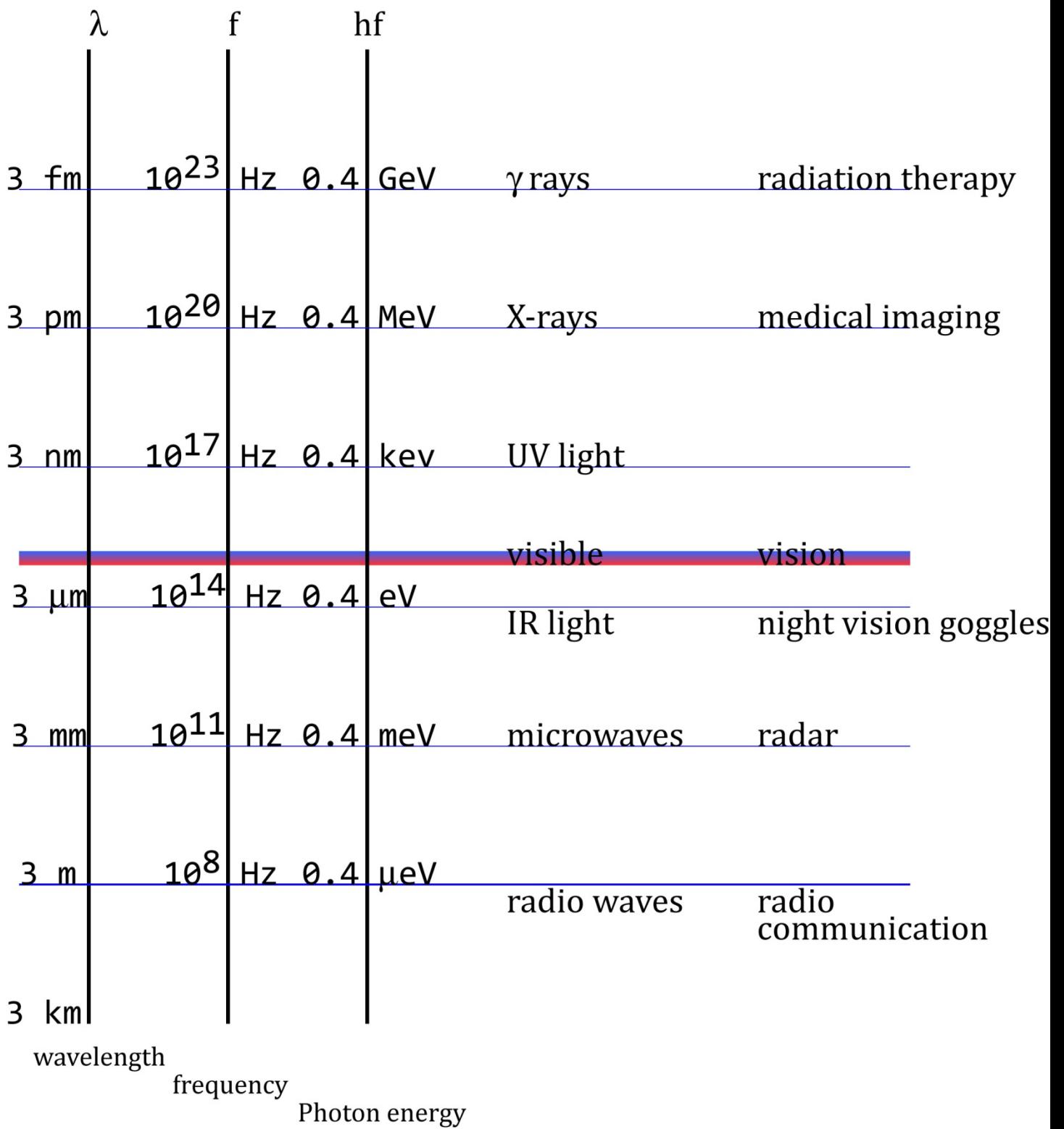


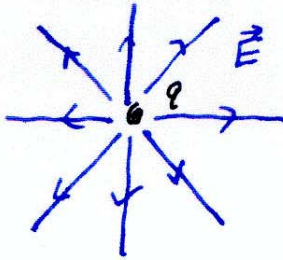
The Electromagnetic Spectrum



Maxwell's Equations

Pictorial

$$\nabla \cdot \vec{E} = \rho / \epsilon_0$$



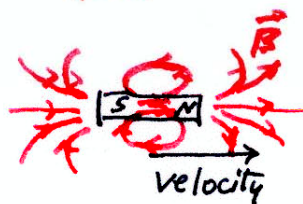
$$\vec{E} = \frac{q\hat{r}}{4\pi\epsilon_0 r^2}$$

$$\nabla \cdot \vec{B} = 0$$



The magnetic field is solenoidal.

