
PHY481: Electromagnetism

HW5

5.3

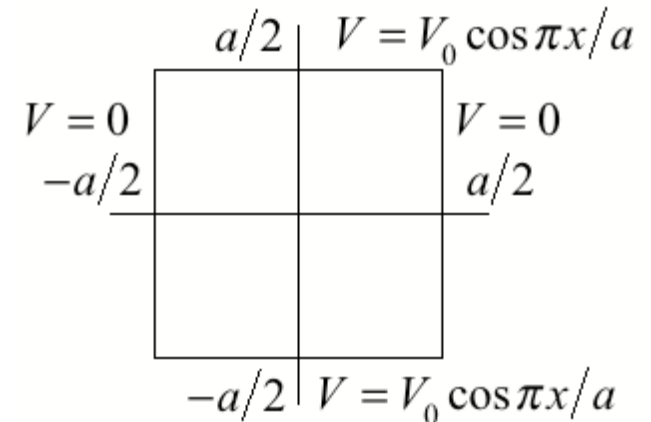
Boundary condition $V = V_0 \cos\left(\frac{\pi x}{a}\right)$ at $y = \pm a/2$

General solution (for even boundary conditions)

$$V(x, y) = \sum_{n=1}^{\infty} C_n \cos\left[(2n+1)\frac{\pi x}{a}\right] \cosh\left[(2n+1)\frac{\pi y}{a}\right]$$

$$\begin{aligned} V(x, \pm a/2) &= V_0 \cos\left(\frac{\pi x}{a}\right) \\ &= C_n \cos\left[(2n+1)\frac{\pi x}{a}\right] \cosh\left[(2n+1)\frac{\pi}{2}\right] \end{aligned}$$

Square pipe cross section



Only non-zero term

$n = 0$, C_0 to be determined

Apply boundary condition to determine coefficient

$$\begin{aligned} V(x, \pm a/2) &= C_0 \cos\left[\frac{\pi x}{a}\right] \cosh\left[\frac{\pi}{2}\right] = V_0 \cos\left[\frac{\pi x}{a}\right] \\ C_0 &= \frac{V_0}{\cosh\left[\frac{\pi}{2}\right]} \end{aligned}$$

Potential inside pipe

$$V(x, y) = \frac{V_0}{\cosh(\pi/2)} \cos(\pi x/a) \cosh(\pi y/a)$$

Phase factors

Prove: $A \sinh(x + \phi) = B \sinh(x) + C \cosh(x)$

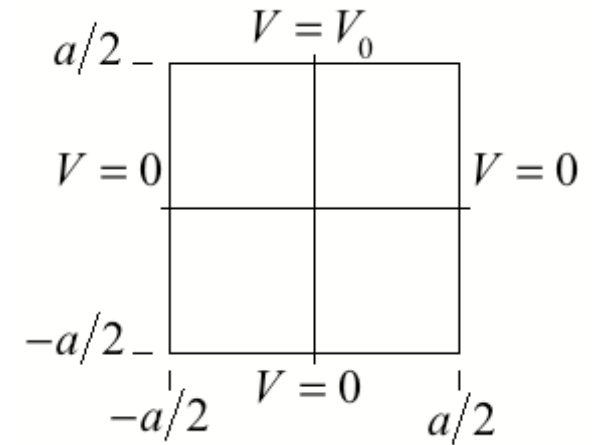
$$\begin{aligned} A \left[\frac{e^{x+\phi} - e^{-(x+\phi)}}{2} \right] &= B \left[\frac{e^x - e^{-x}}{2} \right] + C \left[\frac{e^x + e^{-x}}{2} \right] \\ A \left[\frac{e^\phi e^x - e^{-\phi} e^{-x}}{2} \right] &= (C + B) \left[\frac{e^x}{2} \right] + (C - B) \left[\frac{e^{-x}}{2} \right] \\ C + B &= A e^\phi \\ C - B &= A e^{-\phi} \\ C &= A \frac{e^\phi - e^{-\phi}}{2}; \quad B = A \frac{e^\phi + e^{-\phi}}{2} \\ \sinh(x + \phi) &= \cosh(\phi) \sinh(x) + \sinh(\phi) \cosh(x) \end{aligned}$$

