
PHY492: Nuclear & Particle Physics

Lecture 20

HW-5

Neutrinos and neutrino oscillations

HW-5

11.1 $J^P = 1^- : \rho^0(770) \rightarrow \pi^+ + \pi^-$, what about $\rho^0(770) \rightarrow \pi^0 + \pi^0$?

Isospin spin

$$\begin{aligned} \rho^0(770) \rightarrow \pi^+ + \pi^- & : 1,0 \rightarrow 1,+1; 1,-1 = \sqrt{\frac{1}{2}} \\ \rho^0(770) \rightarrow \pi^0 + \pi^0 & : 1,0 \rightarrow 1,0; 1,0 = 0 \end{aligned}$$

Isospin predicts $\pi^0\pi^0$ is zero.

Spin-Parity: The wave function of the $\pi^0\pi^0$ system must be symmetric, i.e., l even. However, to have $-$ parity, l must be odd. So NO GO.

11.3 In ρ rest frame, let θ be the angle of a pion wrt the ρ momentum direction.

$$J = 1, J_z = 0$$

Wave function of two decay pions contains $Y_{10} = \cos\theta$.

$$\frac{d\sigma}{d\Omega} \sim |Y_{10}|^2 \sim \cos^2\theta$$

$$J = 1, J_z = \pm 1$$

Wave function of two decay pions contains $Y_{11} = \sin\theta$.

$$\frac{d\sigma}{d\Omega} \sim |Y_{11}|^2 \sim \sin^2\theta$$

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11.5

J^{PC} for $\gamma : 1^{--}$; $\pi^0 : 0^{-+}$; $\omega^0 : 1^{--}$; $\eta' : 0^{-+}$; $\rho^0 : 1^{--}$; $J/\psi : 1^{--}$

$\omega^0 \rightarrow \pi^0 + \gamma$; allowed because both sides have $C = -1$

$\eta' \rightarrow \rho + \gamma$ allowed because both sides have $C = +$

$\pi^0 \rightarrow \gamma + \gamma + \gamma$ not allowed, 3 gammas have charge conjugation $-$, pion is $+$

$J/\psi \rightarrow p + \bar{p}$ allowed. J/ψ & $p\bar{p}$ will have Parity $-$, state is 3S_1 (anti-symmetric)

in 1S_0 , $C(p\bar{p}) = +(p\bar{p})$, while in 3S_1 , $C(p\bar{p}) = -(p\bar{p})$

$\rho^0 \rightarrow \gamma + \gamma$ not allowed, 2 gammas have charge conjugation $+$.

HW-5

12.1 Probability vs. time of K^0 beam yielding a \bar{K}^0 .

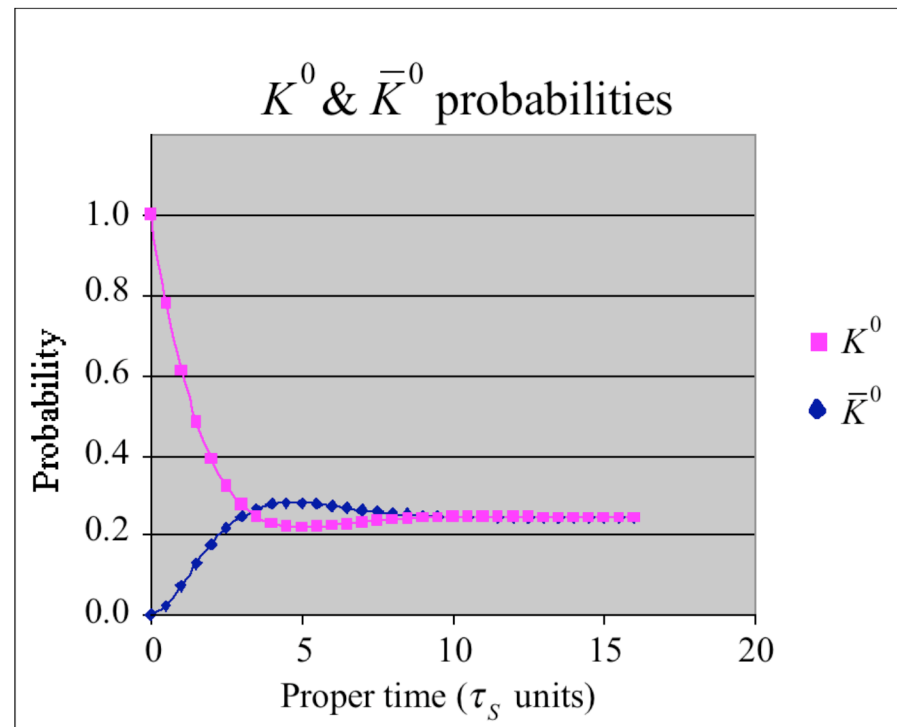
$$\tau_s = 0.9 \times 10^{-10} \text{ s}, \tau_L = 5 \times 10^{-8} \text{ s}$$