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$



$$\mathcal{E}_{\text{mf}} = -\frac{d\Phi}{dt}$$

Faraday, 1831

$$\nabla \times \vec{B} = \mu_0 (\vec{J} + \vec{J}_D)$$

$$\vec{J}_D = \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

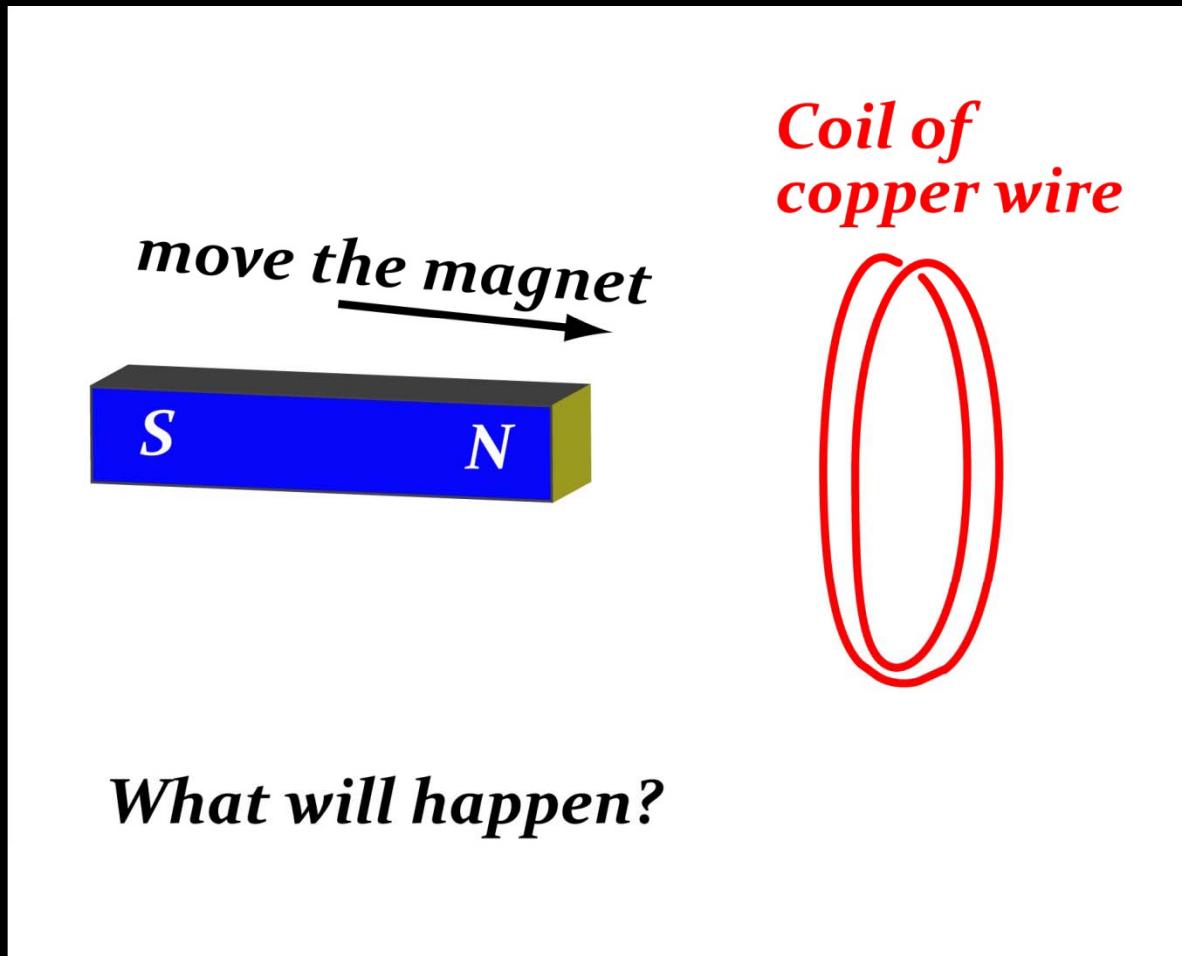


$$\vec{B} = \frac{\mu_0 I}{2\pi r} \hat{\phi}$$

Displacement Current
Maxwell, 1864

These field equations, published in 1864, are still used today. They describe electricity, magnetism, electromagnetism, and optics.

