

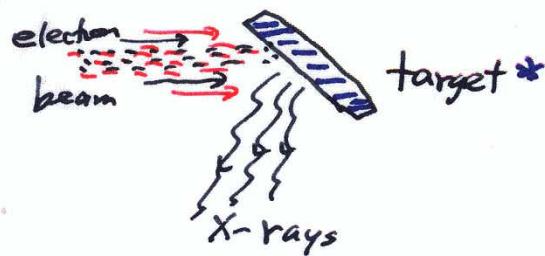
10. Properties of X-rays

C2/1

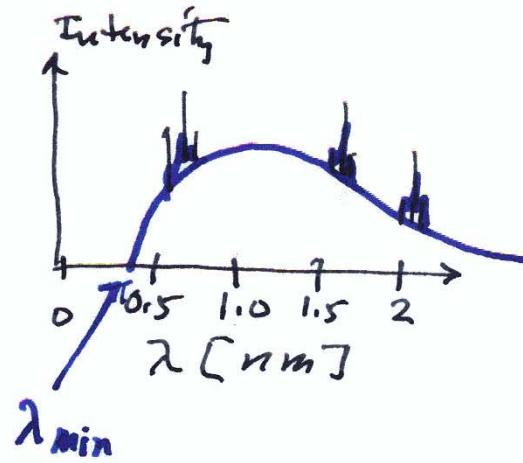
- Production of X-rays
- Interaction of X-rays with matter
(scattering and absorption)

Production of X-rays

X-ray tube



* the anode in the X-ray tube



The X-ray line spectrum

* Energy levels of inner electron shells of target atoms

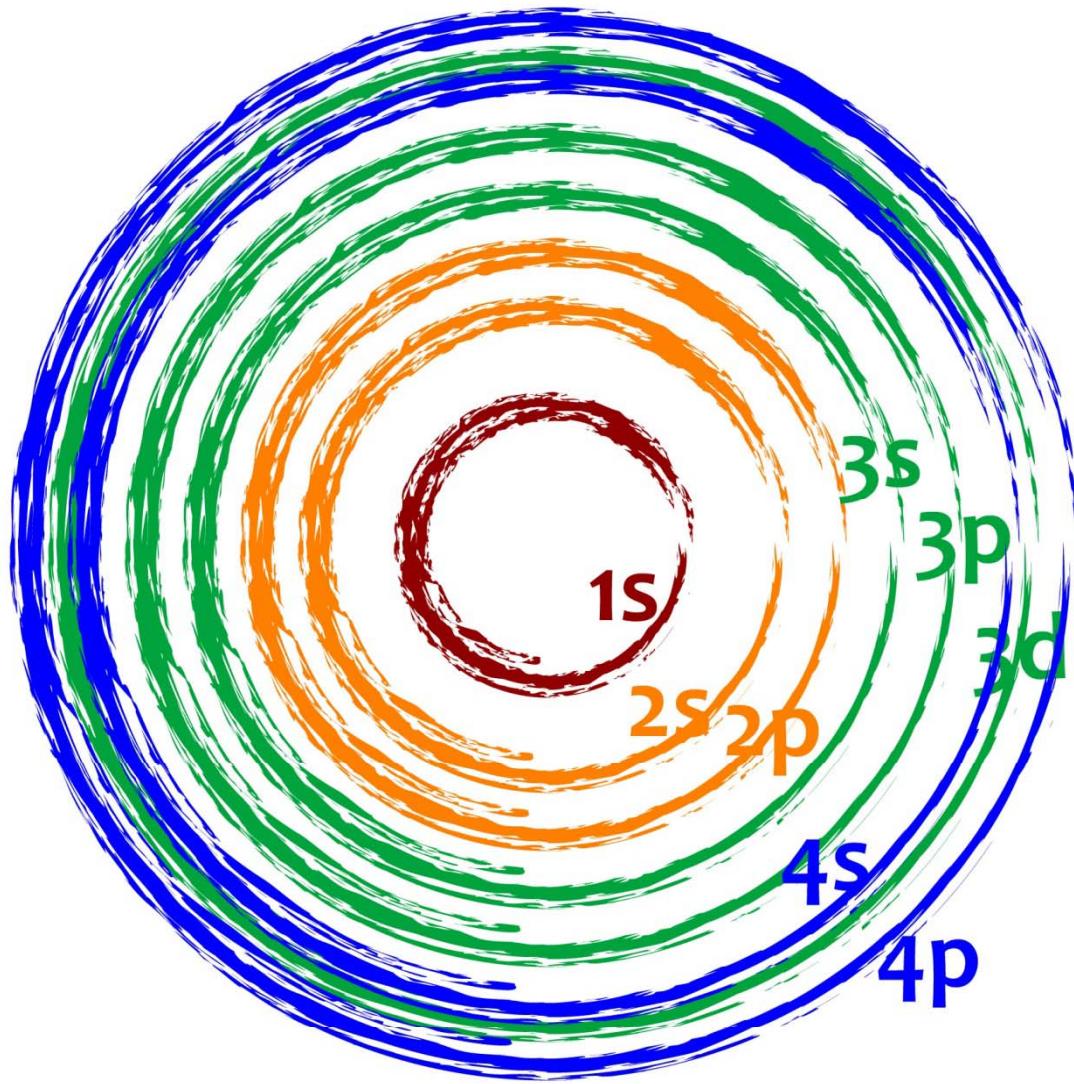
M $3s, 3p_{x_2}, 3p_{y_2}; 3d_{x_2} 3d_{z_2}$