5.7

$$X(x) = B_n \cos[(2n+1)\pi x/a] \quad \text{boundary conditions: } X(\pm a/2) = 0$$

$$Y(y) = C_n \sinh[(2n+1)\pi y/a] + D_n \cosh[(2n+1)\pi y/a] \\ = C'_n \sinh[(2n+1)\pi y/a + \phi_n] \quad \text{instead use phase: } \phi_n$$

$$V(x, y) = \sum_{n=0}^{\infty} g_n \sinh[(2n+1)\pi y/a + \phi_n] \cos[(2n+1)\pi x/a]$$

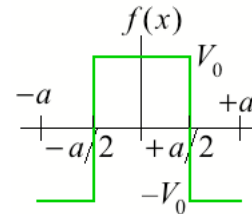


$$V(x, -a/2) = 0 = \sum_{n=0}^{\infty} g_n \sinh[-(2n+1)\pi/2 + \phi_n] \cos[(2n+1)\pi x/a] \quad \text{bc: } Y(-a/2) = 0 \text{ sets phase}$$

$$\phi_n = (2n+1)\pi/2$$

$$V(x, +a/2) = V_0 = \sum_{n=0}^{\infty} g_n \sinh[(2n+1)\pi] \cos[(2n+1)\pi x/a] \quad \text{bc: } Y(+a/2) = V_0 \text{ sets coef. } g_n$$

$$\frac{1}{a} \int_{-a}^{+a} f(x) \cos[(2n+1)\pi x/a] = g_n \sinh[(2n+1)\pi]$$



$$g_n = \frac{4V_0 (-1)^n}{(2n+1)\pi \sinh[(2n+1)\pi]}$$

Be careful to break
integral into pieces !!

$$V(x, y) = \frac{4V_0}{\pi} \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n+1)} \frac{\sinh[(2n+1)\pi(y/a + 1/2)]}{\sinh[(2n+1)\pi]} \cos[(2n+1)\pi x/a]$$

5.11

See lectures 18-20 for derivations of potential for a conducting sphere in an external field

$$V(r, \theta) = + \frac{E_0 a^3 \cos \theta}{r^2} - E_0 r \cos \theta$$

Dipole potential

$$V(r, \theta) = \frac{p \cos \theta}{4\pi\epsilon_0 r^2}$$

Compare $1/r^2$ potentials

$$p = 4\pi\epsilon_0 E_0 a^3$$

Polarizability

$$p = \alpha E_0$$

$$\alpha = 4\pi\epsilon_0 a^3$$

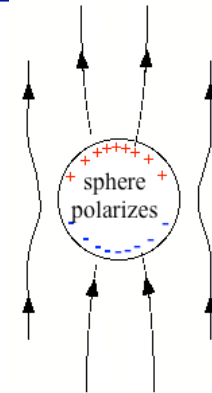
$$\sigma(\theta) = \epsilon_0 E_r(a, \theta) = -\epsilon_0 \left. \frac{\partial V}{\partial r} \right|_{r=a}$$

$$= 2\epsilon_0 \frac{E_0 a^3 \cos \theta}{a^3} + \epsilon_0 E_0 \cos \theta = 3\epsilon_0 E_0 \cos \theta$$

Charge density

$$\sigma(\theta) = 3\epsilon_0 E_0 \cos \theta$$

Dipole moment analogous to $p = qd$



$$\mathbf{r} = a \cos \theta \hat{\mathbf{k}} + a \sin \theta \hat{\boldsymbol{\theta}}$$

$$\mathbf{p} = \int \sigma(\theta) \mathbf{r} dS$$

$$= 2\pi a^3 \int (3\epsilon_0 E_0 \cos \theta) \cos \theta \sin \theta d\theta \hat{\mathbf{k}}$$

$$p_z = 6\pi a^3 \epsilon_0 E_0 \int_0^\pi \cos^2 \theta \sin \theta d\theta$$

$$= 6\pi a^3 \epsilon_0 E_0 \int_1^{-1} -u^2 du$$

Dipole moment

$$p_z = 4\pi a^3 \epsilon_0 E_0$$

5.13

Linear quadrupole exact potential

$$V(r, \theta) = \frac{q}{4\pi\epsilon_0} \left\{ \frac{1}{\sqrt{r^2 + a^2 - 2ar \cos \theta}} + \frac{-2}{r} + \frac{1}{\sqrt{r^2 + a^2 + 2ar \cos \theta}} \right\}$$

