boldsymbol{\ell} = \hat{\mathbf{k}} dz, y = 1$$

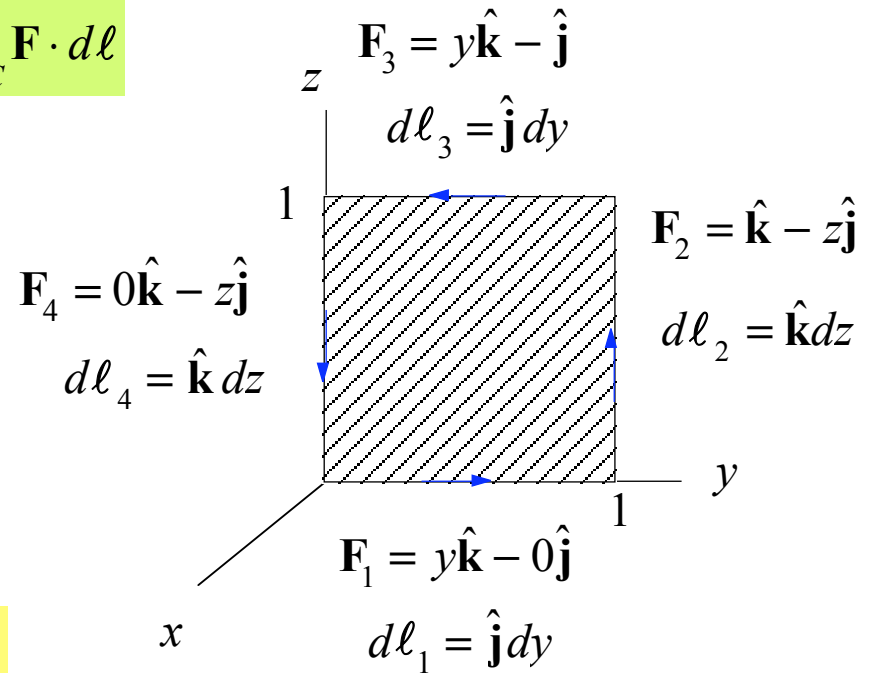
$$3) d\boldsymbol{\ell} = \hat{\mathbf{j}} dy, z = 1, \quad 4) d\boldsymbol{\ell} = \hat{\mathbf{k}} dz, y = 0$$

a) $\int_1 \mathbf{F} \cdot d\boldsymbol{\ell} = -0 \int_0^1 dy = \underline{0}; \quad \int_2 \mathbf{F} \cdot d\boldsymbol{\ell} = 1 \int_0^1 dz = \underline{1};$

$$\int_3 \mathbf{F} \cdot d\boldsymbol{\ell} = -1 \int_1^0 dy = \underline{1}; \quad \int_4 \mathbf{F} \cdot d\boldsymbol{\ell} = -0 \int_1^0 dz = \underline{0}$$

b) $\int_C \mathbf{F} \cdot d\boldsymbol{\ell} = \underline{2}$ $\nabla \times \mathbf{F} = \hat{\mathbf{i}} \left(\frac{\partial F_z}{\partial y} - \frac{\partial F_y}{\partial z} \right) + \hat{\mathbf{j}} \left(\frac{\partial F_x}{\partial z} - \frac{\partial F_z}{\partial x} \right) + \hat{\mathbf{k}} \left(\frac{\partial F_y}{\partial x} - \frac{\partial F_x}{\partial y} \right) = \hat{\mathbf{i}}(2) + \hat{\mathbf{j}}(0) + \hat{\mathbf{k}}(0)$

c) $\hat{\mathbf{n}} = \hat{\mathbf{i}}; \quad dA = dx dy \quad ; \quad \int_S (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} dA = \underline{2} \quad \therefore \quad \int_S (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} dA = \int_C \mathbf{F} \cdot d\boldsymbol{\ell} = \underline{2}$



Exam 1 (cont'd)

