

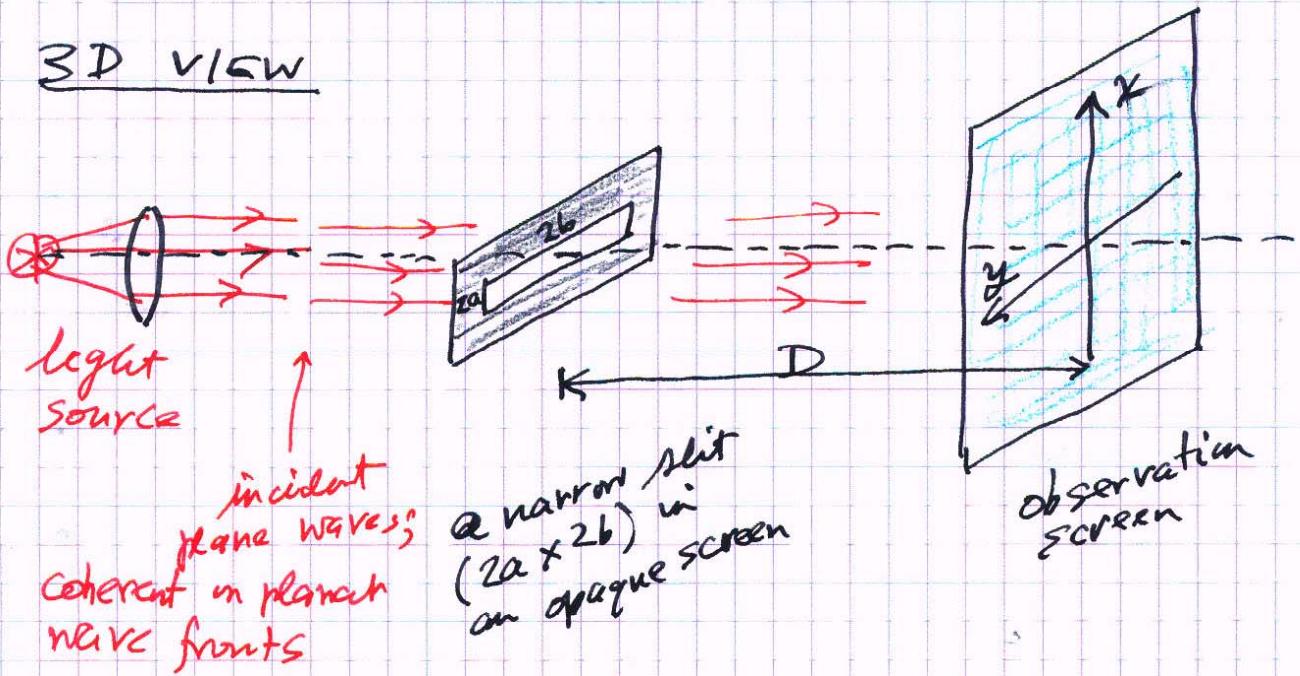
## 7. Diffraction

B3/1

Diffraction: Bending of light, and the related interference, as light waves pass an edge or obstruction.

### Single-slit diffraction

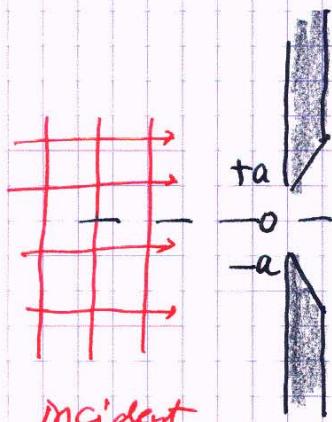
3D view



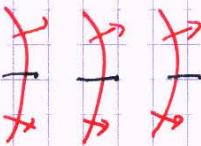
What is the light intensity  $I(x, y; D)$  on the observation screen? The answer depends on  $D$ .