L $2s; 2p_{x_2}, 2p_{y_2}$

K — $1s$

X-rays are produced when an electron falls from higher energy level E_1 to an empty lower energy level E_2 : $hf = \Delta E = E_1 - E_2$

Electron States in an Atom



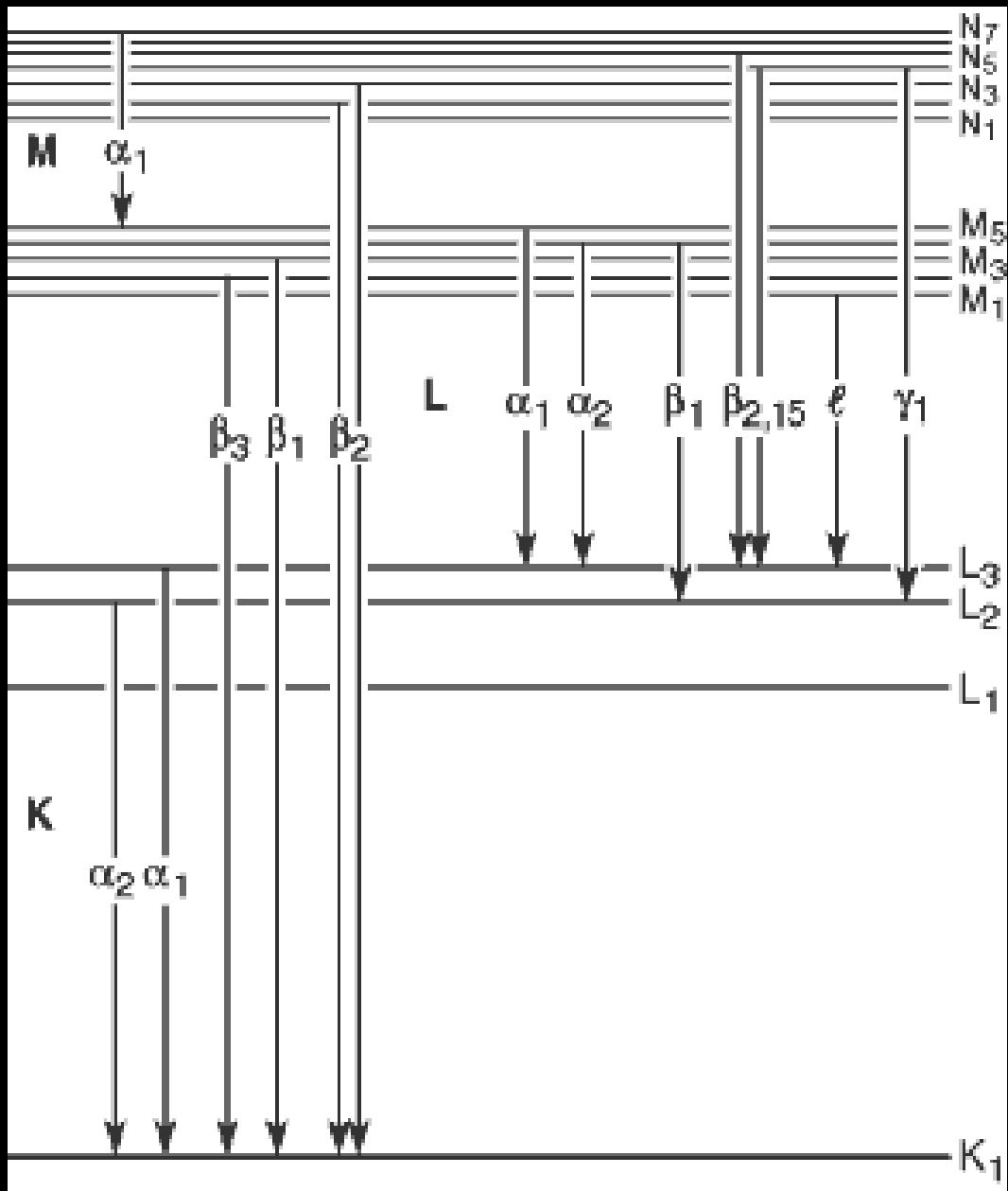
For example, Copper ($Z=29$):

