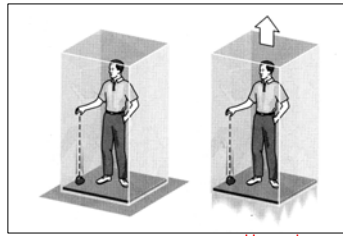


General Relativity (sort of)

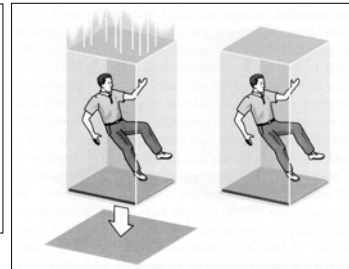
Equivalence Principle:

- Can't tell difference between gravity & acceleration
- ...or between freefall & no gravity.
- So *any* experiment should give same answer in either case.



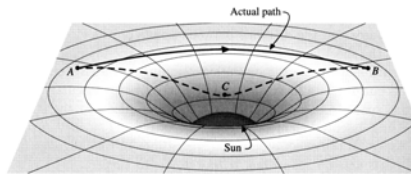
Gravity

Upwards acceleration, no gravity.

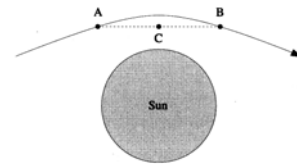


Falling due to gravity

No gravity



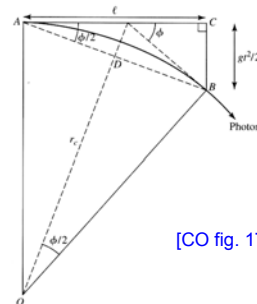
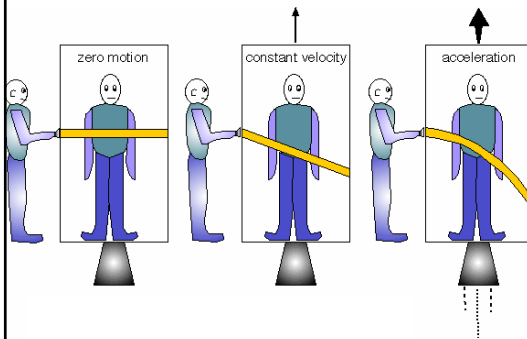
Gravity = Curved space



Objects follow shortest distance through curved space(-time).

The equivalence principle (plus a vigorous waving of one's hands) shows...

- Curved path of light in gravitational field



[CO fig. 17.10]

$$\overline{BC}/\overline{AC} = \overline{BD}/\overline{OD}$$

$$\left(\frac{1}{2}gt^2\right)/\ell = \left[\frac{\ell}{2\cos(\phi/2)}\right]/\overline{OD}$$

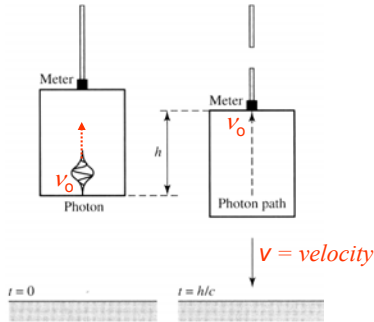
$$r_c = \frac{c^2}{g} = 9.17 \times 10^{15} \text{ m}$$

$$\phi = \frac{\ell}{r_c} = 1.09 \times 10^{-15} \text{ rad}$$



The equivalence principle (plus a vigorous waving of one's hands) shows...

Gravitational Redshift



Equiv. Principle → photon frequency unchanged in free-falling lab.

