

## The Simplest Picture of Galaxy Formation and Why It Fails

- Cosmic Microwave Background is smooth to a few parts in  $10^5$

$$\delta\rho/\rho \sim 10^{-4}$$

- Yet high contrast structures (QSOs, galaxies) by  $z \sim 6$ .

$$\delta\rho/\rho \gg 1$$

- Adiabatic perturbations grow as

$$\delta\rho/\rho \propto t^{2/3} \propto R(t) \propto 1/(1+z)$$

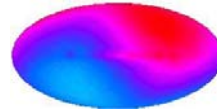
- Expect only

$$\left(\frac{\delta\rho}{\rho}\right)_{QSO} = \frac{(1+z)_{CMB}}{(1+z)_{QSO}} \left(\frac{\delta\rho}{\rho}\right)_{CMB} = \frac{1100}{7} \times 10^{-4} = 0.01$$

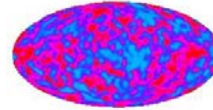
So where did galaxies and clusters come from?



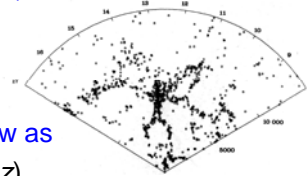
Blue = 0°K  
Red = 4°K



Blue = 2.724°K  
Red = 2.732°K  
Dipole Anisotropy  
~ 1 part in 300



After removing dipole  
Red - blue = 0.0002°K  
~ 1 part in  $10^5$



## In an expanding universe, will a cloud collapse?

The Jeans criterion Version 2:

$$2K < -U$$

Pressure support < gravity

*Collapse if*  $2K < -U$

$$2 \left( \frac{1}{2} M_T v_s^2 \right) < \frac{3}{5} \frac{GM_T^2}{\lambda} \quad \boxed{v_s = \text{sound speed}}$$

$$v_s^2 < \frac{3}{5} \frac{GM_T}{\lambda} = \frac{3G}{5} \frac{(4/3)\pi\lambda^3 \rho_T}{\lambda} = \frac{4}{5} \pi G \lambda^2 \rho_T$$

$$\left( \frac{3M_b}{4\pi\rho_b} \right)^{2/3} = \lambda^2 > \frac{5v_s^2}{4\pi G \rho_T}$$

$$M_{J,b} > \text{const.} \times \frac{\rho_b v_s^3}{\rho_T} = [\text{CO eq. 30.27}]$$

Radiation era

$$v_s = \frac{c}{\sqrt{3}}$$

$$\rho_b \propto R(t)^{-3} \propto T^3$$

$$\rho_T \propto R(t)^{-4} \propto T^4$$

$$M_{J,b} \propto T^{-3}$$

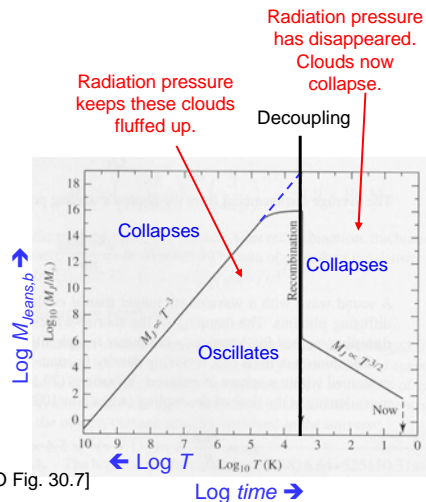
After decoupling

$$v_s = \sqrt{\frac{5kT}{3\mu m_H}}$$

$$\rho_b \propto T^0$$

$$\rho_T \approx \rho_b \propto T^0$$

$$M_{J,b} \propto T^{3/2}$$



[CO Fig. 30.7]