Approximation

Expansions

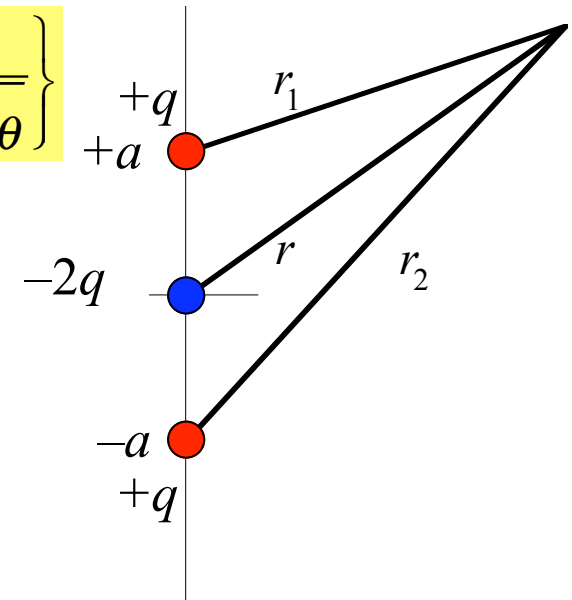
$$\frac{1}{\sqrt{1 + \epsilon}} = 1 - \frac{1}{2}\epsilon + \frac{3}{8}\epsilon^2 + \dots$$

$$\frac{1}{\sqrt{r^2 + a^2 \pm 2ar \cos \theta}} = \frac{1}{r} \left\{ 1 \pm \frac{a \cos \theta}{r} + \frac{3a^2 \cos^2 \theta - a^2}{2r^2} + \dots \right\}$$

Monopole and dipole terms cancel

$$V(r, \theta) = \frac{qa^2}{4\pi\epsilon_0} \left\{ \frac{3\cos^2 \theta - 1}{2r^3} + \dots \right\} = \frac{2qa^2}{4\pi\epsilon_0} \frac{P_2(\cos \theta)}{r^3} + \dots$$

Linear quadrupole



5.15 (see lecture 19)

Half cylinders, radius R . $V = +V_0$ on upper half, and $V = -V_0$ on lower half
Find potential inside and out. (odd function of $\phi \rightarrow$ no cosine terms)

$$V(r, \phi) = A \ln r + B + \sum_{n=1}^{\infty} (A_n r^n + B_n r^{-n}) (C_n \cos n\phi + D_n \sin n\phi)$$

B.C. at $r = 0$ and $r = \infty$

$$V_{Int}(r, \phi) = \sum_{n=1}^{\infty} A_n r^n \sin n\phi$$

$$V_{Ext}(r, \phi) = \sum_{n=1}^{\infty} B_n r^{-n} \sin n\phi$$

B.C. at $r = R$

$$V_{Int}(R, \phi) = V_{Ext}(R, \phi)$$

$$A_n R^n = B_n R^{-n} = c_n \rightarrow A_n = \frac{c_n}{R^n}; \quad B_n = c_n R^n$$

Fourier Integral at $r = R$

Be careful to break Fourier integral into pieces !!

Use Orthogonality

$$\int_0^{2\pi} \cos n\phi \cos m\phi = \pi \delta_{nm}$$

$$\int_0^{2\pi} V(R, \phi) \sin m\phi d\phi = \sum_{n=1}^{\infty} c_n \int_0^{2\pi} \sin n\phi \sin m\phi d\phi$$

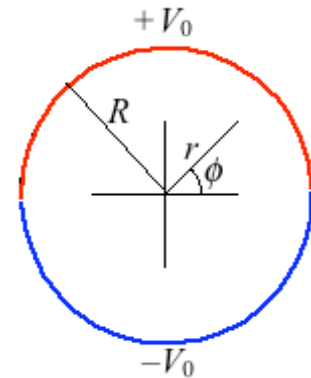
$$-V_0 \cos m\phi \Big|_0^{\pi} + V_0 \cos m\phi \Big|_{\pi}^{2\pi} = c_n \pi \delta_{mn}$$

$$c_m = \frac{4V_0}{\pi m} \quad m = 1, 3, 5, \dots$$

Complete solution

$$V_{Int}(r, \phi) = \frac{4V_0}{\pi} \sum_{m \text{ odd}} \frac{1}{m} \frac{r^m}{R^m} \sin m\phi$$

$$V_{Ext}(r, \phi) = \frac{4V_0}{\pi} \sum_{m \text{ odd}} \frac{1}{m} \frac{R^m}{r^m} \sin m\phi$$



