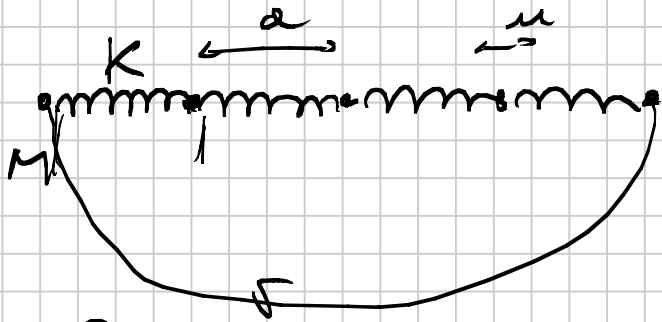


LECTURE # 18

Note Title

11/7/2007

Theory of harmonic crystal



Born von Karman
boundary conditions

$q =$ normal modes
of oscillation

modes = N atoms

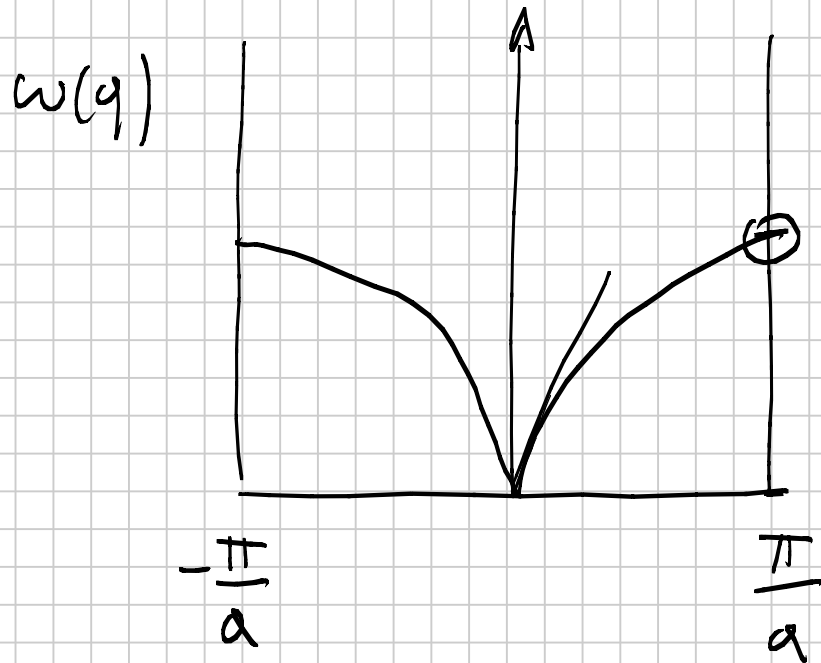
$\omega(q)$

N atoms

K spring constant

$$u_m(q) = e^{iqna}$$

$$q = \frac{2\pi}{N} n \quad 0 < n < N-1$$

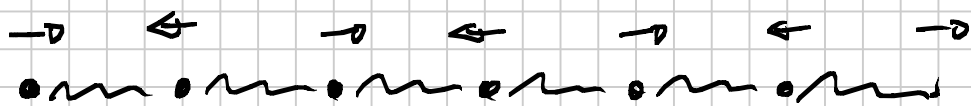
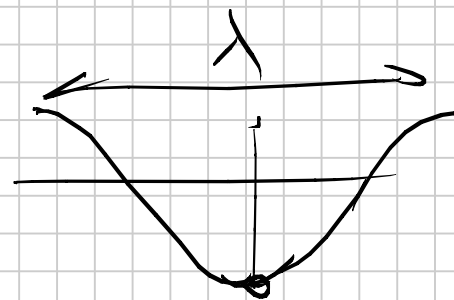


$$\omega(q) = \sqrt{\frac{4K}{M}} \left| \sin \frac{qa}{2} \right|$$

Small q

$$\omega(q) \sim \sqrt{\frac{K}{M}} a q \sim v_s q$$

$$q = \frac{\pi}{a} \Rightarrow \lambda = \frac{2\pi}{q} = \underline{\underline{2a}}$$