B3/2

### SIDE VIEW

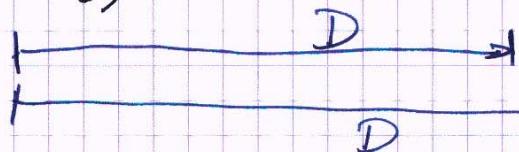


Shadow

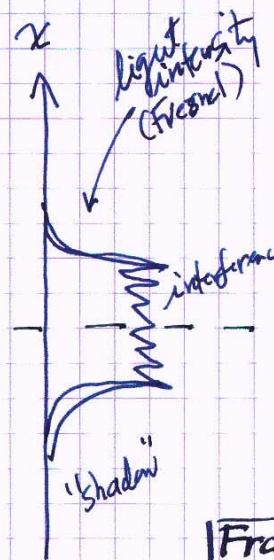


Shadow

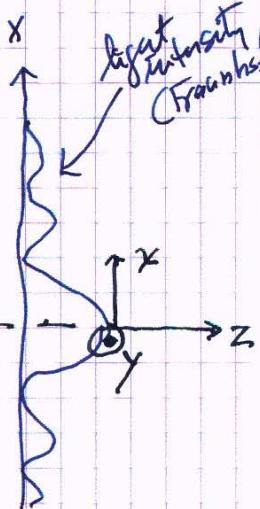
single  
slit  
( $2a \times 2b$ )



