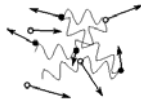


# Black-Body Radiation



$$u_{\lambda} d\lambda = \frac{8\pi hc/\lambda^5}{e^{hc/\lambda kT} - 1} d\lambda$$

[CO Sect. 3.4]

Black body

$$B_{\lambda}(T) = \frac{2hc^2/\lambda^5}{e^{hc/\lambda kT} - 1}$$

$$L = A\sigma T^4$$

$$\lambda_{\max} T = 0.002897755 \text{ m K}$$

for large  $\lambda$ : (Rayleigh-Jeans tail)

$$B_{\lambda}(T) \simeq \frac{2ckT}{\lambda^4}$$

also... dilute black body

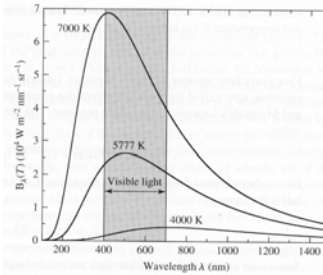
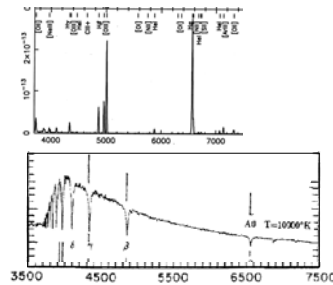
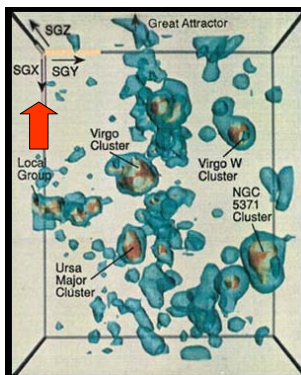


FIGURE 3.8 Blackbody spectrum [Planck function  $B_{\lambda}(T)$ ].



# Isotropy of the Cosmic Microwave Background



Local Group

Black body

$$B_{\lambda}(T) = \frac{2hc^2/\lambda^5}{e^{hc/\lambda kT} - 1}$$

for large  $\lambda$

$$B_{\lambda}(T) \simeq \frac{2ckT}{\lambda^4}$$

Dipole Anisotropy

$\sim 1$  part in 300.

$$T_{\text{moving}} = \frac{T_{\text{rest}} \sqrt{1 - v^2/c^2}}{1 - (v/c) \cos \theta}$$

for  $v \ll c$

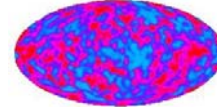
$$T_{\text{moving}} \simeq T_{\text{rest}} \left( 1 + \frac{v}{c} \cos \theta \right) \quad (29.62)$$



Blue = 0°K  
Red = 4°K



Blue = 2.724°K  
Red = 2.732°K  
Dipole Anisotropy  
→ motion of Sun through Universe.



After removing dipole  
Red - blue = 0.0002°K

- Dipole Anisotropy  $\sim 1$  part in 300
  - 600 km/sec motion of Local Group in grav. field of larger scale mass concentrations.

## Isotropy of the Cosmic Microwave Background

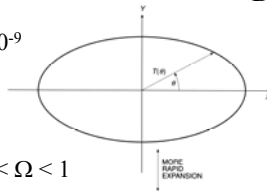
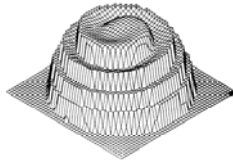
Is the universe really isotropic?

- No quadrupole anisotropy in CMB

$$\rightarrow R_x/R_y = 1 \pm \varepsilon, \quad \varepsilon < 10^{-9}$$

- Vorticity:

$$(\omega/H_0) < 10^{-8} \text{ for } 0.05 < \Omega < 1$$

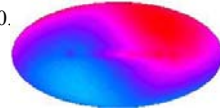


Dipole Anisotropy

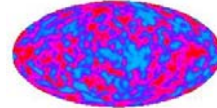
~ 1 part in 300.



Blue = 0°K  
Red = 4°K



Blue = 2.724°K  
Red = 2.732°K  
Dipole Anisotropy  
→ motion of Sun through Universe.



After removing dipole  
Red - blue = 0.0002°K

Fig. 8.6 The pattern of the CBR temperature  $T(\theta, \phi)$  in a Bianchi type  $VII_0$  model. Only one hemisphere, corresponding to  $\pi/2 \leq \theta < \pi$ , is shown. Note the 'spiral' pattern in  $T(\theta, \phi)$  (from Barrow *et al.* (1985), with permission).