## Q. When do the oscillations start?

When

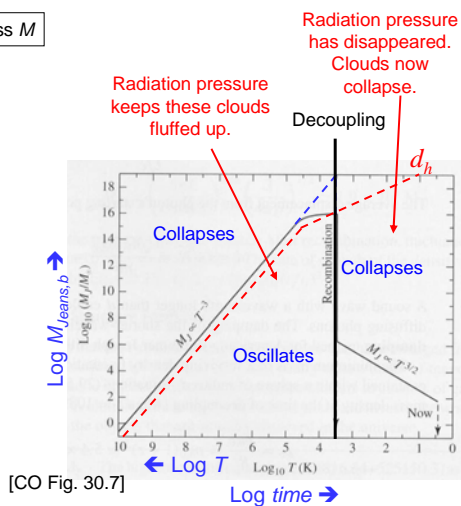
Particle horizon =  $\lambda_M$

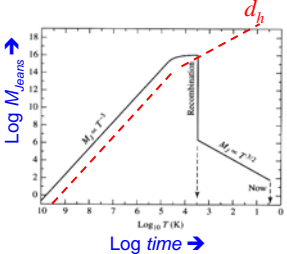
Size scale for mass  $M$

$2K < -U$   
Pressure support < gravity

Before decoupling:

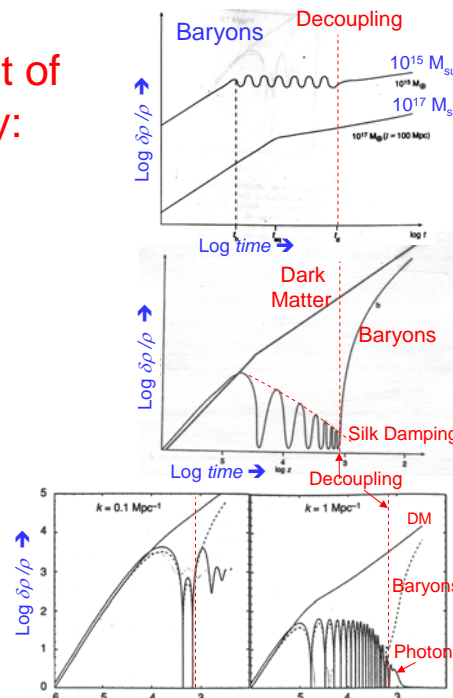
- Particle Horizon  
 $d_h = 2ct \propto R(t)^2 \propto T^{-2}$  (radiation era)
- Proper distance containing mass  $M$   
 $\lambda = (M_b/\rho_b)^{1/3} \propto M_b^{1/3} R(t) \propto M_b^{1/3} T^{-1}$
- Mass for which  $\lambda = d_h$   
 $M_b \propto T^{-3} \propto R^3 \propto t^{3/2}$  (radiation era)  
 $M_b \propto T^{-3/2}$  (matter era)