Fresnel  
Region

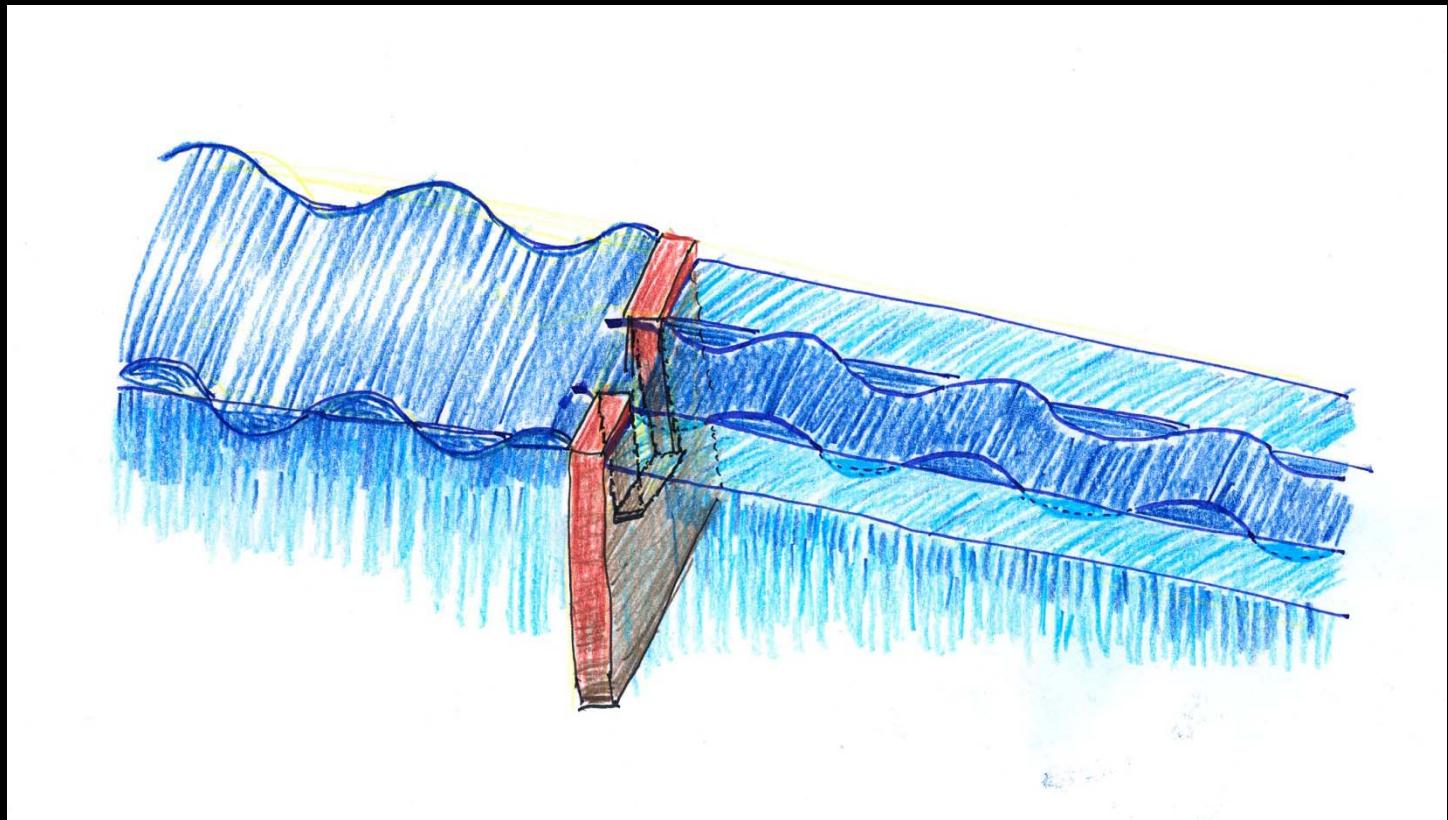


Fraunhofer  
Region



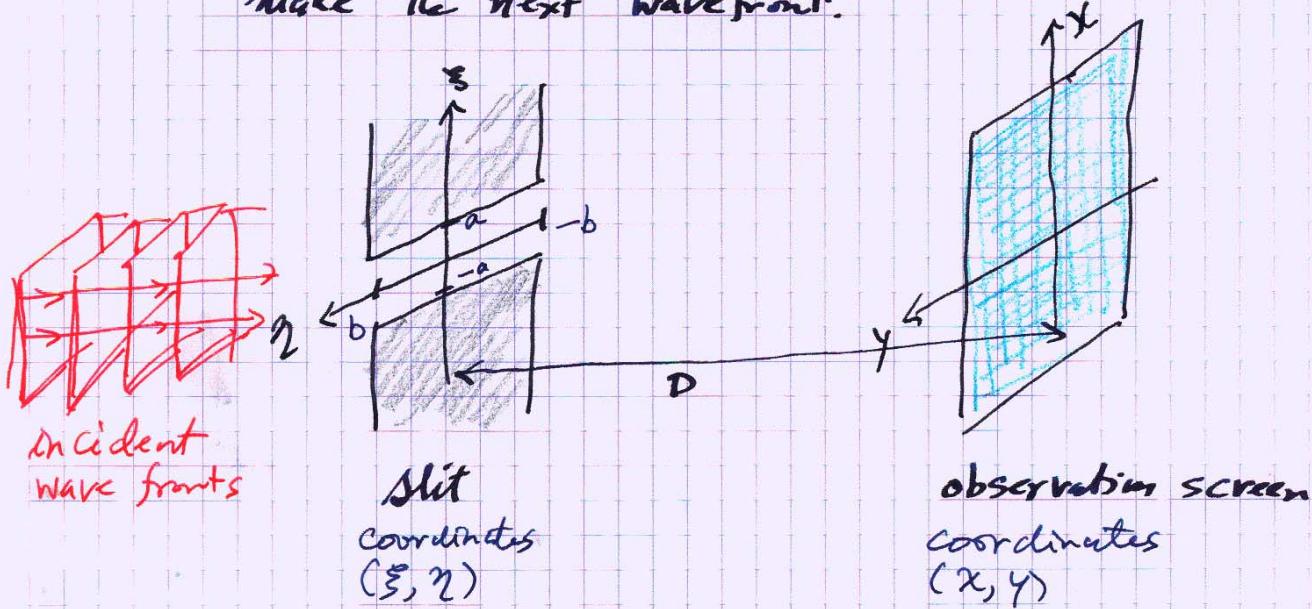
In either case,  $D$  is large compared to  $a$   
or  $\lambda$ .  $D$  is larger in Fraunhofer region.

# Why diffraction must occur!



The medium cannot support a discontinuity.

Huygens' Principle Each point on a wave front acts as a source of "spherical wavelets" which propagate to make the next wavefront.



The Fresnel - Kirchhoff diffraction integral

$$\phi(x, y, D) = C \int_{\text{slit area}} \frac{e^{ikr}}{r} \frac{1 + \cos \theta}{2} dA \cdot e^{-iwt}$$

↑ field  
↑ slit area  
↑ amplitude factor  
↑ spherical wave

Superposition  
"obliquity factor"

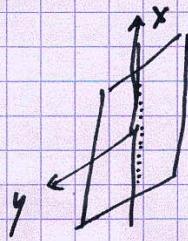
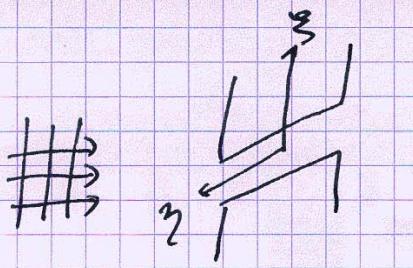
$$r = \sqrt{(x-\xi)^2 + (y-\eta)^2 + D^2}$$

$$r \approx D + \frac{(x-\xi)^2 + (y-\eta)^2}{2D} \quad \text{**}$$

Taylor series  
 $\sqrt{D^2 + \epsilon} \approx D + \frac{\epsilon}{2D}$

\*\*

The wavelets are not spherical (Stokes, 1849)