## When did decoupling occur?

- Saha equation: Collisional ionization rate = Recombination rate

$$\frac{N_{H^+}}{N_{H^0}} = \frac{4}{N_e} \left( \frac{2\pi m_e k T}{h^2} \right)^{3/2} e^{-E_I/kT} \quad (8.8)$$

- Solve for  $\frac{N_{H^+}}{N_{H^0}} = 1$ , using  $T = \frac{T_0}{R}$ ,  $N_e \sim \frac{\rho_0}{m_H} \frac{1}{R^3}$

$$\rightarrow R \sim 7.2 \times 10^{-4}$$

$$T \sim 3800 \text{ K}$$

- Taking composition and radiative transfer into account:

$$T_{dec} = 2970 \text{ K}$$

$$R_{dec} = 9 \times 10^{-4}$$

$$z_{dec} = 1089$$

$$t_{dec} = 118,000 \text{ yrs.}$$

Decoupling also called "recombination"

## The Radiation Era

Energy density (integrated over wavelength):

$$u = aT^4$$

$$\rho_{\text{RAD}} = \frac{u}{c^2} = \frac{aT^4}{c^2} = \frac{\rho_{\text{RAD},0}}{R^4}$$

$$\rho_{\text{matter}} = \frac{\rho_{m,0}}{R^3}$$

$$\Rightarrow \frac{\rho_{\text{RAD}}}{\rho_{\text{matter}}} \propto \frac{1}{R}$$

$\rho_{\text{matter}} = \rho_{\text{rad}}$  at

$$R = \frac{aT_0^4}{\rho_0 c^2} = \frac{8\pi G a T_0^4}{3H_0^2 c^2 \Omega_0} \sim 2.5 \times 10^5 \Omega_0^{-1} h^{-2}$$

$$z = \frac{1}{R} - 1 \sim 4 \times 10^9 \Omega_0 h^2$$

when  $T = \frac{T_0}{R} \sim 1.1 \times 10^5 \Omega_0 h^2 \text{ K}$

$$\text{Age} = \frac{2}{3} \frac{1}{H_0(1+z)^{3/2} \Omega_0^{1/2}} \sim 3200 \text{ years}$$

for  $\Omega_0 = 1, h = 0.71$

$$\begin{aligned} T_0 &= RT \\ T &= \frac{T_0}{R} \end{aligned}$$

## During the Radiation Era:

$$\left[ \left( \frac{1}{R} \frac{dR}{dt} \right)^2 - \frac{8}{3} \pi G \rho_{\text{RAD}} \right] R^2 = -kc^2$$

$\uparrow k \sim 0$

$$\left( \frac{1}{R} \frac{dR}{dt} \right)^2 = \frac{8}{3} \pi G \frac{\rho_{\text{RAD},0}}{R^4}$$

$$\left( \frac{dR}{dt} \right)^2 \propto \frac{1}{R^2}$$

$$RdR \propto dt$$

$$R \propto t^{1/2}$$

instead of  $R \propto t^{1/2}$   
in matter era.

also,  $T_0 = RT \Rightarrow T \propto t^{-1/2}$

## Terminology...

$$\left( \left( \frac{1}{R} \frac{dR}{dt} \right)^2 - \frac{8}{3} \pi G \rho \right) R^2 = -kc^2$$

$$\left( \left( \frac{1}{R} \frac{dR}{dt} \right)^2 - \frac{8}{3} \pi G (\rho_{\text{matter}} + \rho_{\text{radiation}}) \right) R^2 = -kc^2$$

$$\rho_{\text{radiation}} = \frac{u_{\text{radiation}}}{c^2}$$

$$\Omega_{\text{radiation}} = \frac{\rho_{\text{radiation}}}{\rho_{\text{critical}}}$$

## PRIMORDIAL NUCLEOSYNTHESIS

Radiation era:  $R(t) \propto t^{1/2}$ ;  $RT = \text{constant} \Rightarrow T(t) \propto t^{-1/2}$   
 $= 2.726 \text{ }^\circ\text{K}$

$t = 10^{-4} \text{ s}$ ,  $T = 10^{12} \text{ K}$        $p + e^- \rightleftharpoons n + \nu_e$ , etc.      electrons from  $\gamma \rightleftharpoons e^- + e^+$

Statistical equilibrium       $\frac{n_n}{n_p} = e^{-\frac{\Delta E}{kT}} = e^{-\frac{1.5 \times 10^6 \text{ K}}{T}}$

$t \sim 1 \text{ s}$ ,  $T \sim 1 \times 10^9 \text{ K}$

neutrons "freeze out"  
 - redshifting of neutrinos  
 - lower  $\gamma$  energies.

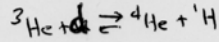
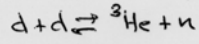
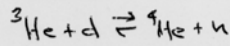
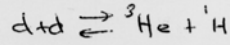
subsequent  $\beta$  decays  $\Rightarrow$

$$\frac{n_n}{n_p} = 0.223$$

$$\frac{n_n}{n_p} \sim 0.164$$

$t \sim 3 \text{ min}, T \sim 10^9 \text{ K}$

${}^4\text{He}$  production:



$\sim$  all neutrons  $\rightarrow {}^4\text{He}$

$\Rightarrow$  mass fraction of  ${}^4\text{He}$

$$= \frac{2n_n}{n_n + n_p} \sim 0.281$$

ROUGHLY INDEPENDENT OF EXACT DENSITY

Observed  $\frac{{}^4\text{He}}{\text{Total}} = 0.228 \pm 0.005$  (Pagel et al. 1992)

or something close to that.