$1s^2 \ 2s^2 \ 2p^6 \ 3s^2 \ 3p^6 \ 3d^{10} \ 4s^1$

Optical Spectrum = outer orbits

X-ray spectrum = inner orbits

Inner electron shells and X-ray emission transitions

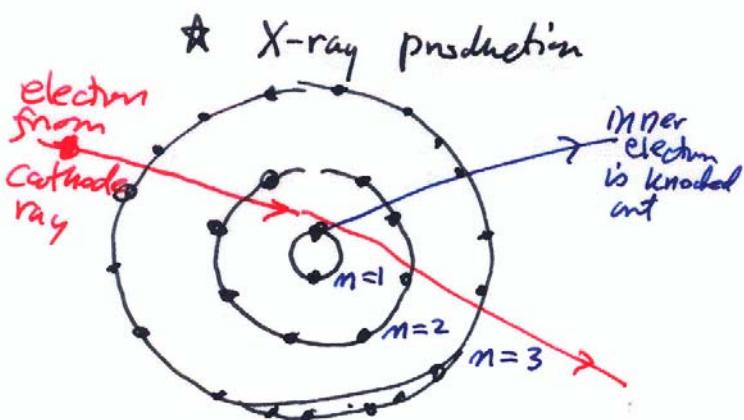


M shell
(3s, 3p $1/2$, 3p $3/2$
3d $3/2$, 3d $5/2$)

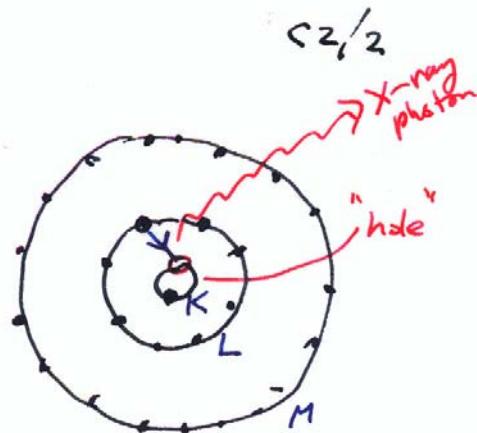
L shell
(2s, 2p $1/2$, 2p $3/2$)

K shell (1s)

X-ray Data Booklet
xdb.lbl.gov



(1) An inner shell electron is knocked out by energetic cathode-ray electron



(2) An outer electron falls into the hole in the inner shell and emits ~~an X-ray~~ a photon.

$$hf = E_L - E_K \quad (L \rightarrow K) \quad \text{or} \quad E_M - E_K \quad (M \rightarrow K)$$

$$\lambda = \frac{c}{f} = \frac{hc}{E_L - E_K} \quad (L \rightarrow K) \quad \text{or} \quad \frac{hc}{E_M - E_K} \quad (M \rightarrow K)$$