5.16

General solution in cylindrical coordinates

$$V(r, \phi) = A \ln r + B + \sum_{n=1}^{\infty} (A_n r^n + B_n r^{-n}) (C_n \cos n\phi + D_n \sin n\phi)$$

$$E_r(R_{ext}) - E_r(R_{int}) = \frac{\sigma}{\epsilon_0} \quad \text{Boundary condition crossing the boundary}$$

$$V(r_{int}, \phi) = A_1 r \cos \phi \quad E_n(r_{int}, \phi) = -A_1 \cos \phi$$

$$V(r_{ext}, \phi) = \frac{B_1}{r} \cos \phi \quad E_n(r_{ext}, \phi) = + \frac{B_1 \cos \phi}{r_{ext}^2} \quad \text{Matching terms}$$

Apply boundary conditions

$$\frac{B_1 \cos \phi}{r_{ext}^2} + A_1 \cos \phi = \frac{\sigma_0 \cos \phi}{\epsilon_0}$$

$$\frac{B_1}{R^2} + A_1 = \frac{\sigma_0}{\epsilon_0}$$

$$V_{int}(R, \phi) = V_{ext}(R, \phi)$$

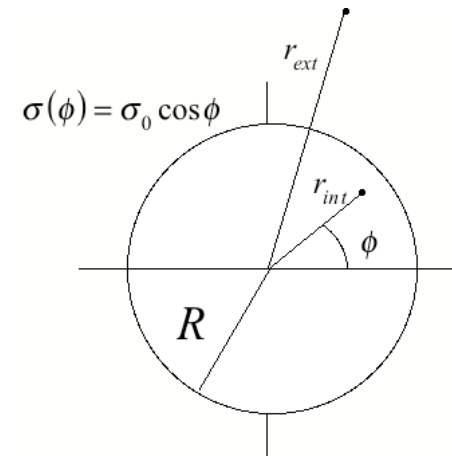
$$A_1 R \cos \phi = \frac{B_1}{R} \cos \phi$$

$$A_1 R = \frac{B_1}{R}$$

$$A_1 = \frac{\sigma_0}{2\epsilon_0} \quad B_1 = \frac{\sigma_0 R^2}{2\epsilon_0}$$

Solution inside and out

$$V(r_{int}, \phi) = \frac{\sigma_0}{2\epsilon_0} r \cos \phi; \quad V(r_{ext}, \phi) = \frac{\sigma_0 R^2}{2\epsilon_0 r} \cos \phi$$

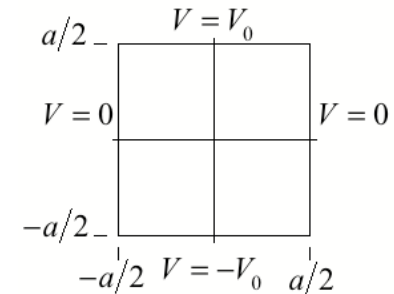


5.28 (see problem 5.7)

$$X(x) = B_n \cos[(2n+1)\pi x/a] \quad \text{boundary conditions: } X(\pm a/2) = 0$$

$$Y(y) = C_n \sinh[(2n+1)\pi y/a] \quad \text{odd B. C. : } Y(\pm a/2) = \pm V_0$$

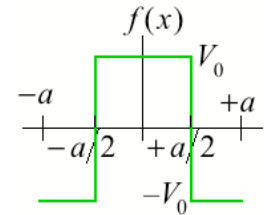
$$V(x, y) = \sum_{n=0}^{\infty} g_n \sinh[(2n+1)\pi y/a] \cos[(2n+1)\pi x/a] \quad g_n = B_n C_n$$



$$V(x, +a/2) = V_0 = \sum_{n=0}^{\infty} g_n \sinh[(2n+1)\pi/2] \cos[(2n+1)\pi x/a] \quad \text{B.C.: } Y(+a/2) = V_0 \text{ sets coef. } g_n$$

$$\frac{1}{a} \int_{-a}^{+a} f(x) \cos[(2n+1)\pi x/a] dx = g_n \sinh[(2n+1)\pi/2]$$

Be careful to break Fourier integral into pieces !!