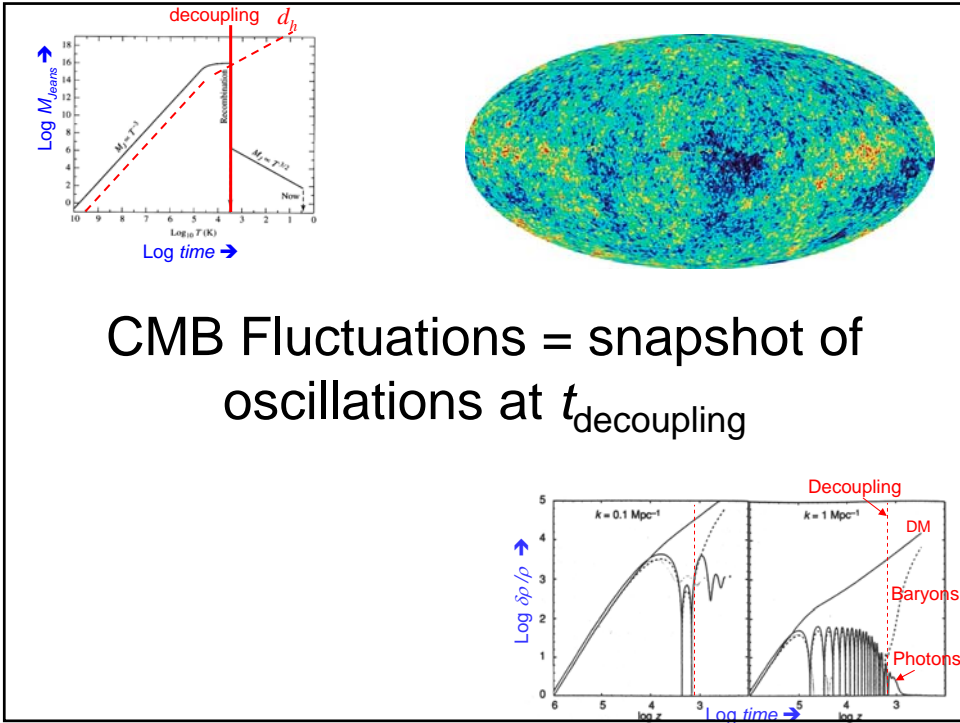




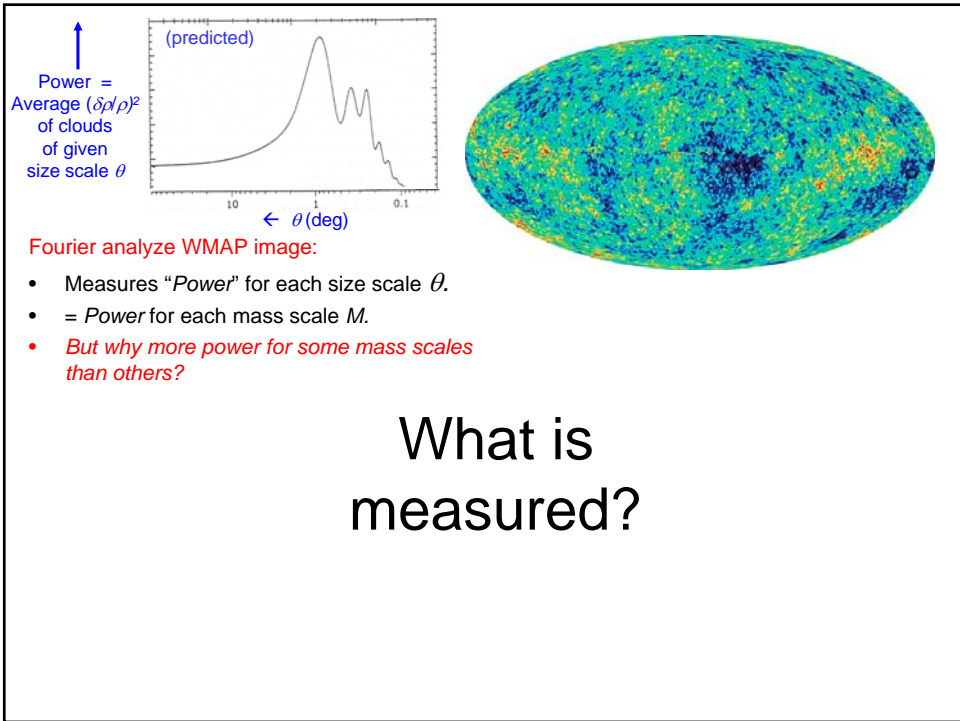
The Rest of the Story:

- Mass for which  $\lambda = d_h$   
 $M \propto T^{-3}$  (rad. era)  
 $M \propto T^{-3/2}$  (matter era)
- At  $t_{decoupling}$  this mass was  $\sim 10^{16} M_\odot$ 
  - $M > 10^{16} M_\odot \rightarrow$  continued growth
  - $M < 10^{16} M_\odot \rightarrow$  oscillations once mass scale comes into particle horizon.
- But Dark Matter not subject to all this.
  - Does not feel radiation pressure.
  - Just collapses away...
- Baryons fall into Dark Matter potential wells as soon as decoupling removes photon pressure support.