B3/4

Consider intensity  $I(x, 0, D)$  on the  $x$  axis ( $y=0$ )

$$\phi(x) = C \int \frac{e^{ikr}}{r} \frac{1 + \cos\theta}{2} d\xi dy \cdot e^{-iwt}$$

Approximations for large  $D$ :

- $r = \sqrt{(x-\xi)^2 + y^2 + D^2} = D + \frac{(x-\xi)^2 + y^2}{2D}$
- $\frac{1 + \cos\theta}{2} = 1$  if  $\theta$  is small
- $y_r = \frac{y}{D}$
- $e^{ikr} = e^{ikD} e^{ik\frac{y^2}{2D}} e^{ik\frac{(x-\xi)^2}{2D}}$

$$\phi(x) = M \int_{-\alpha}^{\alpha} e^{ik\frac{(x-\xi)^2}{2D}} d\xi e^{-iwt}$$

amplitude  
independent of  $x$   $\leftarrow k = 2\pi/\lambda$

$$\text{Let } \frac{\pi}{2} u^2 = \frac{k(x-\xi)^2}{2D} \text{ i.e., } u = \sqrt{\frac{2}{\lambda D}} (x-\xi)$$

$$du = -\sqrt{\frac{2}{\lambda D}} d\xi$$

$$\phi(x) = M \sqrt{\frac{\lambda D}{2}} \int_{\sqrt{\frac{2}{\lambda D}}(x-\alpha)}^{\sqrt{\frac{2}{\lambda D}}(x+\alpha)} e^{i\frac{\pi}{2}u^2} du$$

$$g = \sqrt{\frac{2}{\lambda D}}$$

$$y = \sqrt{\frac{2}{\lambda D}} \quad B3/5$$

$$\phi(x) = \frac{M}{q} \int_{g(x-a)}^{g(x+a)} e^{i \frac{\pi}{\lambda D} u^2} du$$

$$\cos \frac{\pi}{2} u^2 + i \sin \frac{\pi}{2} u^2$$

(Euler)

$$\phi(x) = \frac{M}{q} \left\{ C[g(x+a)] + i S[g(x+a)] - C[g(x-a)] - i S[g(x-a)] \right\}$$

### Fresnel Integrals

$$C(x) = \int_0^x \cos \frac{\pi}{2} u^2 du \quad \text{"Fresnel C"}$$

$$S(x) = \int_0^x \sin \frac{\pi}{2} u^2 du \quad \text{"Fresnel S"}$$

The transcendental functions are available in Mathematica.

Intensity on the observation screen  $x, y, z$

$$I(x) \propto |\phi(x)|^2$$

$$= (x, 0, D)$$

$$I(x) = I_0 \frac{R^2(x)}{R^2(0)}$$

$$R^2(x) = (C[g(x+a)] - C[g(x-a)])^2 + (S[g(x+a)] - S[g(x-a)])^2$$

## Fraunhofer approximation for D very large

$$\phi(x) = \frac{M}{q} \int_{q(x-a)}^{q(x+a)} e^{i\frac{\pi}{2}u^2} du$$

where  $q = \sqrt{\frac{2}{\lambda D}}$  is small

- Write  $u = qx + \xi$  when  $-qa \leq \xi \leq qa$
- approximate

$$u^2 = q^2x^2 + 2qx\xi + \xi^2 \approx q^2x^2 + 2qx\xi$$

- and use

$$\int_{-qa}^{qa} e^{i\frac{\pi}{2} \cdot 2qx\xi} d\xi = \frac{2i \sin[\pi q^2 a x]}{i\pi q x}$$

$$\begin{aligned} \int_{-\theta}^{\theta} e^{i\alpha\phi} d\phi &= \frac{1}{i\alpha} e^{i\alpha\phi} \Big|_{\phi=-\theta}^{\phi=\theta} \\ &= \frac{2i \sin \alpha \theta}{i\alpha} \end{aligned}$$

- note  $\pi q^2 a x = \frac{2\pi}{\lambda} a \frac{x}{D} \equiv k a \sin \theta$

$$\frac{x}{D} = \tan \theta \approx \theta \approx \sin \theta$$

Finally,

$$I(x) = I_0 \left[ \frac{\sin(k a \sin \theta)}{k a \sin \theta} \right]^2$$

Fraunhofer  
approximation  
for single-slit diffraction