2)

a) integral to obtain scalar potential from charge distribution

$$V(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{x}') d^3x'}{|\mathbf{x} - \mathbf{x}'|}$$

b) differential equation relating charge density and the electric field

$$\nabla \cdot \mathbf{E}(\mathbf{x}) = \frac{\rho(\mathbf{x})}{\epsilon_0}$$

c) differential equation relating electric field and potential

$$\mathbf{E}(\mathbf{x}) = -\nabla V(\mathbf{x})$$

d) integral to obtain electric field from the charge distribution

$$\mathbf{E}(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\mathbf{x} - \mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|^3} \rho(\mathbf{x}') d^3x'$$

Exam 1 (cont'd)

3a) Gauss's law: $\oint_S \mathbf{E} \cdot d\mathbf{A} = q_{encl} / \epsilon_0$

b) $\mathbf{E}_T(\mathbf{x}) = E_0 \hat{\mathbf{k}}; \quad \mathbf{E}_B(\mathbf{x}) = -E_0 \hat{\mathbf{k}}$
 $\hat{\mathbf{n}}_T = \hat{\mathbf{k}}; \quad \hat{\mathbf{n}}_B = -\hat{\mathbf{k}}$

$$\int_S \mathbf{E} \cdot d\mathbf{A} = \int_T E_0 \hat{\mathbf{k}} \cdot \hat{\mathbf{k}} dA + \int_B (-E_0 \hat{\mathbf{k}}) \cdot (-\hat{\mathbf{k}} dA)$$

$$2E_0 A = q_{encl} / \epsilon_0 = \sigma A / \epsilon_0$$

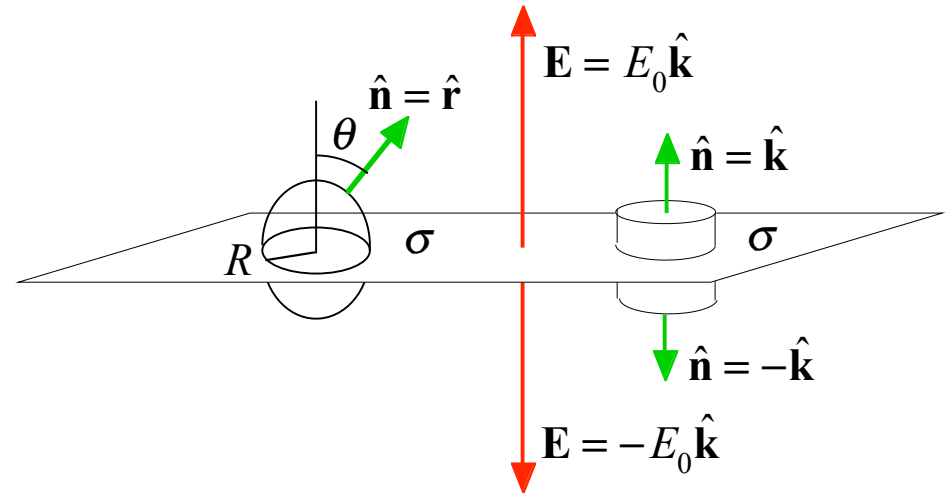
$$E_0 = \sigma / 2\epsilon_0$$

c) $\hat{\mathbf{n}} = \hat{\mathbf{r}}; \quad dA = R^2 \sin\theta d\theta d\phi; \quad q_{encl} = \sigma\pi R^2$

$$\int_S \mathbf{E} \cdot d\mathbf{A} = \int_T (E_0 \hat{\mathbf{k}}) \cdot \hat{\mathbf{r}} dA + \int_B (-E_0 \hat{\mathbf{k}}) \cdot \hat{\mathbf{r}} dA = 2\pi E_0 R^2 \left[\int_0^{\pi/2} d\theta \cos\theta \sin\theta - \int_{\pi/2}^{\pi} d\theta \cos\theta \sin\theta \right]$$

$$= 2\pi E_0 R^2 \left[\int_0^1 x dx - \int_1^0 x dx \right] = 2\pi E_0 R^2 \left[2 \frac{x^2}{2} \Big|_0^1 \right] = 2\pi E_0 R^2$$

$$2\pi E_0 R^2 = q_{encl} / \epsilon_0 = \sigma\pi R^2 / \epsilon_0 \quad \therefore \quad E_0 = \sigma / 2\epsilon_0$$