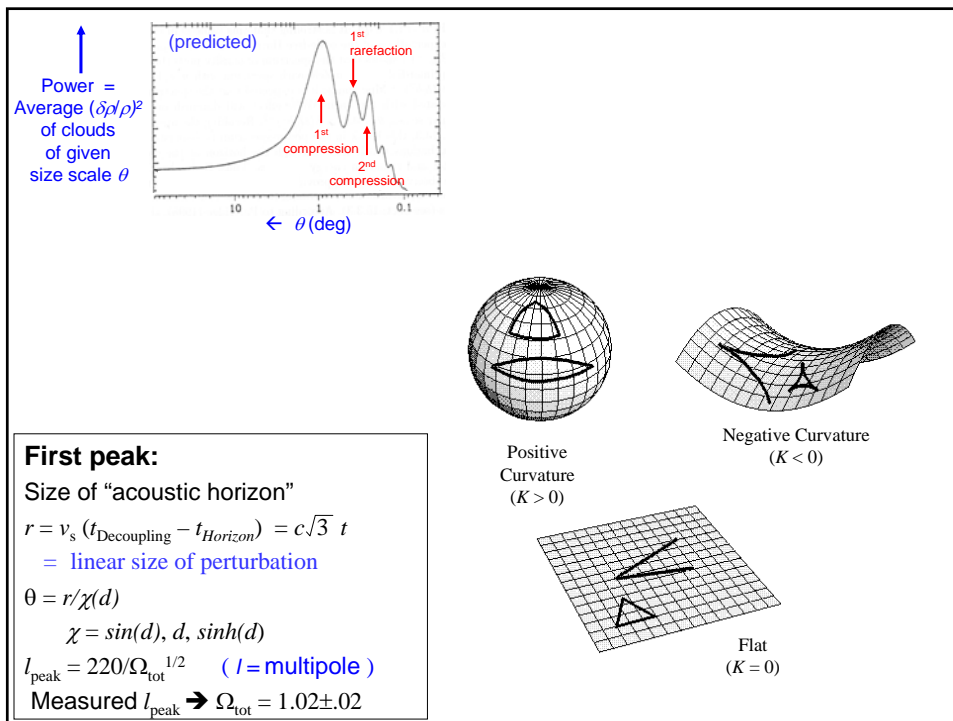
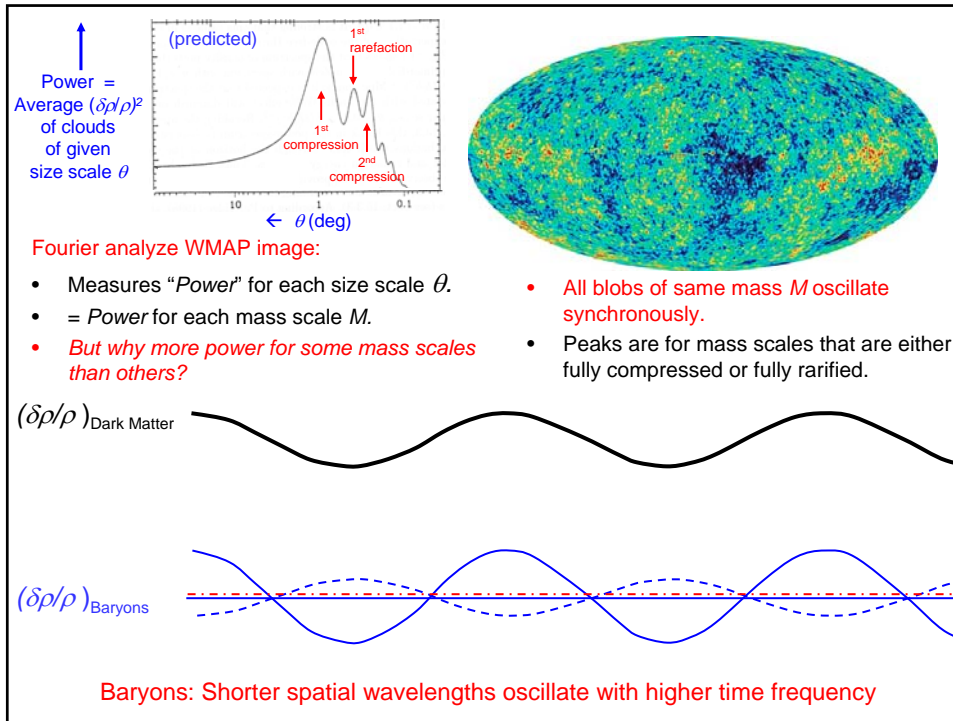