Why doesn't the frequency meter see a blueshift?

$$\frac{\Delta v}{v_0} = \frac{v}{c} = \frac{gh}{c^2}$$

There must be a counteracting Gravitational Redshift:

This redshift is seen by the meter which is *not* free-falling. $\Rightarrow \frac{\Delta v}{v_0} = -\frac{v}{c} = -\frac{gh}{c^2}$

Integrate the effect out to infinite distance:

$$\int_{v_0}^{v_\infty} \frac{dv}{v} \approx - \int_{r_0}^{\infty} \frac{GM}{r^2 c^2} dr$$

$$\frac{v_\infty}{v_0} \approx 1 - \frac{GM}{r_0 c^2}$$

The exact result:

$$\frac{v_\infty}{v_0} = \left(1 - \frac{2GM}{r_0 c^2}\right)^{1/2}$$



Gravitational time dilation

$$\frac{\Delta t_0}{\Delta t_\infty} = \frac{v_\infty}{v_0} = \left(1 - \frac{2GM}{r_0 c^2}\right)^{1/2}$$

Special Relativity The Lorentz Transformation

CO, pg. 90

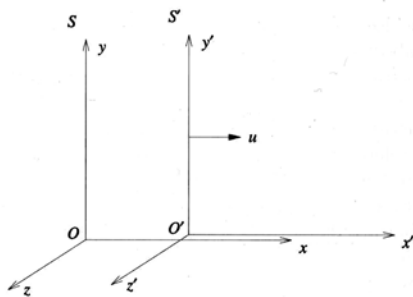


Figure 4.2 Inertial reference frames S and S' .

$$x' = \frac{x - ut}{\sqrt{1 - u^2/c^2}} \quad (4.16)$$

$$y' = y \quad (4.17)$$

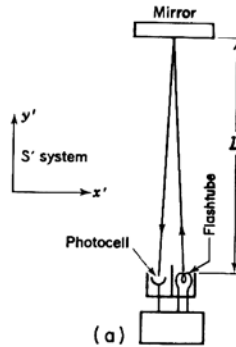
$$z' = z \quad (4.18)$$

$$t' = \frac{t - ux/c^2}{\sqrt{1 - u^2/c^2}} \quad (4.19)$$

Time Dilation in Special Relativity

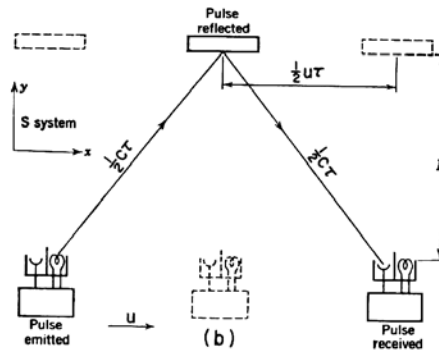
"Light Clock"

velocity = distance / time
speed of light = D / time



As seen
by
moving
observer

Direction of
motion
→



As seen
by
stationary
observer

From the Feynman Lectures

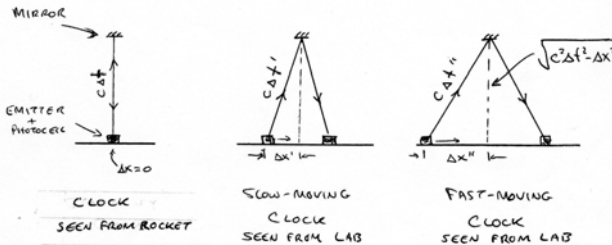
INTERVALS & METRICS

SPECIAL RELATIVITY

IF SPEED OF LIGHT SAME IN ALL REFERENCE FRAMES:
& LAWS OF PHYSICS SAME IN ALL INERTIAL FRAMES:

"CLOCK" THOUGHT-EXPERIMENT SHOWS:

1. TIME RUNS AT DIFFERENT SPEEDS AS SEEN FROM DIFFERENT REFERENCE FRAMES.
2. $c^2\Delta t^2 - \Delta x^2$ IS INVARIANT QUANTITY



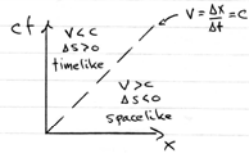
$$\text{INTERVAL}^2 = \Delta S^2 = c^2\Delta t^2 - \Delta x^2 = \text{INVARIANT.}$$