$$\frac{V_0}{(2n+1)\pi} \left\{ \begin{array}{l} -\sin[(2n+1)\pi x/a]_{-a}^{-a/2} + \sin[(2n+1)\pi x/a]_{-a/2}^{a/2} \\ -\sin[(2n+1)\pi x/a]_{a/2}^a \end{array} \right\} = g_n \sinh[(2n+1)\pi/2]$$

$$g_n = \frac{4V_0 (-1)^n}{(2n+1)\pi \sinh[(2n+1)\pi/2]}$$

$$V(x, y) = \frac{4V_0}{\pi} \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n+1)} \frac{\sinh[(2n+1)\pi y/a]}{\sinh[(2n+1)\pi/2]} \cos[(2n+1)\pi x/a]$$

5.32

Dipole potential

$$V = \frac{p \cos \theta}{4\pi\epsilon_0 r^2}$$

General solution in spherical coordinates

$$V(r, \theta) = \sum_{\ell=0}^{\infty} \left(A_{\ell} r^{\ell} + \frac{B_{\ell}}{r^{\ell+1}} \right) P_{\ell}(\cos \theta)$$

Dipole in grounded conducting sphere

$$V(r, \theta) = \left(Ar + \frac{B}{r^2} \right) P_1(\cos \theta) \quad \text{Only one term} \quad \ell = 1$$

Apply boundary conditions

$r \rightarrow 0$

$$V(r, \theta) \rightarrow \frac{B \cos \theta}{r^2} = \frac{p \cos \theta}{4\pi\epsilon_0 r^2}$$

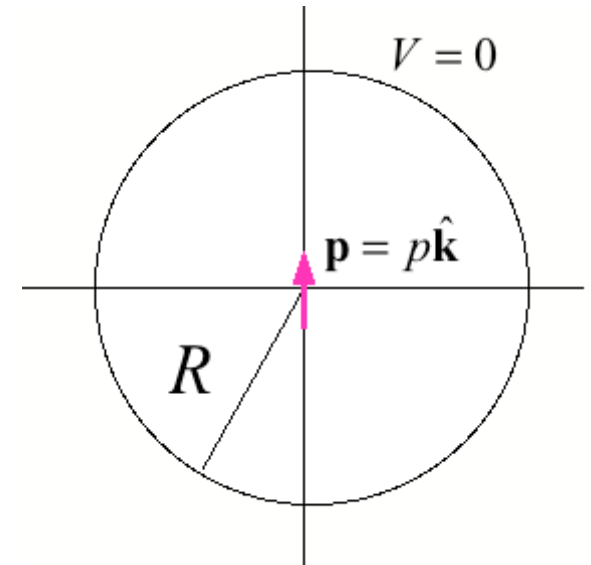
$$B = \frac{p}{4\pi\epsilon_0}$$

$r = R$

$$V(R, \theta) = \left(AR + \frac{B}{R^2} \right) = 0$$

$$A = \frac{-B}{R^3} = \frac{-p}{4\pi\epsilon_0 R^3}$$

Full Solution:
$$V(r, \theta) = \frac{p}{4\pi\epsilon_0} \left(\frac{-r}{R^3} + \frac{1}{r^2} \right) \cos \theta$$



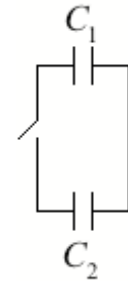
2

Caps C_1 and C_2 carry charges Q_1 and Q_2 .

Close the switch - find new charges Q'_1 and Q'_2

Charge conserved: $Q'_1 + Q'_2 = Q_1 + Q_2$

Potentials become equal: $V' = Q'_1/C_1 = Q'_2/C_2; \quad Q'_2 = (C_2/C_1)Q'_1$



$$Q'_1 + (C_2/C_1)Q'_1 = Q_1 + Q_2$$

$$Q'_1 = \frac{Q_1 + Q_2}{[1 + (C_2/C_1)]}; \quad Q'_2 = \frac{C_2}{C_1} \frac{Q_1 + Q_2}{(1 + (C_2/C_1))} = \frac{Q_1 + Q_2}{(1 + (C_1/C_2))}$$