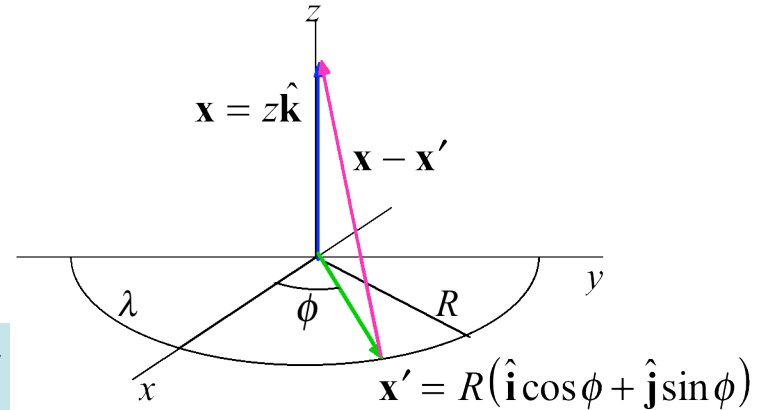


Exam 1 (cont'd)

4a)
$$q = \int \rho(x') d^3 x' = \int \lambda d\ell = \int_{-\pi/2}^{\pi/2} R d\phi = \pi R \lambda$$

b)
$$\mathbf{E}(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\mathbf{x} - \mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|^3} \rho(\mathbf{x}') d^3 x'$$

c)
$$|\mathbf{x} - \mathbf{x}'|^3 = |z\hat{\mathbf{k}} - R(\hat{\mathbf{i}}\cos\phi + \hat{\mathbf{j}}\sin\phi)|^3 = (z^2 + R^2)^{3/2}$$



$$\begin{aligned} \mathbf{E}(\mathbf{x}) &= \frac{1}{4\pi\epsilon_0} \int \frac{z\hat{\mathbf{k}} - R(\hat{\mathbf{i}}\cos\phi + \hat{\mathbf{j}}\sin\phi)}{(z^2 + R^2)^{3/2}} \lambda R d\phi \\ &= \frac{1}{4\pi\epsilon_0} \left(\frac{\lambda R}{(z^2 + R^2)^{3/2}} \pi z \hat{\mathbf{k}} + \frac{\lambda R^2}{(z^2 + R^2)^{3/2}} \int_{-\pi/2}^{\pi/2} (\hat{\mathbf{i}}\cos\phi + \hat{\mathbf{j}}\sin\phi) d\phi \right) \end{aligned}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{\lambda R}{(z^2 + R^2)^{3/2}} (\pi z \hat{\mathbf{k}} - 2R \hat{\mathbf{i}})$$

d)
$$E_z(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \frac{\pi \lambda R z}{(z^2 + R^2)^{3/2}} = \frac{q}{4\pi\epsilon_0} \frac{z}{z^3 (1 + (R/z)^2)^{3/2}}$$

$$R/z \rightarrow 0 \quad (\text{i.e., } z \gg R)$$

$$E_z(\mathbf{x}) \approx \frac{q}{4\pi\epsilon_0 z^2}$$

Exam 1 (cont'd)

5) Identity for $\nabla \cdot (\mathbf{A} \times \mathbf{B})$ using $\nabla \times \mathbf{A}$ and $\nabla \times \mathbf{B}$

$$\mathbf{C} = \mathbf{A} \times \mathbf{B} \text{ where } C_i = \epsilon_{ijk} A_j B_k$$

$$\nabla \cdot (\mathbf{C}) = \frac{\partial C_i}{\partial x_i} = \frac{\partial}{\partial x_i} (\epsilon_{ijk} A_j B_k) = B_k \epsilon_{ijk} \frac{\partial A_j}{\partial x_i} + A_j \epsilon_{ijk} \frac{\partial B_k}{\partial x_i}$$

Make replacements $\epsilon_{kij} = \epsilon_{ijk}$ and $-\epsilon_{jik} = \epsilon_{ijk}$

$$\begin{aligned} \nabla \cdot (\mathbf{C}) &= B_k \left(\epsilon_{kij} \right) \frac{\partial A_j}{\partial x_i} + A_j \left(-\epsilon_{jik} \right) \frac{\partial B_k}{\partial x_i} \\ &= B_k (\nabla \times \mathbf{A})_k - A_j (\nabla \times \mathbf{B})_j \end{aligned}$$

$$= \mathbf{B} \cdot (\nabla \times \mathbf{A}) - \mathbf{A} \cdot (\nabla \times \mathbf{B})$$