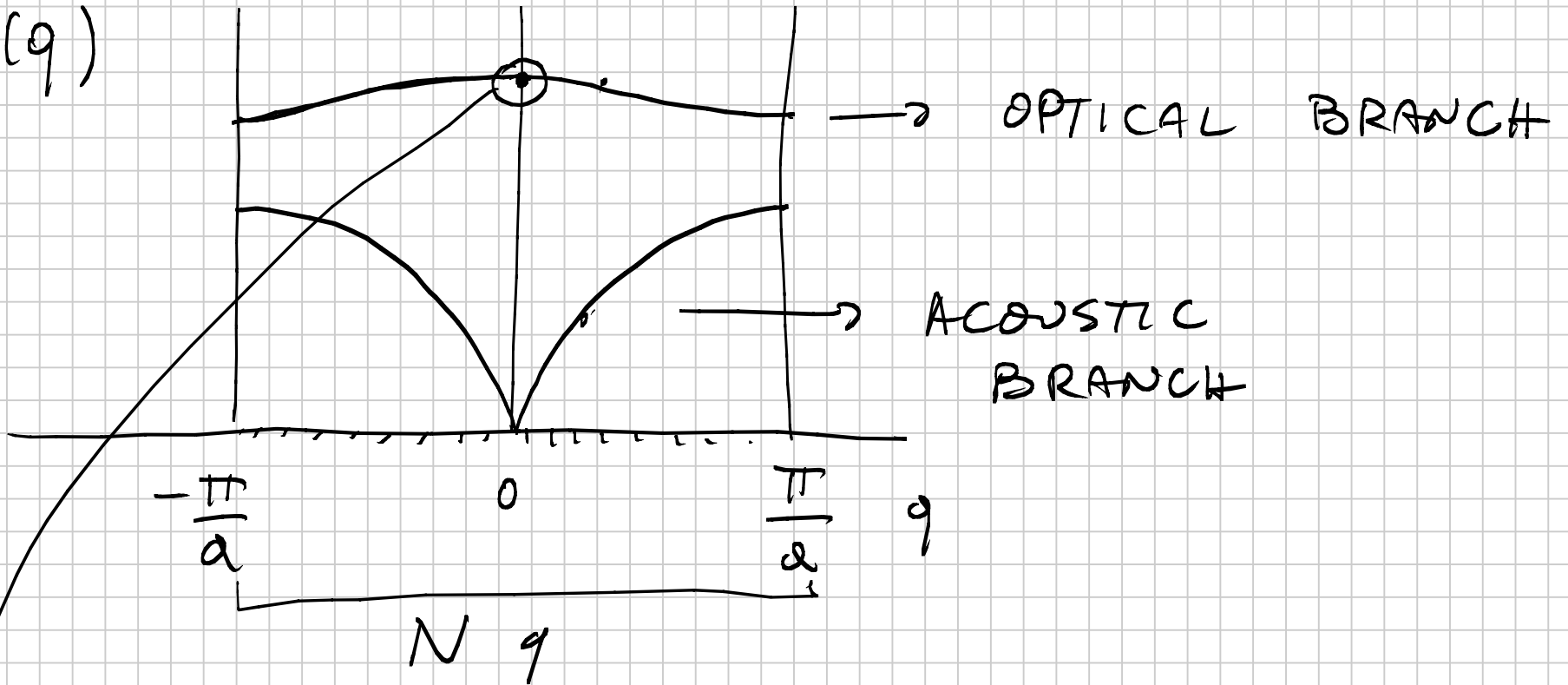
Crystal with basis



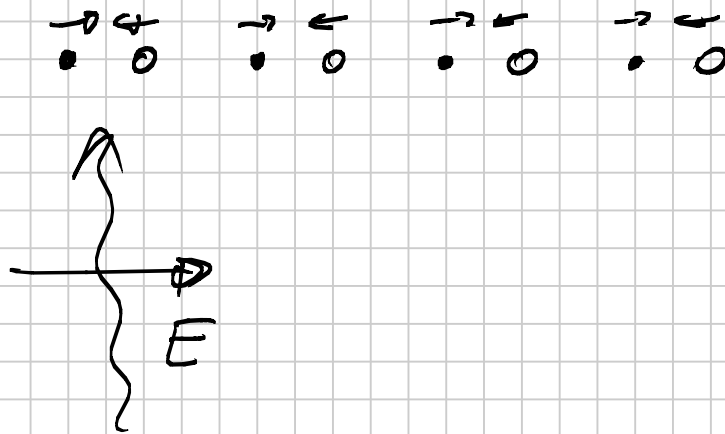
N unit cells
2 atoms on each cell

$2N$ DEGREES OF FREEDOM

$\omega(q)$



$q=0$
OPTICAL
BRANCH



1D Crystals

$1D \rightarrow 3P$

No basis

FCC

Pb

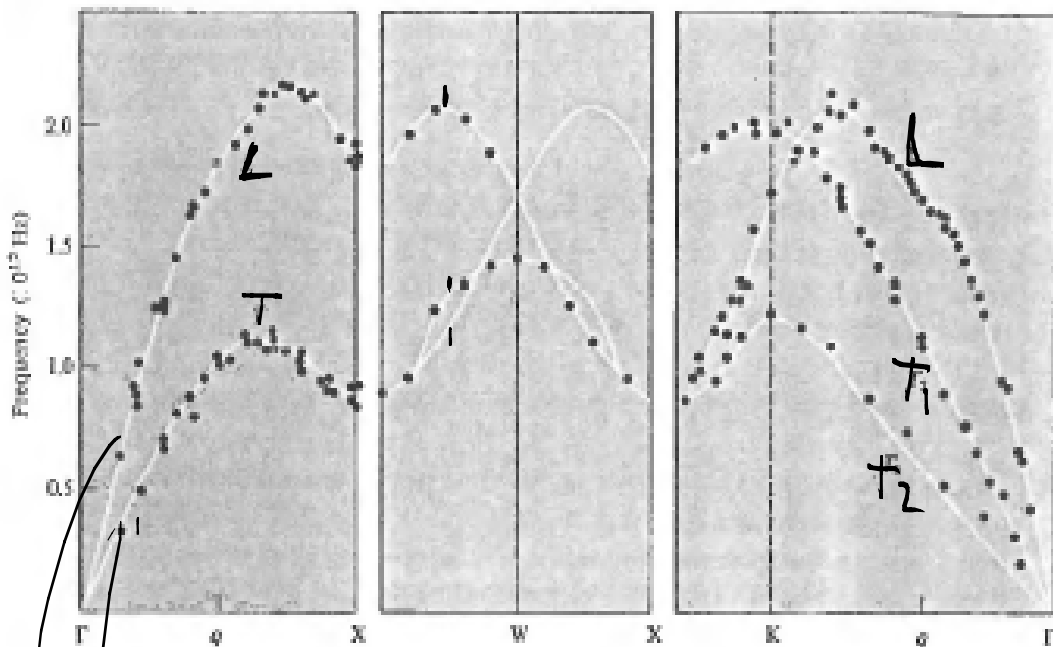
N TOT # UNIT CELLS = TOT # OF Pb IONS