$$p = 1 + \varepsilon; \quad q = -(1 - \varepsilon)$$

$$\varepsilon = 2.29 \times 10^{-3}$$

$$\Delta mc^2 = 3.5 \times 10^{-12} \text{ MeV}$$

$$P(\bar{K}^0, t) = \left| \frac{q}{p} \right|^2 \left[\frac{1}{4} \exp\left(-\frac{t}{\tau_s}\right) + \frac{1}{4} \exp\left(-\frac{t}{\tau_L}\right) - \frac{1}{2} \exp\left(-\frac{1}{2} \left(\frac{1}{\tau_s} + \frac{1}{\tau_L} \right) t\right) \cos \frac{\Delta mc^2}{\hbar} t \right]$$



Neutrino mass

- Many attempts to see if neutrinos have mass
 - Tritium beta decay end-point
 - electron neutrino mass $< 3 \text{ eV}/c^2$
 - Meson decays
 - muon neutrino mass $< 200 \text{ keV}/c^2$
 - tau neutrino mass $< 18 \text{ MeV}/c^2$
 - Cosmological limits
 - lots of assumptions about big bang
 - nuclear-synthesis, more assumptions
 - (number density) \times (mass) in universe constraint
 - neutrino masses $< 1 \text{ eV}$
- Theoretical bias for non-zero neutrino mass
 - Why are all neutrinos left-handed?
 - What happened to all the right hand neutrinos?
 - "See-Saw" mechanism
 - right hand neutrinos forced to be VERY MASSIVE
 - and left hand neutrinos to be very light (but not zero).

Propagation of neutrino mass states

- Remember, Schrodinger wave equation solutions ?
- Remember, time dependent Schrodinger wave equation solutions ?

$$e^{-i(\omega t - kx)}; \quad \omega = \frac{E}{\hbar}; \quad k = \frac{p}{\hbar}$$

$$\psi(x, t) = \psi(x, 0) e^{-\frac{i}{\hbar}(Et - px)}; \quad \text{let } x = L, \quad t = L / c$$

$$\psi(L) = \psi(0) e^{-\frac{i}{\hbar c}(E - pc)L}$$

$$\psi(L) = \psi(0) e^{-\frac{i}{\hbar c} \frac{m^2 c^4}{2E} L}$$

$$pc = E \sqrt{1 - \frac{m^2 c^4}{E^2}} \approx E - \frac{m^2 c^4}{2E}, \quad E \gg mc^2$$

$$E - pc \approx \frac{m^2 c^4}{2E}$$

- Phase factor $e^{-\frac{i}{\hbar c} \frac{m^2 c^4}{2E} L}$ depends on distance L from production, particle energy, and mass **squared !!**

Neutrino masses and flavors

- Neutrino mass states (H_{free}) and flavor states (H_{weak}) are **different!**
- Remember Cabibbo mixing of charge $-1/3$ quark states?

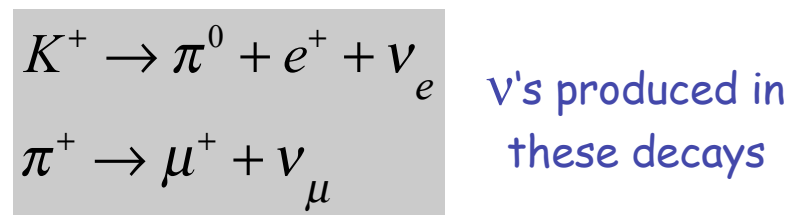
$$\begin{array}{c} \text{weak} \\ \text{eigenstates} \end{array} \begin{pmatrix} d' \\ s' \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} d \\ s \end{pmatrix} \begin{array}{c} \text{mass} \\ \text{eigenstates} \end{array} \quad \text{Cabibbo angle } \theta$$

- A neutrino flavor state is a mixture of 3 mass states (**simplify: use 2**)

$$\begin{array}{c} \text{weak} \\ \text{eigenstates} \end{array} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} \begin{array}{c} \text{mass, } m_1, m_2 \\ \text{eigenstates} \end{array} \quad \text{Mixing angle } \theta$$