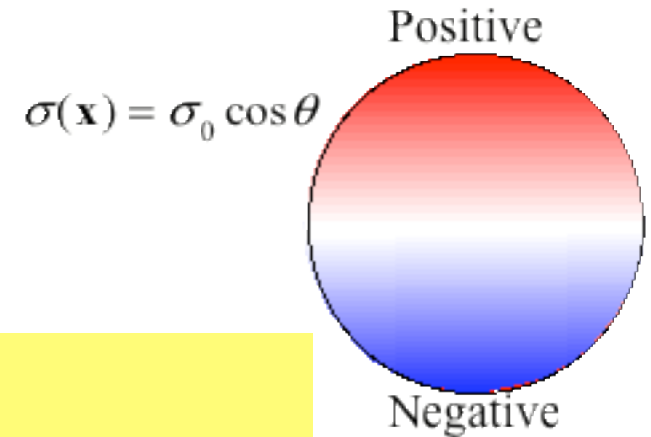
Exam 1 (cont'd)

6a) Dipole moment

$$\mathbf{p} = \int \mathbf{r} \rho(\mathbf{x}) d^3x$$

$$\hat{\mathbf{r}} = \hat{\mathbf{k}} \cos \theta + \sin \theta (\hat{\mathbf{i}} \cos \phi + \hat{\mathbf{j}} \sin \phi)$$

$$dA = R^2 \sin \theta d\theta d\phi$$



$$\begin{aligned} \mathbf{p} &= 2\pi R^3 \int \hat{\mathbf{r}} (\sigma_0 \cos \theta) \sin \theta d\theta \\ &= 2\pi R^3 \sigma_0 \left[\int_0^\pi \hat{\mathbf{k}} \cos^2 \theta \sin \theta d\theta \right] + R^3 \sigma_0 \int_0^\pi \cos \theta \sin^2 \theta d\theta \int_0^{2\pi} d\phi (\hat{\mathbf{i}} \cos \phi + \hat{\mathbf{j}} \sin \phi) \\ &= 2\pi R^3 \sigma_0 \hat{\mathbf{k}} \int_1^{-1} -x^2 dx = 2\pi R^3 \sigma_0 \hat{\mathbf{k}} \left[-\frac{x^3}{3} \right]_1^{-1} \end{aligned}$$

$$= \frac{4\pi R^3 \sigma_0}{3} \hat{\mathbf{k}}$$

b) Top is positive, bottom is negative, the same as a dipole

c) From far away, potential is that of a point dipole : $V(r, \theta) = \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{4\pi\epsilon_0 r^2}$

Exam 1 (cont'd)

7a) Prove $\int_{-\infty}^{\infty} \delta(ax + b) f(x) dx = \frac{f(-b/a)}{|a|}$ (d^3x should be dx , answer OK)

$x' = ax + b; \quad dx = dx'/|a|; \quad 2^{\text{nd}} \text{ integral is } f \text{ at } x' = 0$

$$\int_{-\infty}^{\infty} \delta(ax + b) f(x) dx = \frac{1}{a} \int_{-\infty}^{\infty} \delta(x') f(x'/a - b/a) dx' = \frac{f(-b/a)}{|a|}$$

b) Prove $\int_{-\infty}^{\infty} \delta^3(\mathbf{Ax} + \mathbf{b}) f(\mathbf{x}) d^3x = f(-\mathbf{A}^{-1}\mathbf{b}) / |\det \mathbf{A}|$ **Jacobian** $|\det \mathbf{A}|$

$$\mathbf{x} = \mathbf{A}^{-1}\mathbf{x}' - \mathbf{A}^{-1}\mathbf{b}$$

$$d^3x = d^3x' / |\text{Det } \mathbf{A}|$$

$$\begin{aligned} \int_{-\infty}^{\infty} \delta(\mathbf{Ax} + \mathbf{b}) f(\mathbf{x}) d^3x &= \int_{-\infty}^{\infty} \delta(\mathbf{x}') f(\mathbf{A}^{-1}\mathbf{x}' - \mathbf{A}^{-1}\mathbf{b}) d^3x' / |\text{Det } \mathbf{A}| \\ &= f(-\mathbf{A}^{-1}\mathbf{b}) / |\text{Det } \mathbf{A}| \end{aligned}$$

Conductors in static equilibrium

Why these statements are true ?