Compare stored energy before and after switch is closed

$$Q_1 = Q'_1 + \Delta Q; \quad Q_2 = Q'_2 - \Delta Q$$

$$U'_1 + U'_2$$

Potentials become equal: $V'_1 - V'_2 = 0$

$$U_1 + U_2 = \frac{1}{2} \left(\frac{Q_1^2}{C_1} + \frac{Q_2^2}{C_2} \right)$$

$$= \frac{1}{2} \left(\frac{Q_1'^2}{C_1} + \frac{Q_2'^2}{C_2} \right)$$

$$+ \frac{1}{2} \left(\frac{(\Delta Q)^2}{C_1} + \frac{(\Delta Q)^2}{C_2} \right)$$

$$+ \Delta Q \left(\frac{Q'_1}{C_1} - \frac{Q'_2}{C_2} \right)$$

$$U_1 + U_2 = U'_1 + U'_2 + \frac{(\Delta Q)^2}{2C}; \quad C = \frac{C_1 C_2}{C_1 + C_2}$$

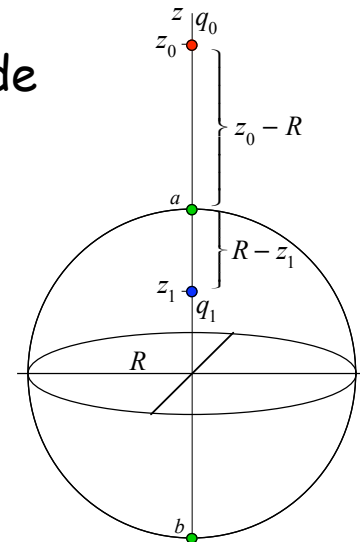
3 a,b

Charge q_0 above a grounded conduction sphere

Determine a) image charge and position, b) potential outside

Potential at points (a) and (b) must be zero.

$$\begin{aligned}\frac{q_1}{R - z_1} &= \frac{-q_0}{z_0 - R}; & q_1 &= \frac{-q_0(R - z_1)}{z_0 - R} \\ \frac{q_1}{R + z_1} &= \frac{-q_0}{z_0 + R}; & q_1 &= \frac{-q_0(R + z_1)}{z_0 + R} \\ \frac{(R - z_1)}{z_0 - R} &= \frac{(R + z_1)}{z_0 + R}; & 2R^2 &= 2z_1z_0\end{aligned}$$



$$a) \quad z_1 = \frac{R^2}{z_0}; \quad q_1 = -q_0 \frac{R}{z_0}$$

$$b) \quad V(r, \theta) = \frac{1}{4\pi\epsilon_0} \left[\frac{q_0}{\sqrt{r^2 + z_0^2 - 2rz_0 \cos \theta}} + \frac{q_1}{\sqrt{r^2 + z_1^2 - 2rz_1 \cos \theta}} \right]$$

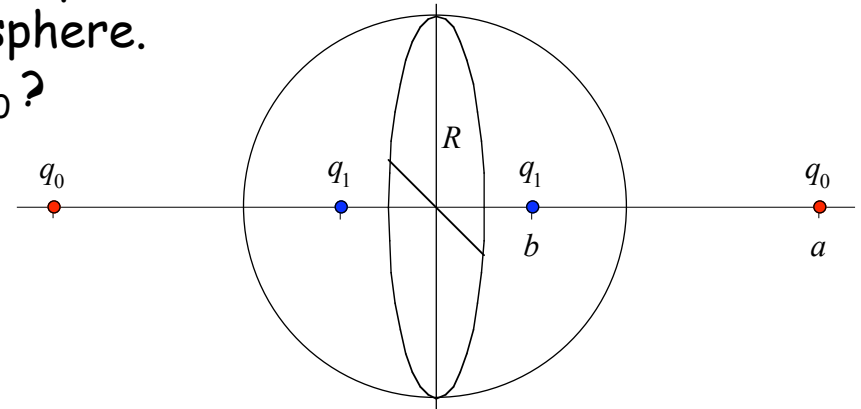
3 c

Two charges q_0 outside grounded conduction sphere

Image: Two charges q_1 inside make $V=0$ on sphere.