• 3 $3N$ NORMAL MODES

N q WAVES $\Rightarrow \forall q \rightarrow 3$ MODES

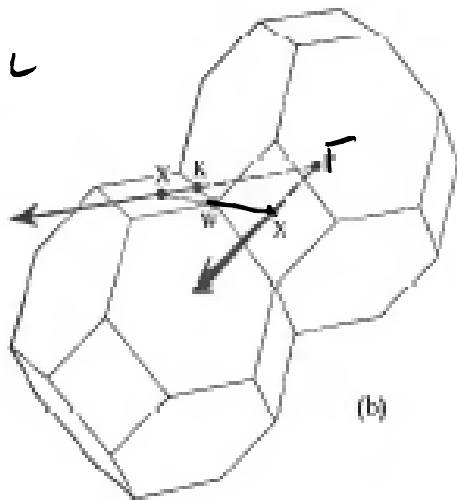
$\rightarrow 1$ LONGITUDINAL MODE $\vec{u} \parallel \vec{q}$

$\rightarrow 2$ TRANSVERSAL MODES $\vec{u} \perp \vec{q}$



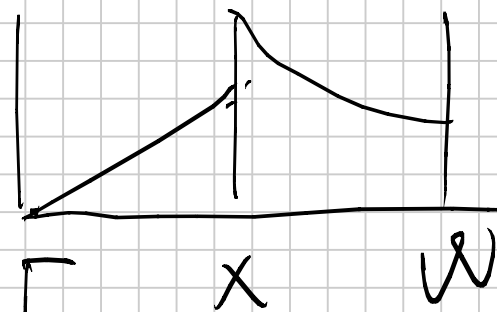
(a)

LONGITUDINAL BRANCH
 TRANSVERSAL BRANCH



(b)