↕
LORENTZ TRANSFORM

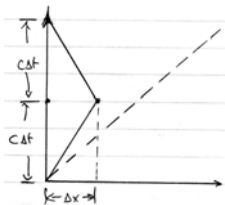
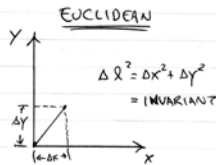
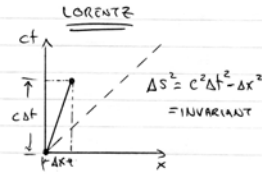
$$\begin{cases} x' = \frac{x - ut}{\sqrt{1 - u^2/c^2}} \\ y' = y \\ z' = z \\ t' = \frac{t - ux/c^2}{\sqrt{1 - u^2/c^2}} \end{cases}$$

SPACE-TIME (LORENTZ GEOMETRY)

$$\Delta S^2 = c^2 \Delta t^2 - \Delta x^2 = \text{INVARIANT}$$



$\Delta S = 0$ for light
 $0 = c^2 \Delta t^2 - \Delta x^2$
 $V = \frac{\Delta x}{\Delta t} = c$



CURVED PATHS ARE SHORTER INTERVALS IN LORENTZ GEOMETRY

STRAIGHT: $\Delta S = 2c\Delta t$
 CURVED: $\Delta S = 2\sqrt{c^2 \Delta t^2 - \Delta x^2}$

$$\Delta s^2 = c^2 \Delta t^2 - \Delta x^2$$

Generalize to:
(Minkowski space)

$$ds^2 = c^2 dt^2 - dx^2 - dy^2 - dz^2$$

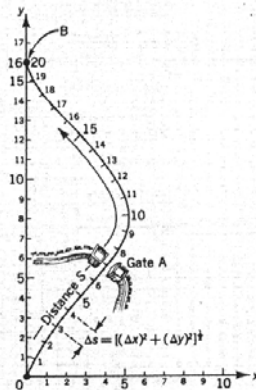
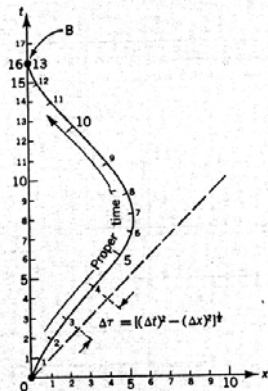
World line: a particle's path through space-time

Proper time = $\int ds$ along time-like world-line

Proper time interval along world line is smaller than distance covered along time axis!

Proper time = elapsed time on clock moving with particle

Proper distance = $\int \sqrt{-(ds^2)}$ along space-like world-line.



Euclidean geometry, for comparison.

Figures from Taylor & Wheeler, "Spacetime Physics"

Metrics

From special relativity:

$$ds^2 = c^2 dt^2 - dx^2 - dy^2 - dz^2$$

For General Relativity:

$$ds^2 = \sum_{ij} g_{ij} dx^i dx^j$$

$$g = \begin{bmatrix} c^2 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

Field equation

$$G_{\mu\nu} = -\frac{8\pi G}{c^4} T_{\mu\nu}$$

Einstein Tensor:

Curvature of space. g_{ij} & derivatives.

Mass + energy Stress-energy tensor

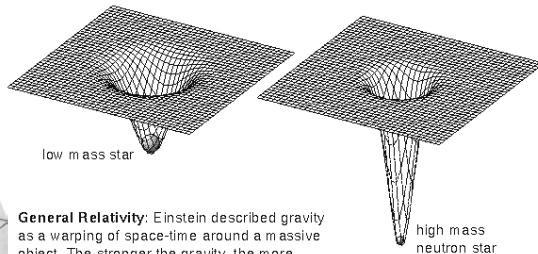
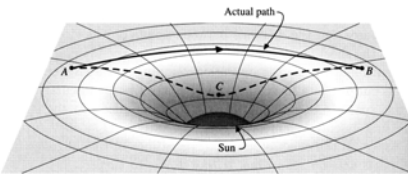
Schwarzschild metric (black holes, etc)

$$(ds)^2 = \left(c dt \sqrt{1 - \frac{2GM}{rc^2}} \right)^2 - \left(\frac{dr}{\sqrt{1 - \frac{2GM}{rc^2}}} \right)^2 - (r d\theta)^2 - (r \sin\theta d\phi)^2$$

Robertson-Walker metric (the Universe)

$$(ds)^2 = (c dt)^2 - R^2(t) \left[\left(\frac{dr}{1-kr^2} \right)^2 + (d\theta)^2 + (r \sin\theta d\phi)^2 \right]$$

What curves into where?