What radius sphere makes force $F = 0$ on q_0 ?

$$a = z_0, b = R^2/a; \quad q_1 = -q_0 R/a$$



$$\begin{aligned} F &= \frac{q_0^2}{4a^2} + \frac{q_1 q_0}{(a+b)^2} + \frac{q_1 q_0}{(a-b)^2} \\ &= \frac{q_0^2}{4a^2} - \frac{R}{a} \frac{q_0^2}{(a+R^2/a)^2} - \frac{R}{a} \frac{q_0^2}{(a-R^2/a)^2} \end{aligned}$$

$$\begin{aligned} 0 &= \frac{1}{4a^2} - \frac{R}{a^3} \left(\frac{1}{(1+R^2/a^2)^2} + \frac{1}{(1-R^2/a^2)^2} \right) \\ &= \frac{1}{4a^2} - \frac{R}{a^3} \left[\left(1 - 2\frac{R^2}{a^2} + 3\frac{R^4}{a^4} + \dots \right) + \left(1 + 2\frac{R^2}{a^2} + 3\frac{R^4}{a^4} + \dots \right) \right] \end{aligned}$$

$$\frac{1}{4a^2} = \frac{2R}{a^3} \left[1 + 3\frac{R^4}{a^4} + \dots \right]; \quad \frac{1}{4a^2} \approx \frac{2R}{a^3} \quad \text{assuming } R^2/a^2 \ll 1$$

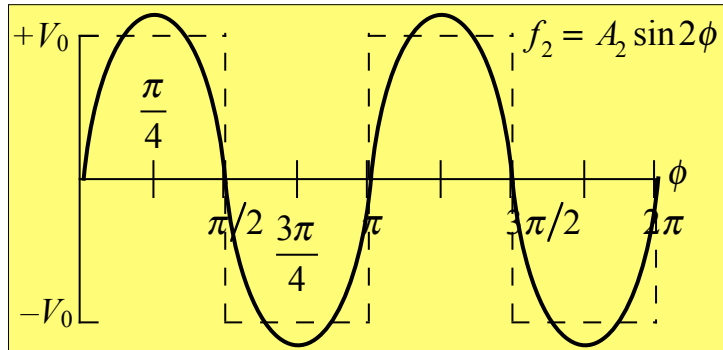
$$\frac{1}{(1 \pm x)^2} = 1 \mp 2x + 3x^2 + \dots$$

$$R \approx \frac{a}{8}$$

4

Use this function: $f_n(\phi) = A_n \sin n\phi$

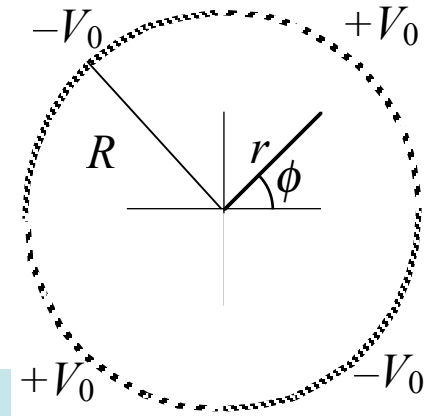
a) What is the best n to match this cylindrical potential?



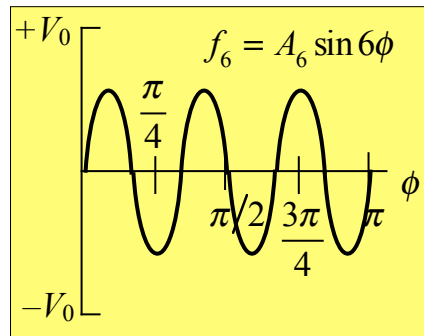
$$f_{n=2} = A_2 \sin 2\phi$$

zeros: $\pi/2, \pi, 3\pi/2, 2\pi, \dots$

critical points: $\pi/4, 3\pi/4, 5\pi/4, 7\pi/4, \dots$



b) Next integer m where f_m includes these zeros and critical points?



zeros: $0, \pi/6, \pi/3, \pi/2, 2\pi/3, 5\pi/6, \pi, \dots$

critical points: $\pi/12, \pi/4, 5\pi/12, 7\pi/12, 3\pi/4, 3\pi/4, 11\pi/12, \dots$

c) Find $V(r, \phi)$ that satisfies Laplace's equation and compare first two terms with the functions above.