- Electric field $\mathbf{E}$ inside a conductor is zero.
- Potential V inside a conductor is a constant.
- Inside a conductor the charge density ρ is zero.
- Charge Q on a conductor resides only on the surface.
- A net charge Q_{net} here is always paired with charge $-Q_{net}$ elsewhere.
- On a conductor either V or Q_{net} , but not both, can be specified.
- At a conductor's surface, field component $E_t = 0$, component $E_n = \sigma/\epsilon_0$,
but in force calculations, $E_n = \sigma/2\epsilon_0$ due to only "distant" charges.
- In an empty cavity in a conductor, $\mathbf{E} = 0$ and surface $\sigma_{cavity} = 0$.
- The potential V of a conductor can be set by a battery V_B .
- The earth is an "infinite" source of charge at a constant V .
- V or Q_{net} of conductors (and $Q_{external}$) determine V everywhere.

Boundary value problems: 1-parallel plates

Parallel plates (find potential, field & surface charge)

The ground is a constant potential, so let $V_1 = 0$.

Battery draws electrons from plate 2 and puts them on plate 1. Process "continues" until $V_2 = V_B$.

Potential between the plates

Between the plates the potential depends only on x .

Laplace's equation: $\nabla^2 V(x) = 0$

Boundary conditions: $V(0) = 0; \quad V(d) = V_B$

General solution (ODE)

$$\nabla^2 V = \frac{d}{dx} \left(\frac{dV}{dx} \right) = 0; \quad \frac{dV}{dx} = c_1$$

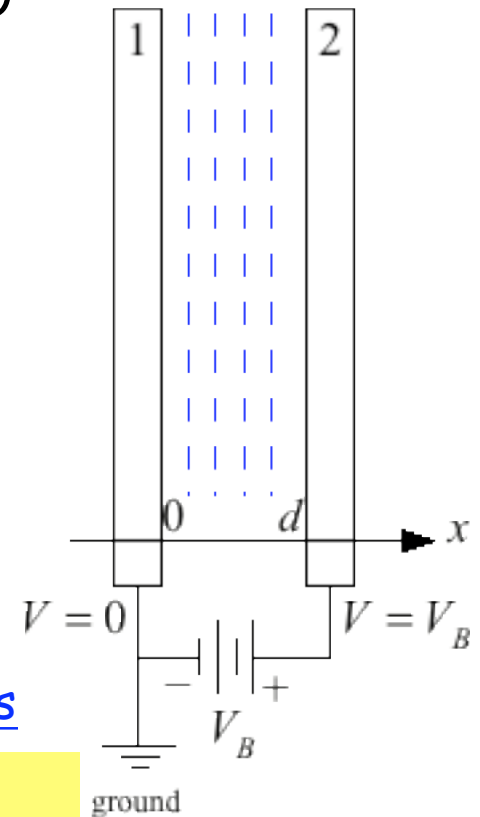
$$V(x) = \underline{c_1 x + c_2}$$

Apply boundary conditions

$$V(0) = \underline{c_2 = 0}$$

$$V(d) = c_1 d = V_B; \quad \underline{c_1 = \frac{V_B}{d}}$$

equipotentials



Potential

$$V(x) = \underline{V_B \frac{x}{d}}$$

Parallel plates (cont'd)

Electric field

Electric field

$$\mathbf{E} = -\nabla V = -\nabla \left(V_B \frac{x}{d} \right) = -\frac{V_B}{d} \hat{\mathbf{i}}$$

