Part I: Thin Converging Lens

This experiment is a classic exercise in geometric optics. The goal is to measure the radius of curvature and focal length of a single converging lens from which you can calculate the index of refraction $n$. We shall explicitly consider the errors that accompany any measurement and how errors are analyzed to yield a quantitative estimate of uncertainty. This includes quantities derived from measurements, in this case the index of refraction.

In the procedures for this lab, you are asked to estimate the uncertainties. Please see Appendix I and II for reference material and some relevant equations. The questions, labeled Q1, Q2, … should be explicitly addressed in your write-up in the Analysis & Discussion section.

Procedures

A. By definition, the focal length $f$ of a lens is the image distance from the lens center for an infinitely distance object. To obtain a rough estimate of $f$, project an image of the trees outside the lab onto the white paper near the door. Use a 3” diameter lens labeled A, B, C, or D.

Q1. Why do the trees appear upside down?

B. Use a spherometer to measure the radius of curvature of both surfaces of your lens. See Appendix I. You begin by finding the “zero” position, $x_0$, using a scratch-free spot on your bench (a good approximation to a flat surface). Then perform the measurement with your lens in place, $x_1$; the distance $h$ is then $|x_0 - x_1|$. Repeat the measurements a few times to obtain an estimate of the spherometer’s precision.

Rulers, calipers, graph paper, and markers are available.

Q2. Having estimated measurement uncertainties $\delta x_0$ and $\delta x_1$, write an expression for $\sigma_h$ and evaluate it using your data.

C. Arrange an object (the T on the lamp window) and screen with a separation greater than $4f$ on the optical rail. Locate the lens position which gives a sharp image on the screen. Record the object and image distances measured from the center of the lens including uncertainties. Use the thin lens equation to calculate $f$ and $\sigma_f$. Repeat this for 4 positions of the screen increasing the object-screen separation in units of about 2 cm. Find the best value for the focal length.

Q3. How does the best value for $f$ compare to your original rough estimate?

D. Insert a variable iris before/after the lens. Observe the image as the aperture size is changed. Specifically note whether it affects your ability to focus the image.

Q4. What is the meaning of the term “depth of field”?

E. Place the light source a distance less than $f$ from the lens. Try to position the screen to bring the object into focus.

Q5. How do you explain your observations?

F. Calculate the index of refraction (including uncertainty) for the glass of your lens using the Lensmaker’s Equation. Compare with known values for $n$. 


Appendix I: Lens Equations

Thin Lens Equation

\[ \frac{1}{s} + \frac{1}{s'} = \frac{1}{f} \]

Lensmaker’s Equation

\[ \frac{1}{f} = (n - 1) \left( \frac{1}{R_1} - \frac{1}{R_2} \right) \]

assuming that \( n_{\text{air}} = 1.000 \).

Spherometer Equation

\[ R = \frac{b^2}{2h} + \frac{h}{2} \]

From Pythagoras:

\[ R^2 = (R - h)^2 + b^2 \]

\[ R^2 = R^2 - 2Rh + h^2 + b^2 \]

\[ R = \frac{b^2}{2h} + \frac{h}{2} \]
Appendix II  Error Analysis

Random fluctuations in the measurement process lead to a Gaussian distribution about the true value. This distribution gives us a parameter, $\sigma$, called the “standard deviation”. (Systematic errors lead to a non-Gaussian distribution.) Essentially, if many measurements $x_i$ are taken, 68% of the data points lie within $\bar{x} \pm \sigma$, where $\bar{x}$ is the mean value of $x$.