- At $x = 0, t = 0$ an electron or muon neutrino is a mass state mixture.

$$\begin{aligned} |\nu_e\rangle &= |\nu_1\rangle \cos \theta + |\nu_2\rangle \sin \theta \\ |\nu_\mu\rangle &= |\nu_2\rangle \cos \theta - |\nu_1\rangle \sin \theta \end{aligned}$$



- Attempts (really can't do it) to measure the mass of the electron neutrino near $t = 0$ would give one of two possible answers:

m_1 , with probability $\cos^2 \theta$, or m_2 , with probability $\sin^2 \theta$.

Propagation of mass eigenstates

$$e^{-\frac{i}{\hbar c} \frac{m^2 c^4}{2E} L}$$

- What happens to pure muon neutrinos a distance L away?

$$|v_\mu(L)\rangle = e^{-\frac{i}{\hbar c} \frac{m_2^2 c^4}{2E} L} |v_2\rangle \cos \theta - e^{-\frac{i}{\hbar c} \frac{m_1^2 c^4}{2E} L} |v_1\rangle \sin \theta$$

$$\begin{aligned} |v_1\rangle &= |v_e\rangle \cos \theta - |v_\mu\rangle \sin \theta \\ |v_2\rangle &= |v_\mu\rangle \cos \theta + |v_e\rangle \sin \theta \end{aligned}$$

- Masses are different, so phases will be **different!**
- Neutrinos become a **mixture** of muon and electron neutrinos.

$$|v_\mu(L)\rangle = (e^{-\frac{i}{\hbar c} \frac{m_2^2 c^4}{2E} L} \cos^2 \theta + e^{-\frac{i}{\hbar c} \frac{m_1^2 c^4}{2E} L} \sin^2 \theta) |v_\mu\rangle + (e^{-\frac{i}{\hbar c} \frac{m_2^2 c^4}{2E} L} \cos \theta \sin \theta - e^{-\frac{i}{\hbar c} \frac{m_1^2 c^4}{2E} L} \sin \theta \cos \theta) |v_e\rangle$$

$$\langle v_\mu | v_\mu(L) \rangle = (e^{-\frac{i}{\hbar c} \frac{m_2^2 c^4}{2E} L} \cos^2 \theta + e^{-\frac{i}{\hbar c} \frac{m_1^2 c^4}{2E} L} \sin^2 \theta)$$

$\sin^2 2\theta$ is amplitude

wavelength

$$\lambda(\text{km}) = 2\pi E / (1.27 \Delta m^2)$$

surviving
muon neutrinos

$$|\langle v_\mu | v(L) \rangle|^2 = 1 - \sin^2 2\theta \sin^2 (1.27 \Delta m^2 L / E)$$

$$\begin{aligned} \Delta m^2 &\text{ in eV}^2 \\ L &\text{ in km} \\ E &\text{ in GeV} \end{aligned}$$

appearing
electron neutrino

$$|\langle v_e | v(L) \rangle|^2 = \sin^2 2\theta \sin^2 (1.27 \Delta m^2 L / E)$$

Calculate muon neutrino survival probability

Messy trigonometry and algebra:

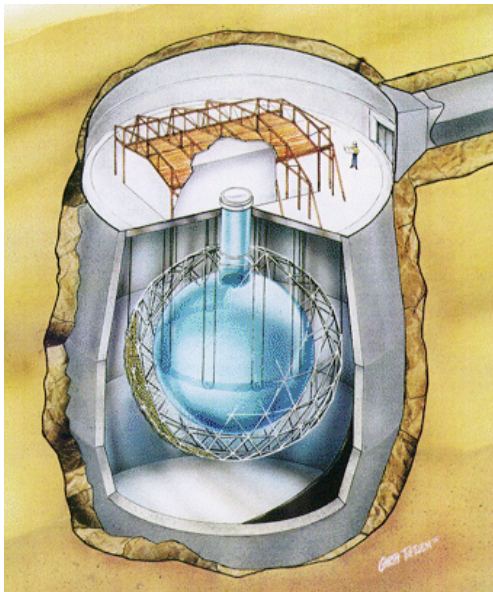
$$\begin{aligned} \left| \langle \nu_\mu | \nu_\mu(L) \rangle \right|^2 &= \left(e^{-\frac{i}{\hbar c} \frac{m_2^2 c^4}{2E} L} \cos^2 \theta + e^{-\frac{i}{\hbar c} \frac{m_1^2 c^4}{2E} L} \sin^2 \theta \right) \\ &\quad \times \left(e^{+\frac{i}{\hbar c} \frac{m_2^2 c^4}{2E} L} \cos^2 \theta + e^{+\frac{i}{\hbar c} \frac{m_1^2 c^4}{2E} L} \sin^2 \theta \right) \\ &= \cos^4 \theta + \sin^4 \theta + \left(e^{+\frac{i}{\hbar c} \frac{(m_2^2 - m_1^2) c^4}{2E} L} + e^{-\frac{i}{\hbar c} \frac{(m_2^2 - m_1^2) c^4}{2E} L} \right) \cos^2 \theta \sin^2 \theta \\ &= \cos^2 \theta (1 - \sin^2 \theta) + \sin^2 \theta (1 - \cos^2 \theta) + \frac{2}{2} \cos^2 \theta \sin^2 \theta \\ &= 1 - 2 \sin^2 \theta \cos^2 \theta \left(1 - \cos \left(\frac{(m_2^2 - m_1^2) c^4 L}{2E \hbar c} \right) \right) \\ &= 1 - \sin^2 2\theta \left(\sin^2 \left(\frac{(m_2^2 - m_1^2) c^4 L}{4E \hbar c} \right) \right) \end{aligned}$$

Solar and cosmic ray oscillation experiments

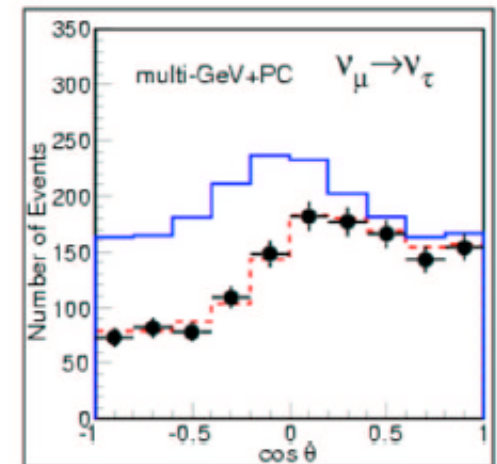
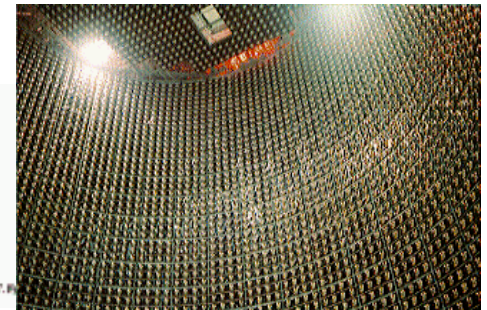
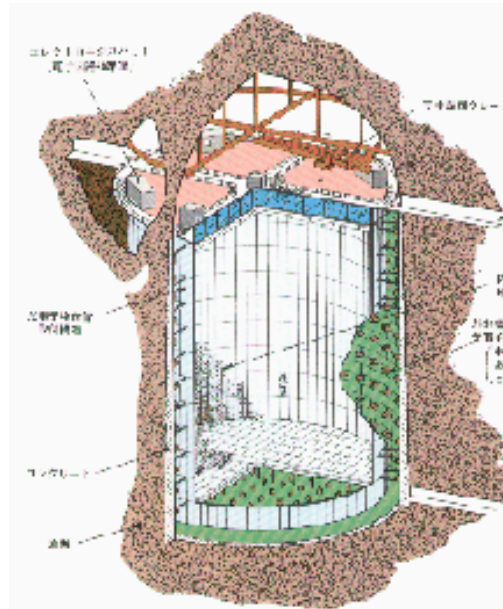
- Three neutrino flavors, three mass states, the mixing matrix is 3×3
- **Three Δm^2** and **two angles** (one can't be measured this way) determined by solar and cosmic ray neutrino experiments

solar neutrinos
Sudbury (SNO)

1kT D_2O



solar and cosmic ray
Super-kamiokanda 50 kT H_2O



Three ν flavors and three ν masses

- As is the case for quarks, the mixing matrix is 3×3
- Very different from quarks, off-diagonal mixing angles are LARGE.