* Selection Rules

$$\Delta l = \pm 1 \quad \text{and} \quad \Delta j = 0, \pm 1$$

* Approximate energy levels

$$\text{Bohr model : } H = \frac{p^2}{2m} - \frac{ze^2}{r}$$

$$E_n = -\frac{Z^2 Ry}{n^2} \quad \text{where } Ry = 13.6 \text{ eV}$$

This neglects the outer electrons, which is a pretty good approximation. (*Why?*) More accurately,

$$E_n = -\frac{(Z-\delta)^2 Ry}{n^2} \quad \text{where } \delta \approx 1 \quad \text{"screening approximation"}$$

Examples

↳ from the X-ray Data Booklet

Carbon $Z=6$ $K\alpha_1 = 274 \text{ eV}$ $2p_{3/2} \rightarrow 1s$

$$\text{Bohr model } E_1 = -Z^2 R_y \text{ and } E_2 = -\frac{Z^2}{4} R_y$$

$$E_2 - E_1 = \frac{3}{4} Z^2 R_y = 367 \text{ eV}$$

$$\text{Screening model } E_1 = -(Z-1)^2 R_y \text{ and } E_2 = -\frac{(Z-1)^2}{4} R_y$$

$$E_2 - E_1 = \frac{3}{4}(Z-1)^2 R_y = 255 \text{ eV}$$

not very accurate

Uranium $Z=92$ $K\alpha_1 = 98.4 \text{ keV}$ $2p_{3/2} \rightarrow 1s$

$$\text{Bohr model } E_2 - E_1 = \frac{3}{4} Z^2 R_y = 86.3 \text{ keV}$$

But the inner electrons of uranium
are relativistic; $\frac{v}{c} \approx Z\alpha = \frac{92}{137} = 0.67$
so the Bohr model is not accurate.

Iron $Z=26$ $K\alpha_1 = 6.40 \text{ keV}$ $2p_{3/2} \rightarrow 1s$

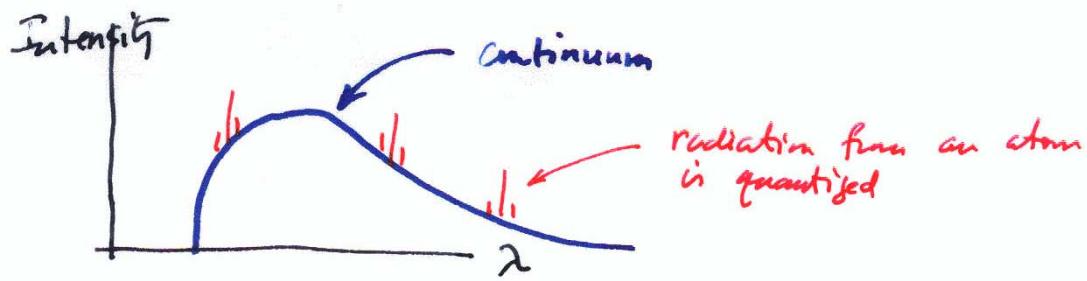
$$\text{Bohr model } E_2 - E_1 = \frac{3}{4} Z^2 R_y = 6.90 \text{ keV}$$

$$\text{Screening model } E_2 - E_1 = \frac{3}{4}(Z-1)^2 R_y = 6.38 \text{ keV}$$

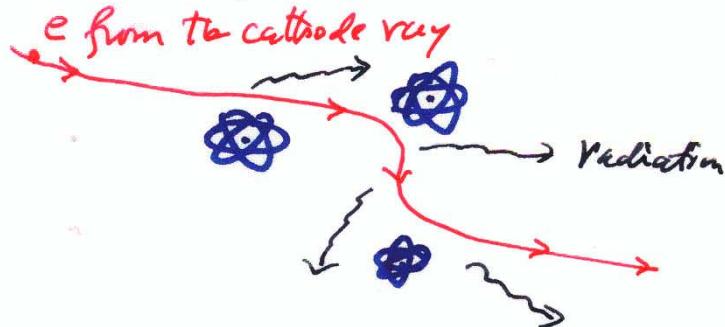
for a wide range of intermediate Z values,
the Bohr model and screening model are quite accurate