3D



3D FCC LATTICE

BRILLOUIN ZONE

TRUNCATED OCTAHEDRON

$$\Gamma = (0, 0, 0)$$

$$X = \left(\frac{2\pi}{a}, 0, 0\right)$$

Figure 22.13

(a) Typical dispersion curves for the normal-mode frequencies in a monatomic Bravais lattice. The curves are for lead (face-centered cubic) and are plotted in a repeated-zone scheme along the edges of the shaded triangle shown in (b). Note that the two transverse branches are degenerate in the $[100]$ direction. (After Brockhouse et al., *Phys. Rev.* 128, 1099 (1962).)

$$W = \left(\frac{2\pi}{a}, \frac{\pi}{a}, 0 \right)$$

Calculate the
total E of crystal
3D crystal
 N atoms

$$\langle E \rangle = \langle T \rangle + \langle V \rangle$$

$$\langle T \rangle = \frac{k_B T}{2} \times 3N \quad \begin{array}{l} \nearrow \\ \# \text{ degrees} \\ \text{of freedom} \end{array} = \frac{3}{2} k_B T N$$

$\langle V \rangle$ VIRIAL THEOREM

$$V = \frac{1}{2} k x^2 \quad V(x) = x^\lambda$$

$$\text{THEOREM} \Rightarrow 2\langle T \rangle = \lambda \langle V \rangle$$

HARMONIC POTENTIAL

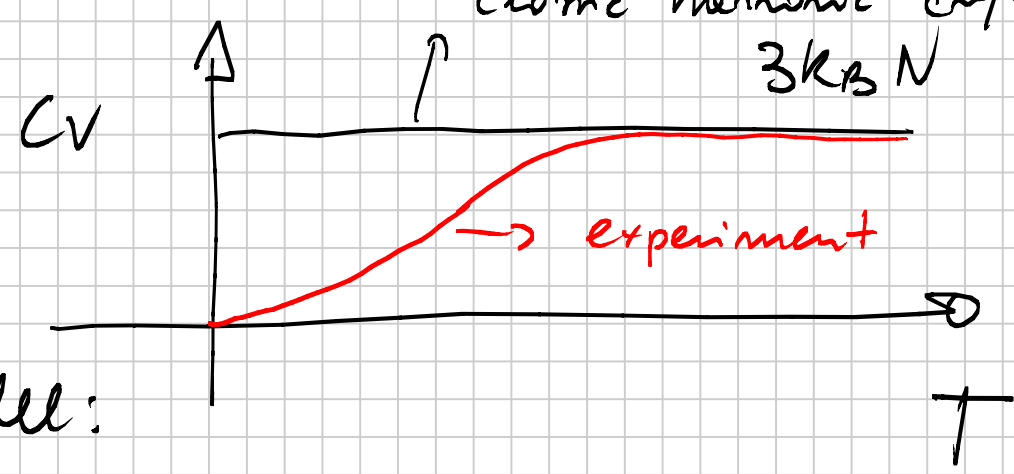
$$\langle T \rangle = \langle V \rangle$$

$$\langle E \rangle = N \frac{3}{2} k_B T + N \frac{3}{2} k_B T = 3 k_B N T$$

SPECIFIC HEAT

$$C_V = \frac{d\langle E \rangle}{dT} = 3 k_B N \rightarrow \text{CONSTANT}$$

classical harmonic crystal
 $3 k_B N$



We need new model:

Quantum Harmonic Crystal

Each $q \rightarrow 1$ Harmonic Oscillator

$$H = \frac{p^2}{2M} + \frac{1}{2} M \omega^2 x^2 \longrightarrow \frac{\hat{p}^2}{2m} + \frac{1}{2} m \omega^2 \hat{x}^2$$

$$\omega = \sqrt{\frac{k}{M}}$$

$$[\hat{x}, \hat{p}] = i\hbar$$

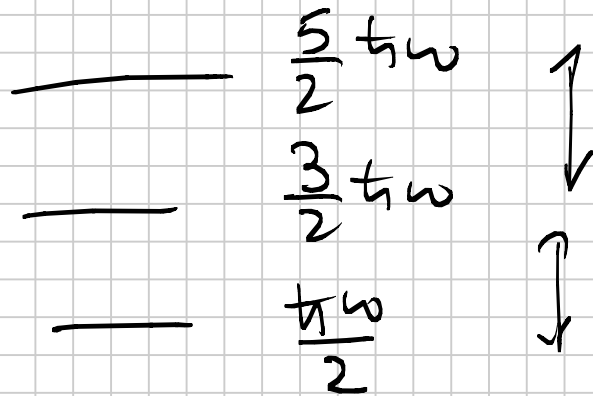
$$\hat{a} = \left(\frac{\hat{x}}{x_0} \right) + i \left(\frac{\hat{p}}{p_0} \right)$$