4 c

$$V(r, \phi) = \sum_{n=1}^{\infty} A_n r^{-n} \sin n\phi; \quad \text{Boundary: } V(R, \phi) = \sum_{n=1}^{\infty} A_n R^{-n} \sin n\phi$$

$$V(R, \phi) = \pm V_0$$

$$\frac{1}{\pi} \int_0^{2\pi} V(R, \phi) \sin m\phi d\phi = \frac{1}{\pi} \sum_{n=1}^{\infty} A_n R^{-n} \int_0^{2\pi} \sin n\phi \sin m\phi d\phi = A_m R^{-m}$$

$$\int_0^{2\pi} \sin n\phi \sin m\phi d\phi = \pi \delta_{mn}$$

$$= \frac{2V_0}{m\pi} \left\{ -\cos m\phi \Big|_0^{\pi/2} + \cos m\phi \Big|_{\pi/2}^{\pi} \right\} = \frac{2V_0}{m\pi} \{ -2\cos m\pi/2 + 1 + \cos m\pi \}$$

$$\begin{aligned} \left\{ -2\cos m\frac{\pi}{2} + 1 + \cos m\pi \right\} &= 0 \quad (m=1) \\ &= 4 \quad (m=2) \\ &= 0 \quad (m=3, 4, 5) \\ &= 4 \quad (m=6) \end{aligned}$$

$$\frac{1}{\pi} \int_0^{2\pi} V(R, \phi) \sin m\phi d\phi = \frac{8V_0}{m\pi} = A_m R^{-m}$$

$$\rightarrow A_m = \frac{8V_0 R^m}{m\pi}; \quad m = 4n - 2, \quad n = 1, 2, 3, \dots$$

$$V(r, \phi) = \frac{8V_0}{\pi} \sum_{n=1}^{\infty} \frac{1}{4n-2} \frac{R^{4n-2}}{r^{4n-2}} \sin[(4n-2)\phi]$$

5

Solve Laplace's equation for a single line charge (radius a), charge density λ . You may need to use Gauss's Law to determine one of the unknown constants.

$$V(r, \phi) = A \ln r + B; \quad \mathbf{E} = -\nabla V = -\frac{A}{r} \hat{r}$$

Gauss's Law

Potential on the surface of line charge

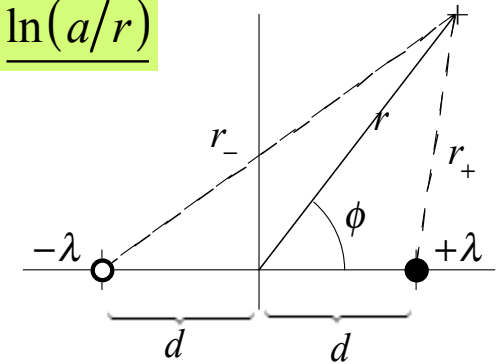
$$2\pi r L E_r = \frac{\lambda L}{\epsilon_0}; \quad E_r = \frac{\lambda}{2\pi\epsilon_0 r} \rightarrow A = -\frac{\lambda}{2\pi\epsilon_0}$$

$$V(a, \phi) = 0 = -\frac{\lambda}{2\pi\epsilon_0} \ln a + B; \quad B = \frac{\lambda}{2\pi\epsilon_0} \ln a$$

$$V(r, \phi) = (-\lambda/2\pi\epsilon_0) \ln r + (\lambda/2\pi\epsilon_0) \ln a = (\lambda/2\pi\epsilon_0) \ln(a/r)$$

Find potential of Line Dipole:

$$V_d(r, \phi) = \frac{\lambda}{2\pi\epsilon_0} \ln \frac{a}{r_+} + \frac{-\lambda}{2\pi\epsilon_0} \ln \frac{a}{r_-} = \frac{\lambda}{2\pi\epsilon_0} (\ln r_- - \ln r_+)$$



$$r_{\pm} = r \left(1 \mp 2d \cos \phi / r + d^2 / r^2 \right)^{1/2}$$

$$\ln r_+ = \ln r + \frac{1}{2} \ln \left(1 - 2d \cos \phi / r + d^2 / r^2 \right)$$

$$\ln r_- = \ln r + \frac{1}{2} \ln \left(1 + 2d \cos \phi / r + d^2 / r^2 \right)$$

$$\begin{aligned} V_d(r, \phi) &= \frac{\lambda}{2\pi\epsilon_0} (\ln r_- - \ln r_+) \\ &\approx \frac{2\lambda d \cos \phi}{2\pi\epsilon_0 r} = \frac{p_{\text{line}} \cos \phi}{2\pi\epsilon_0 r}, \text{ where } p_{\text{line}} = 2\lambda d \end{aligned}$$