The Continuum Spectrum

C2/4



In the classical theory of electromagnetic radiation, a charged particle radiates when it accelerates



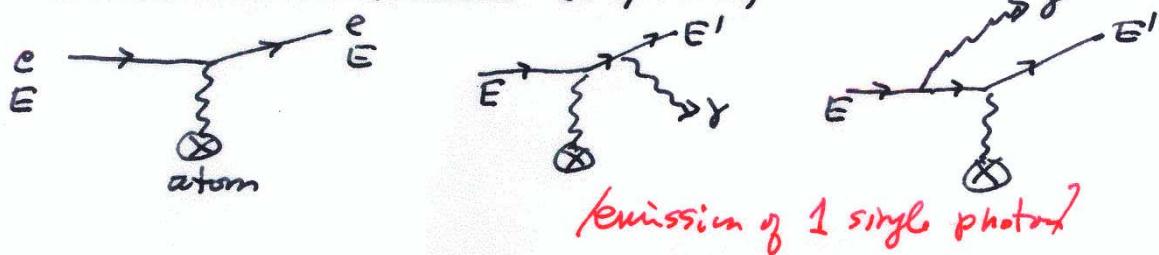
Classical e.m. Radiation

- continuous spectrum

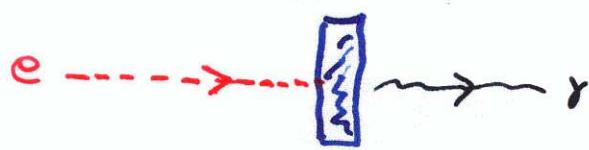
- Larmor's Formula: $P_{\text{radiated}} = \frac{e^2 a^2}{6\pi\epsilon_0 c^3}$

Instantaneous power; $a = \text{acceleration}$

Quantum electrodynamics (Reynman)

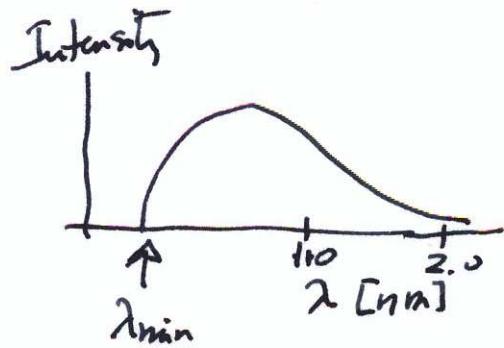


The X-ray spectrum is continuous, but radiation is probabilistic.



c2/5

As the electron slows down in the anode, about 1% of its kinetic energy is transferred to X-ray radiation. ^(on average) The rest goes to heat.

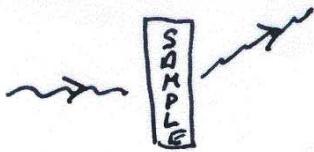


$$\lambda_{\min} = \frac{hc}{E}$$

(why?)

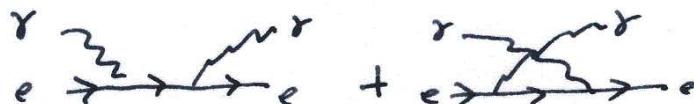
Interaction of X-rays with matter

Scattering



- Thomson scattering

$$\gamma + e \rightarrow \gamma + e \quad \text{with} \quad E_\gamma \ll mc^2$$



$$\frac{d\sigma}{d\Omega} = \frac{1}{2} r_e^2 (1 + \cos^2 \theta) \quad \text{where} \quad r_e = \frac{e^2}{4\pi \epsilon_0 m c^2}$$

$$\sigma_T = \int \frac{d\sigma}{d\Omega} \sin \theta d\theta d\phi = \frac{8\pi}{3} r_e^2$$

(J.J. Thomson, Theory
published ~1906)

$$r_e = 2.8 \times 10^{-15} \text{ m}$$

$$\sigma_T = 6.6 \times 10^{-29} \text{ m}^2 = 0.66 \text{ barns}$$

- Compton scattering ($E_\gamma \gtrsim mc^2$)