Electromagnetic Induction



$$\text{EMF} = - d\Phi / dt$$

Electromagnetic Waves in Field Theory

We consider \vec{E} and \vec{B} in empty space

$$\nabla \cdot \vec{E} = 0 \quad \text{and} \quad \nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \text{and} \quad \nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Now, construct solutions w/ $\vec{E}(\vec{x}, t) = \hat{z} F(z, t)$

i.e., propagating in the z direction, polarized in the x direction.

$$\bullet \quad \nabla \cdot \vec{E} = 0 \rightarrow 0 = 0 \quad \checkmark$$

$$\bullet \quad \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \rightarrow \begin{vmatrix} \hat{z} & \hat{y} & \hat{x} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ F & 0 & 0 \end{vmatrix} = +\hat{y} \frac{\partial F}{\partial z} = -\frac{\partial \vec{B}}{\partial t}$$

$$\therefore \text{require } \vec{B}(\vec{x}, t) = \hat{y} G(z, t) \quad \text{and} \quad \frac{\partial F}{\partial z} = -\frac{\partial G}{\partial t}$$

$$\bullet \quad \nabla \cdot \vec{B} = 0 \rightarrow 0 = 0 \quad \checkmark$$

$$\bullet \quad \nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \rightarrow \begin{vmatrix} \hat{z} & \hat{y} & \hat{x} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ 0 & G & 0 \end{vmatrix} = -\hat{x} \frac{\partial G}{\partial z} = \mu_0 \epsilon_0 \hat{z} \frac{\partial F}{\partial t}$$

$$\therefore \text{requires } \frac{\partial G}{\partial z} = -\mu_0 \epsilon_0 \frac{\partial F}{\partial t}$$

Electromagnetic Waves in Field Theory

$$\vec{E}(\vec{x}, t) = \hat{i} F(z, t) \quad \text{and} \quad \vec{B}(\vec{x}, t) = \hat{j} G(z, t)$$

$$\frac{\partial F}{\partial z} = -\frac{\partial G}{\partial t} \quad \text{and} \quad \frac{\partial G}{\partial z} = -\mu_0 \epsilon_0 \frac{\partial F}{\partial t}$$

Solutions

Consider $\frac{\partial^2 F}{\partial z^2} = -\frac{\partial^2 G}{\partial z \partial t} = -\frac{\partial^2 G}{\partial t \partial z} = \mu_0 \epsilon_0 \frac{\partial^2 F}{\partial t^2}$

$$\frac{\partial^2 F}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2 F}{\partial t^2} \quad (\text{wave equation}) \quad \text{w/} \quad c^2 = \frac{1}{\mu_0 \epsilon_0}$$

wavespeed = c ↗

A general solution is $F(z, t) = f(z - ct)$

where $f(\cdot)$ is any function.

$$\left[\frac{\partial^2 F}{\partial z^2} = f''(z-ct) \quad \text{and} \quad \frac{\partial^2 F}{\partial t^2} = c^2 f''(z-ct) \quad \checkmark \right]$$

OK, now what is $G(z, t)$?

$$\frac{\partial G}{\partial t} = -\frac{\partial F}{\partial z} \Rightarrow -f'(z-ct)$$

$$\therefore G(z, t) = \frac{1}{c} f(z-ct)$$

$$G(z, t) = \frac{1}{c} F(z, t)$$

$$B = \frac{1}{c} E$$

Harmonic Solutions

$$f(z) = E_0 \sin(kz)$$

$$\vec{E}(\vec{x}, t) = \hat{i} E_0 \sin(kz - \omega t)$$

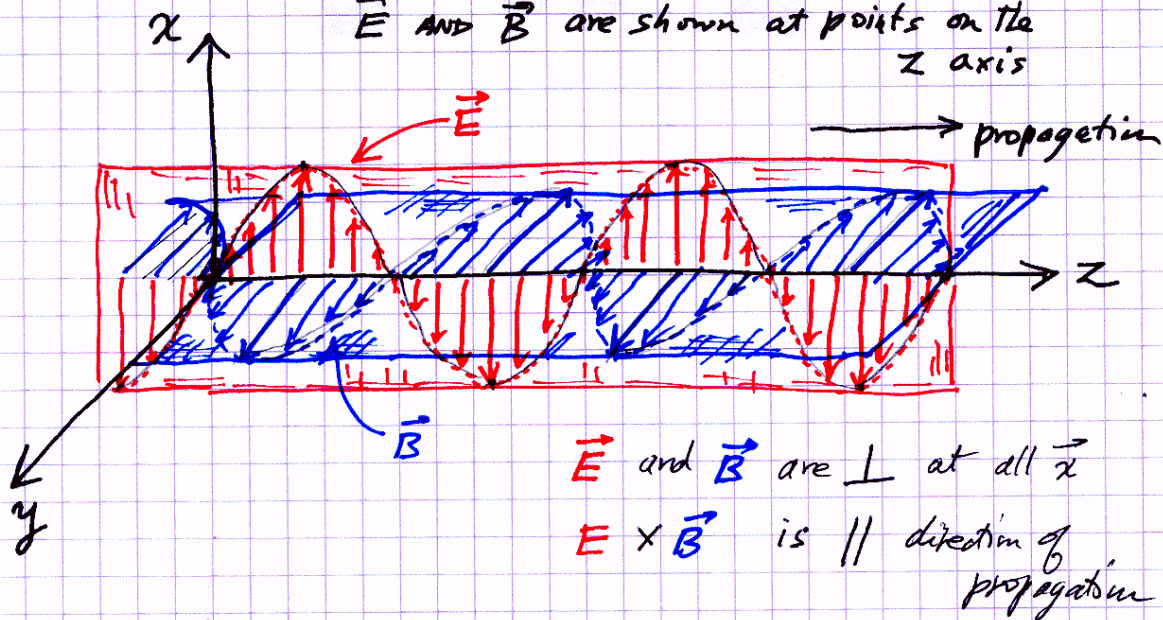
$$\vec{B}(\vec{x}, t) = \hat{j} \frac{E_0}{c} \sin(kz - \omega t)$$