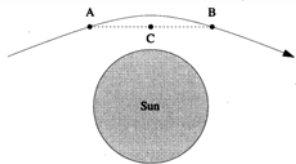
General Relativity: Einstein described gravity as a warping of space-time around a massive object. The stronger the gravity, the more space-time is warped.

Schwarzschild metric (black holes, etc)

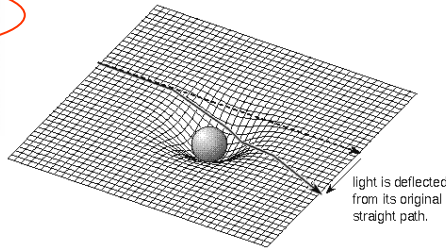
$$(ds)^2 = \left(c dt \sqrt{1 - \frac{2GM}{rc^2}} \right)^2 - \left(\frac{dr}{\sqrt{1 - \frac{2GM}{rc^2}}} \right)^2 - (r d\theta)^2 - (r \sin\theta d\phi)^2$$

Robertson-Walker metric (the Universe)

$$(ds)^2 = (c dt)^2 - R^2(t) \left[\left(\frac{dr}{1-kr^2} \right)^2 + (d\theta)^2 + (r \sin\theta d\phi)^2 \right]$$



Objects follow shortest distance through curved space(-time).



light is deflected from its original straight path.

General Relativity: Light travels along the curved space taking the shortest path between two points. Therefore, light is deflected toward a massive object! The stronger the local gravity is, the greater the light path is bent.

Paths of Objects through Curved Space-time

- *Geodesic* = straightest possible worldline
 - = maximum value of $\int ds$ in most examples.
 - But actually = extremum (max or min).
- Special Relativity:
 - Objects with no forces acting on them take straightest path through space-time.
 - Conservation of Energy-Momentum.
- General Relativity:
 - Objects follow Geodesics through curved space-time.
 - This follows directly from the equivalence principle.
- Light follows null geodesic
 - $ds = 0$ at every point along path
 - $\rightarrow \int ds = 0$
 - $\rightarrow d(\text{length})/dt = c$

Does the Schwarzschild metric actually tell us anything?

Now for some incredibly vigorous hand waving....

- The orbit of a satellite

Schwarzschild metric: $(ds)^2 = \left(c dt \sqrt{1 - 2GM/rc^2} \right)^2 - \left(\frac{dr}{\sqrt{1 - 2GM/rc^2}} \right)^2 - (r d\theta)^2 - (r \sin \theta d\phi)^2$

Assume:

- Circular orbit $\rightarrow dr = 0$
- In plane where $d\phi = 0$
- Moving at constant angular velocity $\omega \rightarrow d\theta = \omega dt$

$$(ds)^2 = \left[\left(c \sqrt{1 - 2GM/rc^2} \right)^2 - r^2 \omega^2 \right] dt^2 = \left(c^2 - \frac{2GM}{r} - r^2 \omega^2 \right) dt^2$$

$$\Delta s = \int_0^{2\pi/\omega} \sqrt{c^2 - \frac{2GM}{r} - r^2 \omega^2} dt.$$

Find r that makes Δs be an extremum:

$$\frac{d}{dr}(\Delta s) = \frac{d}{dr} \left(\int_0^{2\pi/\omega} \sqrt{c^2 - \frac{2GM}{r} - r^2 \omega^2} dt \right) = 0$$

$$\frac{d}{dr} \sqrt{c^2 - \frac{2GM}{r} - r^2 \omega^2} = 0$$

$$\frac{2GM}{r^2} - 2r\omega^2 = 0$$

$$v = r\omega = \sqrt{\frac{GM}{r}}$$

A familiar
Newtonian result!