The X-ray loses some energy

$$\frac{E'_\gamma}{E_\gamma} = \frac{mc^2}{mc^2 + E(1 - \cos \theta)}$$

$$\lambda' - \lambda = \frac{h}{mc} (1 - \cos \theta)$$

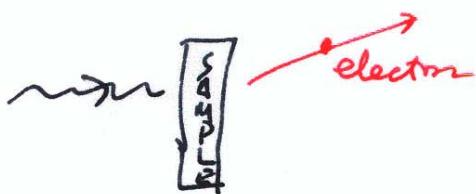
$$\underline{E = hf = \frac{hc}{\lambda}}$$

(Arthur Compton, 1923)

Interactions of X-rays with matter

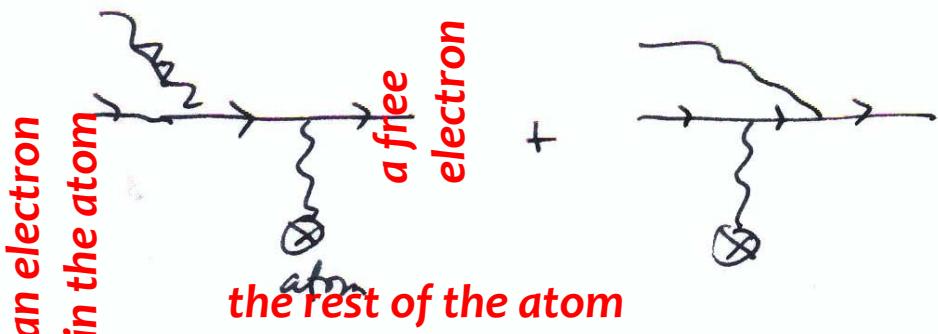
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Absorption



- Photoelectric effect

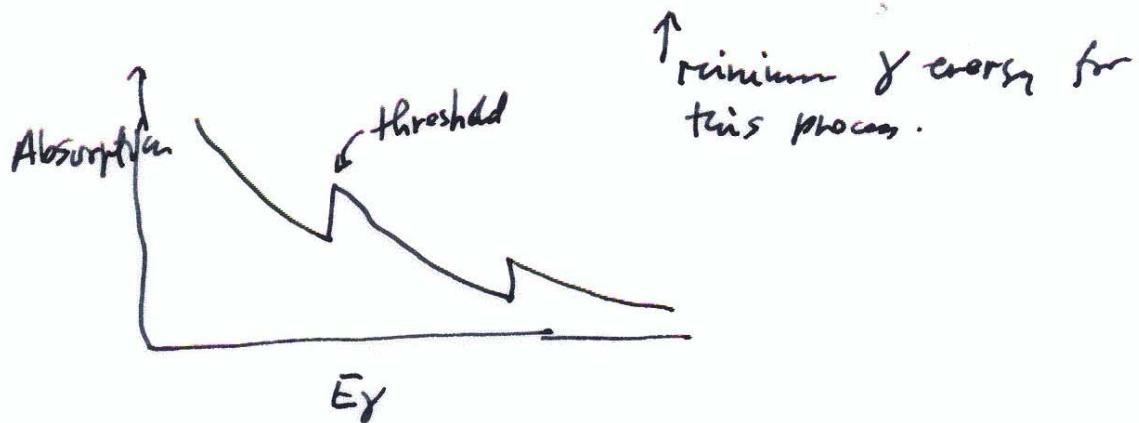
... on an atomic scale



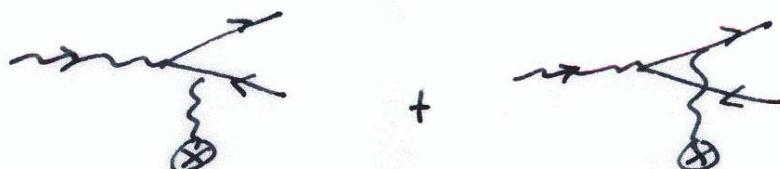
$$hf - B = \frac{1}{2}mv^2$$

(neglect recoil energy of the atom)