$$\left. \begin{array}{l} \omega = ck \\ B_0 = E_0/c \end{array} \right\}$$

The Harmonic Plane Wave

HARMONIC ELECTROMAGNETIC WAVE

\vec{E} AND \vec{B} are shown at points on the z axis



$$\vec{E}(\vec{x}, t) = \hat{i} E_0 \sin(kz - \omega t)$$

$$\vec{B}(\vec{x}, t) = \hat{j} B_0 \sin(kz - \omega t)$$

Wave properties

λ : wavelength

$$k\lambda = 2\pi \quad \text{or} \quad \lambda = \frac{2\pi}{k}$$

T : period

$$\omega T = 2\pi \quad \text{or} \quad T = \frac{2\pi}{\omega}$$

f : frequency

$$f = \frac{1}{T} = \frac{\omega}{2\pi}$$

c : velocity

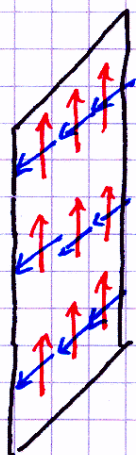
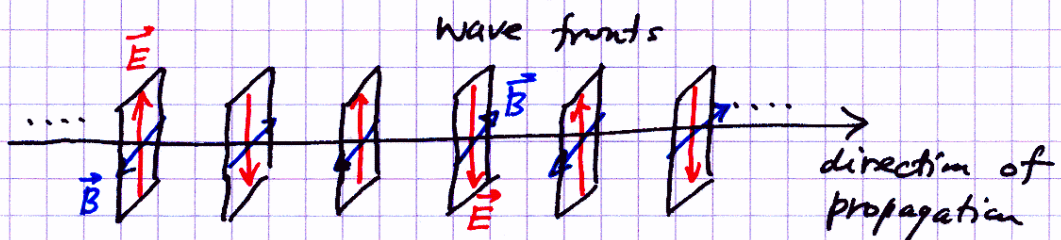
$$c = \frac{\lambda}{T} = \frac{\omega}{k}$$

Maxwell's equations require

- $B_0 = E_0/c$
- $\vec{E}, \vec{B}, \vec{k}$ form an orthogonal triad

The plane wave is an idealization —

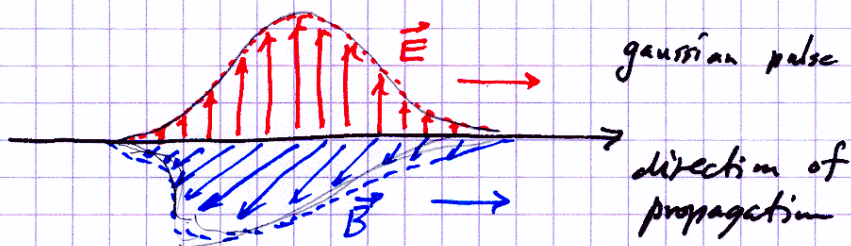
\vec{E} and \vec{B} are coherent on infinite planar wave fronts



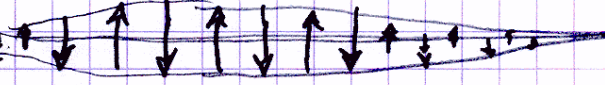
\vec{E} and \vec{B} are coherent on infinite planar wave fronts
 $\vec{E} \times \vec{B}$ is \parallel to \vec{k}

A finite electromagnetic wave; a "pulse"

e.g., $\vec{E}(\vec{x}, t) = \hat{i} f(z-ct)$ and $\vec{B}(\vec{x}, t) = \hat{j} f(z-ct)/c$



But

If ~~the~~ \vec{E} and \vec{B} oscillate  with ~~the same phase~~, then the plane wave is a good approximation.