Missing one crucial angle, $\sin\theta_{13}$, and a CP violating phase δ

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = M \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$|M| \approx \begin{pmatrix} 0.83 & 0.55 & s_{13}e^{i\delta} \\ 0.39 & 0.59 & .7 \\ 0.39 & 0.59 & .7 \end{pmatrix}; \quad s_{13} < 0.25 \\ \delta \sim ?$$

- Result for muon neutrinos oscillating to electron neutrinos is the same as in the two neutrino case.

appearing

electron neutrinos

$$\left| \langle \nu_e | \nu(L) \rangle \right|^2 = \sin^2 2\theta_{13} \sin^2 \left(1.27 \Delta m_{23}^2 L / E \right)$$

$$\Delta m_{23}^2 \sim 3 \times 10^{-3} \text{ eV}^2; \quad \text{pick } E \sim 2 \text{ GeV}$$

$$\lambda = \frac{2\pi E}{1.27 \Delta m_{23}^2} = \frac{4\pi}{4 \times 10^{-3}} \text{ km} \approx 3000 \text{ km}$$

First maximum at $\lambda / 4 \approx 750 \text{ km}$

Fermilab neutrino beam
points to Sudan MN which
is 750 km away

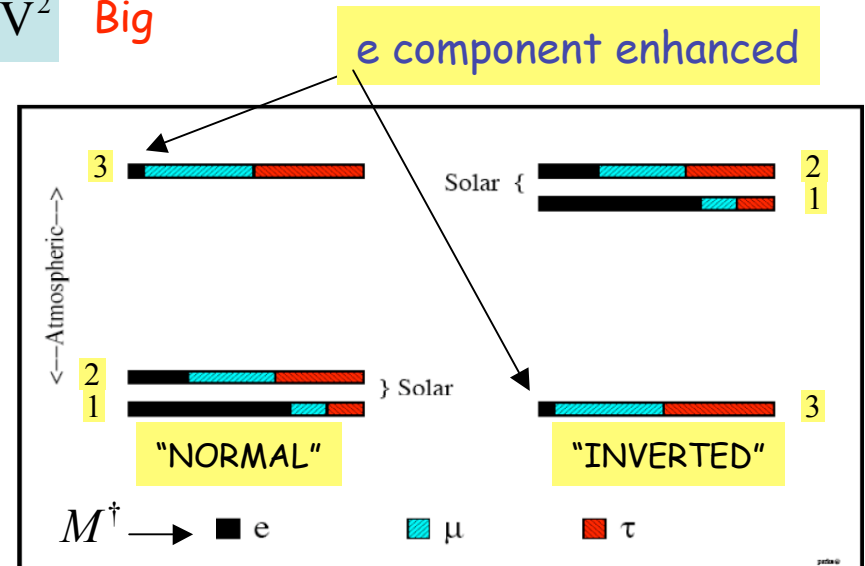
Mass hierarchy of three mass states

- Mass splitting $\Delta m_{ij}^2 = |m_j^2 - m_i^2|$

Solar: $\Delta m_{12}^2 \sim 5 \times 10^{-5} \text{ eV}^2$ **Small**

Atmospheric: $\Delta m_{23}^2 \sim 2.5 \times 10^{-3} \text{ eV}^2$ **Big**

- Mass hierarchy **unknown**
 - "NORMAL" as in quark masses
 - "INVERTED" as in P-states?
 - Can't do CP violation on earth without hierarchy resolution



- MSW effect
 - In a long baseline experiment ν_e and $\bar{\nu}_e$ passing through the earth have oscillation probability affected with opposite signs.
 - Good for determination of mass hierarchy, but complicates CP violation analysis. Unavoidable, even in Japan.

Rate for muon to electron neutrinos

Appearing
electron neutrino

$$\left| \langle \nu_e | \nu(L) \rangle \right|^2 = \sin^2 2\theta_{13} \sin^2 \left(1.27 \Delta m_{23}^2 L / E \right)$$

- Muon neutrino (anti-neutrino) beam from Chicago to northern MN
- Put 20 kT of liquid scintillator there in the beam
- Look for neutrino interactions that produce **electrons**.

$$\nu_e + N \rightarrow e^- + \text{hadrons}$$

$$\bar{\nu}_e + N \rightarrow e^+ + \text{hadrons}$$

- Will need 5 years to do R&D, design, and construct detector
 - Will need 3-4 years to accumulate data on neutrinos
 - Will need 5 years to accumulate data on anti-neutrinos
 - So look for results on CP violation in leptons 2020.
- New technology: Liquid Argon Time Projection Chamber
 - May be able to do it faster -- I hope.