$$hf = B + \frac{1}{2}mv^2 \geq B$$

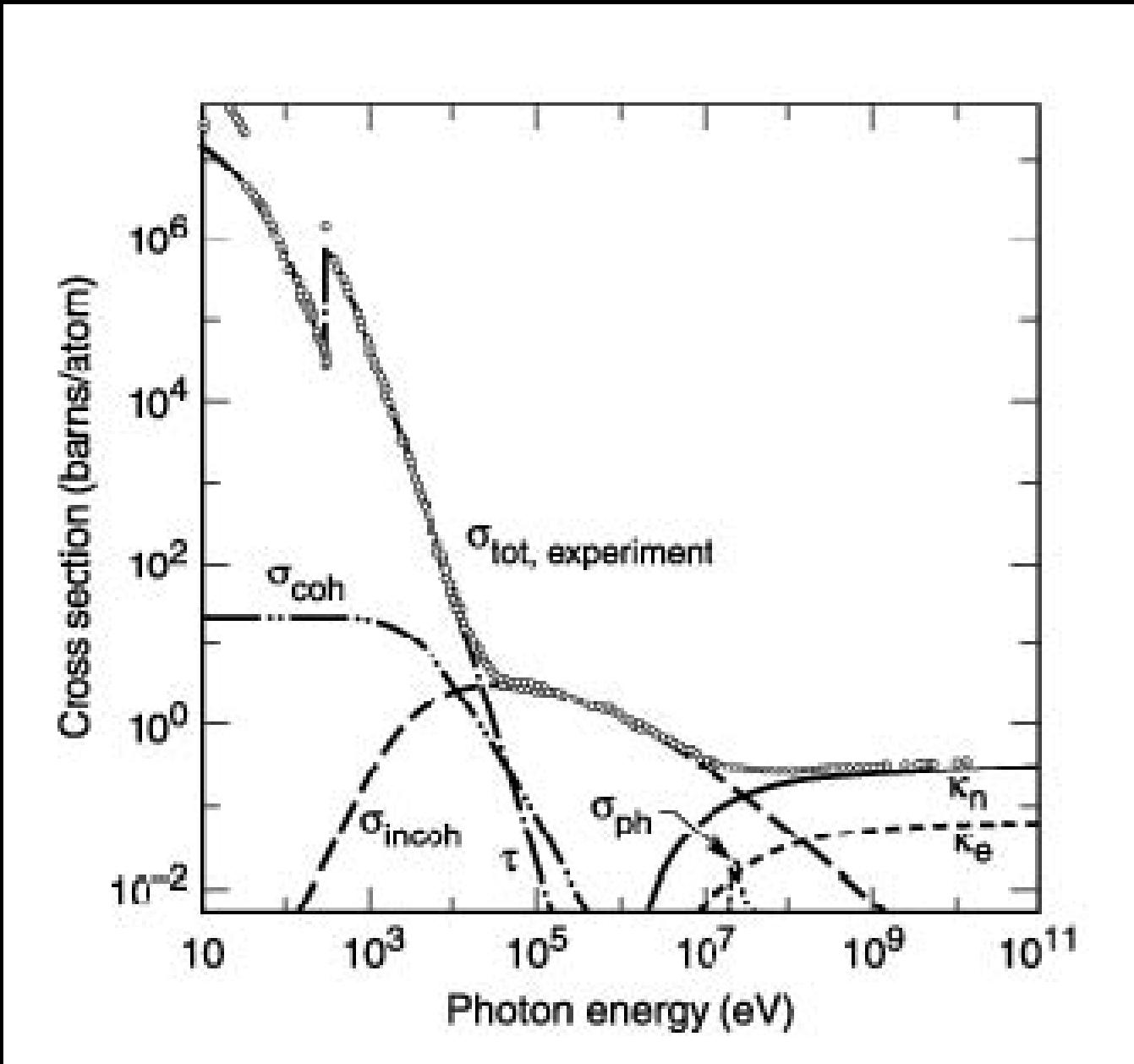


- Pair Production: $\gamma + \text{atom} \rightarrow e^- e^+ + \text{atom}$



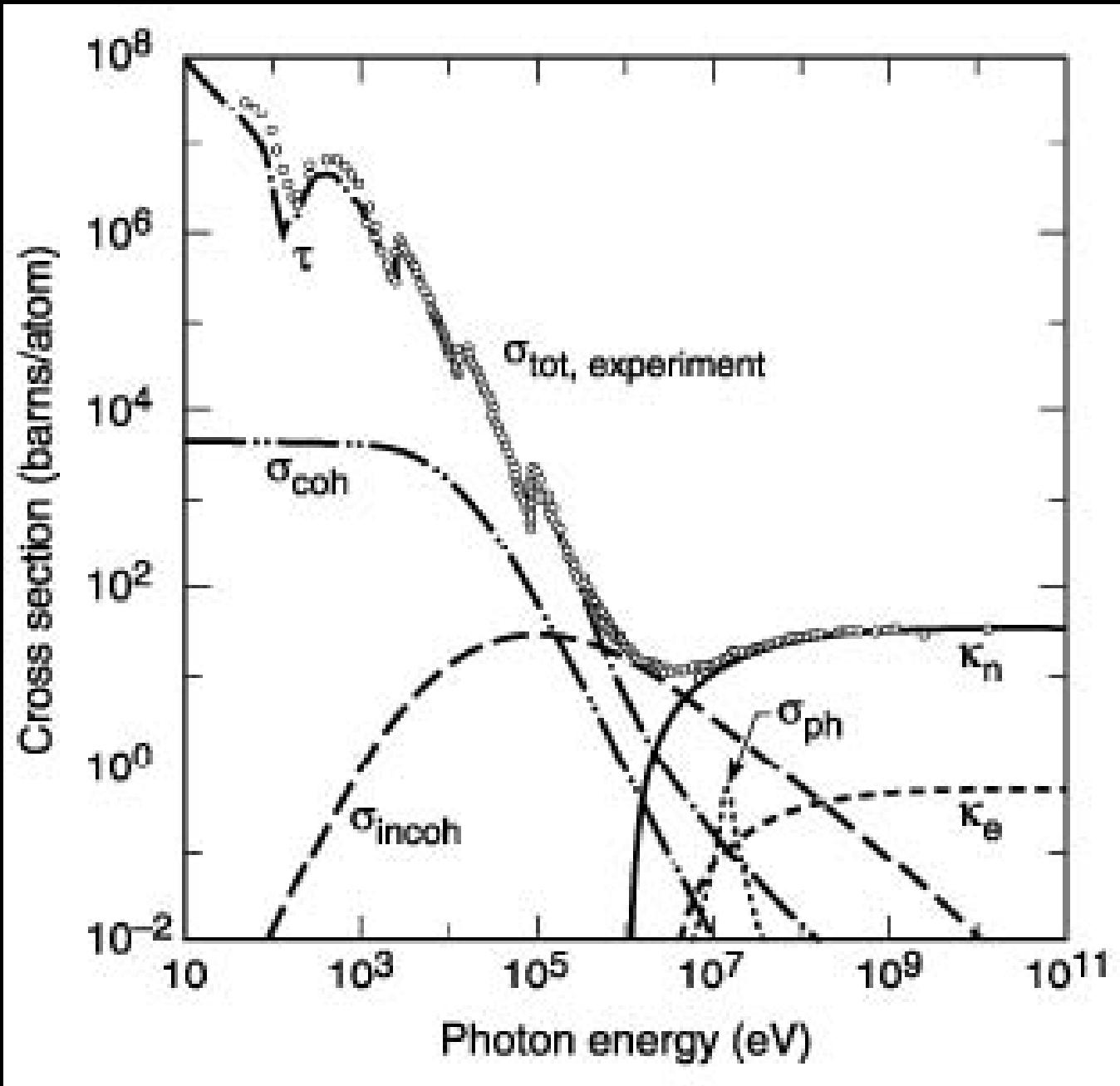
Can occur if $E_\gamma > 2mc^2 = 1.02 \text{ MeV}$

Carbon cross sections



X-ray Data Booklet
xdb.lbl.gov

Lead cross sections



X-ray Data Booklet
xdb.lbl.gov