$$x_0 = \sqrt{2m\hbar\omega}$$

$$\hat{a}^{\dagger} = \left(\frac{\hat{x}}{x_0} \right) - i \left(\frac{\hat{p}}{p_0} \right)$$

$$p_0 = \frac{2\hbar}{M\omega}$$

$$H = \hbar\omega \left(\hat{a}^{\dagger} \hat{a} + \frac{1}{2} \right)$$



$\hbar\omega(q)$

FOR EACH q

$$H = \sum_{q \in BZ} \hbar\omega(q) \left(\hat{a}_q^\dagger \hat{a}_q + \frac{1}{2} \right)$$

$\hat{a}_q^\dagger \hat{a}_q \rightarrow$ # of quanta of the harmonic oscillator at q

↓
PHONONS

$$[\hat{a}_q, \hat{a}_q^\dagger] = 1$$

$$[\hat{a}_q, \hat{a}_{q'}^\dagger] = \delta_{qq'}$$

BOSONS

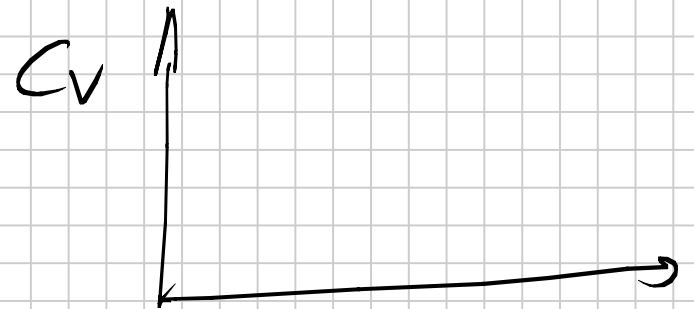
$$\langle E \rangle = \langle H \rangle = \sum_q \hbar \omega(q) \left(\langle \hat{a}_q^\dagger \hat{a}_q \rangle + \frac{1}{2} \right)$$

$$\langle \hat{a}_q^\dagger \hat{a}_q \rangle = \langle N_q \rangle = \frac{1}{\frac{\hbar \omega(q)}{k_B T} - 1}$$

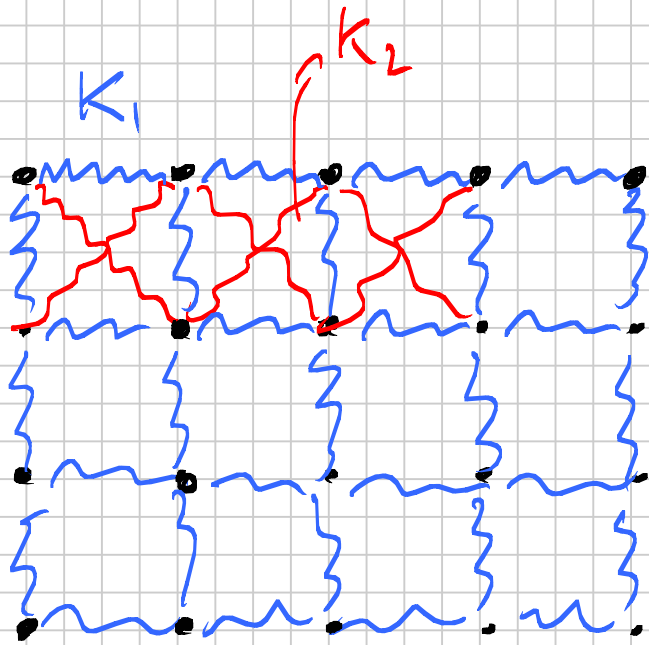
Bose
Statistics

$$\langle E \rangle = \sum_{q \in BZ} \frac{\hbar \omega(q)}{e^{\frac{\hbar \omega(q)}{k_B T}} - 1} \rightarrow E(T)$$

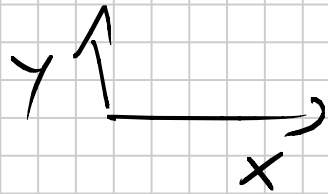
EINSTEIN METHOD
DEBYE METHOD



FR1



ω of normal
mode of excitation
for
 $\vec{q} = \left(\frac{\pi}{a}, 0\right)$



→ LONGITUDINAL ω_L
→ TRANSVERSAL ω_T