Now, suppose an arbitrary function $f(x,y)$ depends on the variables $x$ and $y$, assumed to be independent of each other. How do we compute the uncertainty in $f$, $\sigma_f$, given $\sigma_y$ and $\sigma_x$? Under the assumption that the uncertainties are small compared to the absolute value of the quantities in question

$$\sigma_f = \sqrt{\left(\frac{\partial f}{\partial x}\right)^2 \sigma_x^2 + \left(\frac{\partial f}{\partial y}\right)^2 \sigma_y^2}$$

For errors that are much smaller than the measured values, specific functions yield:

- $f = ax + by$
  $$\sigma_f = \sqrt{a^2 \sigma_x^2 + b^2 \sigma_y^2}$$

- $f = cxy$
  $$\frac{\sigma_f}{f} = \sqrt{\left(\frac{\sigma_x}{x}\right)^2 + \left(\frac{\sigma_y}{y}\right)^2}$$

- $f = cx^a y^b$
  $$\frac{\sigma_f}{f} = \sqrt{\left(a \frac{\sigma_x}{x}\right)^2 + \left(b \sigma_y \frac{y}{y}\right)^2}$$

- $f = ce^{bx}$
  $$\frac{\sigma_f}{f} = b \sigma_x$$

- $f = ce^{ax}$
  $$\frac{\sigma_f}{f} = (b \ln a) \sigma_x$$

When we make $N$ measurements of the same quantity $x$, each with an uncertainty $\delta_x$, we expect that after averaging the measurement will have uncertainty smaller than $\delta_x$. In fact, the value of $\sigma_x$ varies as $1/\sqrt{N}$ when $N$ is large.
Part II: Thin Divergent Lens

For a divergent lens, all principles and conventions used for a convergent lens will apply. The key difference is that a divergent lens cannot by itself form a real image of a real object. Hence, in this experiment we will measure \( f \) using a virtual object. The virtual object and real image are on the same side of the lens. You will measure the radius of curvature and focal length, then calculate the index of refraction of the glass. Be sure to consider error propagation in your final results.

You will use white light and laser light sources. The narrow emission spectrum of the laser eliminates a source of blurring called chromatic aberration. This will allow for more precise measurements. The laser also produces a well-defined beam of light rays. You will measure the degree of collimation, or spreading, of the He-Ne laser beam. The beam can be expanded or contracted with a two-lens telescope.

Procedure:

A. Use a spherometer to measure the radius of curvature of a divergent lens.

B. Use a convergent lens \( L_1 \) to form a sharp image \( i_1 \) of your object on a screen using the lamp as a source. Next, place a divergent lens \( L_2 \) between \( L_1 \) and \( i_1 \) as shown below. Measure the distances to \( i_1 \) and \( i_2 \) to calculate \( f \) for the divergent lens. Repeat this for 3 positions of \( i_2 \) by changing the lens-screen separation in units of about 1 cm. Find the best value for the focal length using the thin-lens equation.

Q1. Make a quantitative comparison of the \( L_2 \) focal length obtained by the two different methods. Do they agree within experimental error? How could you improve either measurement?

C. Calculate the index of refraction (including uncertainty) for the glass of your lens using the lensmaker’s equation. Compare this with the value found in Part I.
D. Switch to the laser source. If the beam from the laser is not perfectly collimated, the diverging rays must spread over some angle $\theta$. Aim the beam from the He-Ne laser onto a distant wall and measure the radius of the maximum spot size that can be discerned.

![Diagram of laser with angle $2\theta$]

**Q2.** Is the spot of uniform brightness? Calculate $2\theta$ in degrees and radians.

E. Using the converging and diverging lenses and the equation for a Galilean telescope make a laser beam expander/reducer. Measure the beam diameter before and at several distances beyond the lenses.

**Q3.** Does the beam diverge? Estimate its divergence angle.

**Q4.** What is the magnification of the beam? Show whether this is consistent with the two-lens equation.

Before leaving the laboratory, make rough estimates of all quantities that need to be calculated or included in your report.

### Appendix II: Misc Equations

The image position for a two lens combination is

$$ s_i = \frac{f_2 d - \frac{f_2 s_o f_1}{s_o - f_1}}{d - f_2 - \frac{s_o f_1}{s_o - f_1}} $$

where $s_o$ is the position of the object (before either lens), and $d$ is the distance between the lenses. If $s_o = \infty$, and $d = f_1 + f_2$, then $s_i = \infty$. This is a Galilean telescope.