$E_n \rightarrow$ Surface charges

Surface $1r$ charge density

$$E_{1n} = \frac{\sigma_{1r}}{\epsilon_0} = -\frac{V_B}{d}$$

$$\sigma_{1r} = -\frac{\epsilon_0 V_B}{d}$$

Surface 2ℓ charge density

$$E_{2n} = \frac{\sigma_{2\ell}}{\epsilon_0} = \frac{V_B}{d}$$

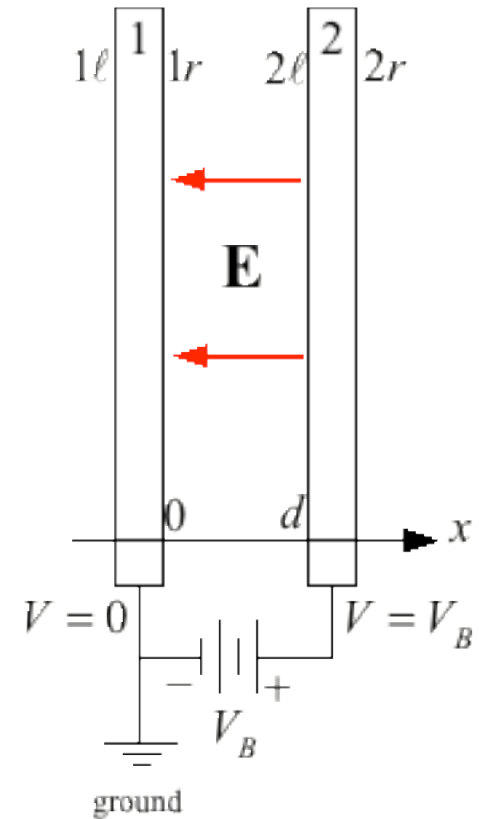
$$\sigma_{2\ell} = +\frac{\epsilon_0 V_B}{d}$$

Surface charge

$$Q_{2\ell} = -Q_{1r} = \sigma A = \frac{\epsilon_0 V_B A}{d}$$

Force on each plate

$$F_2 = Q_2 (E_{2n})_{non-q_2} = \frac{1}{2} \frac{\sigma_2^2}{\epsilon_0} A$$



Parallel plates (cont'd)

Capacitance

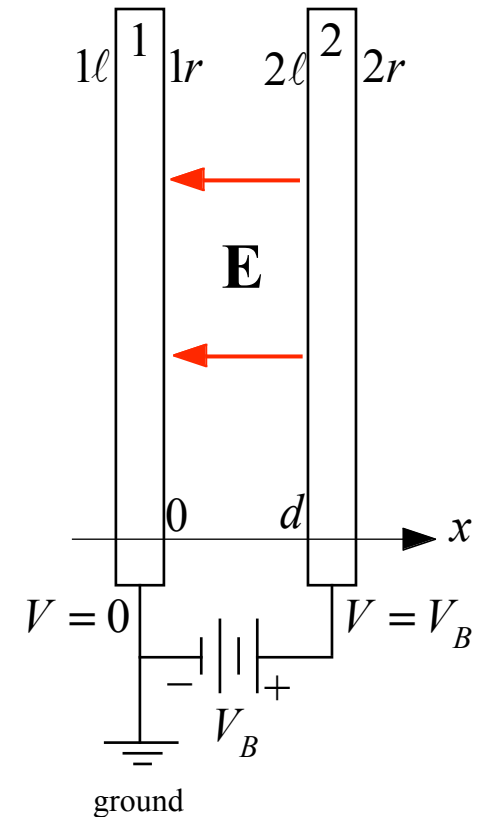
Two ways to know the electric field

$$E = \frac{\sigma}{\epsilon_0}; \quad E = \frac{V_B}{d}$$

Capacitance is a geometrical factor relating the charge on a conductor and its potential.

$$Q = \sigma A = \frac{\epsilon_0 A}{d} V_B = C V_B$$

$$C = \frac{\epsilon_0 A}{d}$$



Energy storage

$$U = \frac{\epsilon_0}{2} \int E^2 d^3 x'$$

$$= \frac{\epsilon_0}{2} \frac{V_B^2}{d^2} A d = \frac{1}{2} \frac{\epsilon_0 A}{d} V_B^2$$

$$U = \frac{1}{2} C V_B^2$$