CMB Fluctuations = snapshot of oscillations at  $t_{\text{decoupling}}$



What is measured?

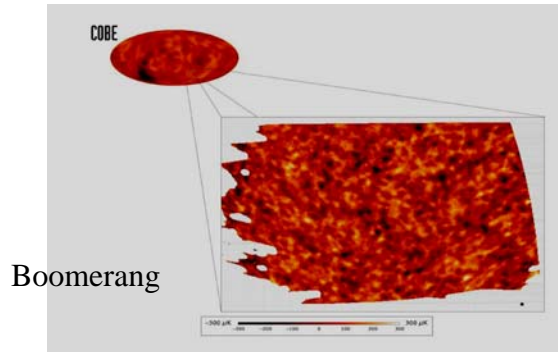


# Boomerang balloon flight (1999)

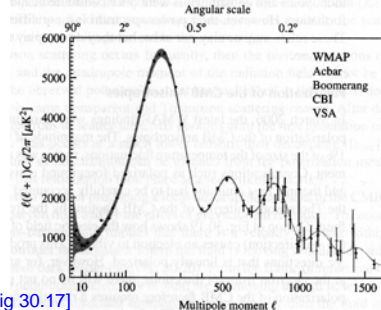
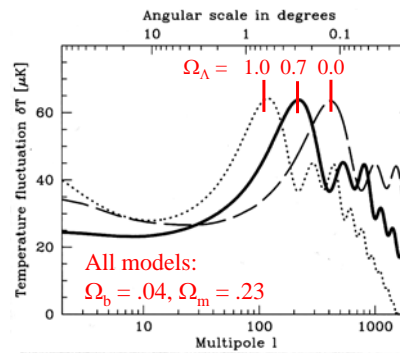
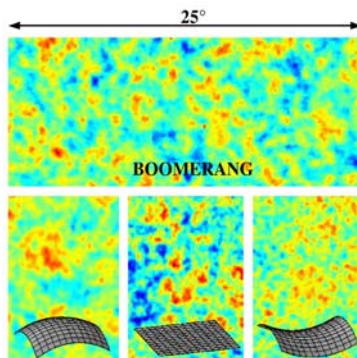
Mapped Cosmic Background Radiation with far higher angular resolution than previously available.



Launch near Mt. Erebus in Antarctica



## Position of 1<sup>st</sup> peak measures curvature



### First peak:

Size of "acoustic horizon"

$$r = v_s (t_{\text{Decoupling}} - t_{\text{Horizon}}) = c\sqrt{3} t$$

= linear size of perturbation

$$\theta = r/\chi(d)$$

$$\chi = \sin(d), d, \sinh(d)$$

$$l_{\text{peak}} = 220/\Omega_{\text{tot}}^{1/2} \quad (l = \text{multipole})$$

Measured  $l_{\text{peak}} \rightarrow \Omega_{\text{tot}} = 1.02 \pm 0.02$

[CO fig 30.17]

## The "Concordance" Cosmology (or $\Lambda$ CDM)

- Type Ia Supernovae as "standard candles"
  - accelerating expansion
  - $q_0 = \Omega_m/2 - \Omega_\Lambda$
- CMB anisotropy →  $\Omega_{\text{total}} = \Omega_m + \Omega_\Lambda$
- Can solve for  $\Omega_m$ ,  $\Omega_\Lambda$

Another independent measure:  
Rate of galaxy cluster evolution

