## Physics 472 - Spring 2009

## Homework \#9, due Friday, March 27

(Point values are in parentheses.)

1. [7] The isotropic 2-dimensional harmonic oscillator is easily solved by writing the Hamiltonian as a sum of $x$ and $y$ Hamiltonians:

$$
\hat{H}^{0}=\frac{\hat{P}_{x}^{2}+\hat{P}_{y}^{2}}{2 m}+\frac{1}{2} m \omega^{2}\left(x^{2}+y^{2}\right)=\hat{H}_{x}^{0}+\hat{H}_{y}^{0} \text { with }\left\lfloor\hat{H}_{x}^{0}, \hat{H}_{y}^{0}\right\rfloor=0 .
$$

Simultaneous eigenstates of $\hat{H}_{x}^{0}$ and $\hat{H}_{y}^{0}$ obey $\hat{H}^{0}\left|n_{x}, n_{y}\right\rangle=\left(n_{x}+n_{y}+1\right) \hbar \omega\left|n_{x}, n_{y}\right\rangle$.
a) Consider the perturbation $\lambda \hat{H}^{\prime}=\lambda m \omega^{2} x y$. Calculate the first and second order energy shifts of the ground state. In class we used $\hat{X}=\sqrt{\frac{\hbar}{2 m \omega}}\left(\hat{a}_{x}^{+}+\hat{a}_{x}\right)$ to evaluate the matrix elements: $\left\langle n_{x}^{\prime}\right| x\left|n_{x}\right\rangle=\sqrt{\frac{\hbar}{2 m \omega}}\left(\sqrt{n_{x}} \delta_{n_{x}^{\prime}, n_{x}-1}+\sqrt{n_{x}+1} \delta_{n_{x}^{\prime}, n_{x}+1}\right)$. The same holds for $\hat{Y}=\sqrt{\frac{\hbar}{2 m \omega}}\left(\hat{a}_{y}^{+}+\hat{a}_{y}\right)$. To evaluate the matrix elements of $\hat{H}^{\prime}$ in the $\left|n_{x} n_{y}\right\rangle$ tensor product states, use:

$$
\left\langle n_{x}^{\prime} n_{y}^{\prime}\right| x y\left|n_{x} n_{y}\right\rangle=\left\langle n_{x}^{\prime}\right| x\left|n_{x}\right\rangle\left\langle n_{y}^{\prime}\right| y\left|n_{y}\right\rangle
$$

b) Use degenerate P.T. to calculate the first-order energy shifts of the first excited states, as well as the "correct" linear combinations of those two states that diagonalize $\hat{H}^{\prime}$.
c) The full Hamiltonian, $\hat{H}=\hat{H}^{0}+\lambda \hat{H}^{\prime}$, is exactly solvable if you make the coordinate transformation $u=(x+y) / \sqrt{2}, v=(x-y) / \sqrt{2}$. Express $\hat{H}$ in terms of $u$, $v$, and their conjugate momenta $P_{u}$, and $P_{v}$. You should find that the harmonic oscillator in the " $u$ " direction has a higher frequency than before, while in the " $v$ " direction the frequency is lower. Calculate the exact energies of the new basis states $\left|n_{u}, n_{v}\right\rangle$. For the ground state, expand the energy to second order in $\lambda$. For the next two higher states, expand the energies to first order in $\lambda$. Compare your results with those you obtained in parts (a) and (b).
2. [7] Consider an electron in a 3-dimensional isotropic harmonic oscillator potential, in the presence of a uniform magnetic field $\vec{B}=B_{\text {ext }} \hat{k}$. The full Hamiltonian for the system is:

$$
\hat{H}=\frac{\hat{P}_{x}^{2}+\hat{P}_{y}^{2}+\hat{P}_{z}^{2}}{2 m}+\frac{1}{2} m \omega^{2}\left(x^{2}+y^{2}+z^{2}\right)+\frac{e B_{e x t}}{2 m}\left(L_{z}+2 S_{z}\right)
$$

This problem is exactly solvable, but we'll use the machinery of perturbation theory to get organized. Treat the first two terms of $H$ as $H^{0}$, and the Zeeman term as $\lambda H^{\prime}$. The eigenstates of $H^{0}$ satisfy: $\hat{H}^{0}\left|n_{x}, n_{y}, n_{z}, m_{s}\right\rangle=\left(n_{x}+n_{y}+n_{z}+\frac{3}{2}\right) \hbar \omega\left|n_{x}, n_{y}, n_{z}, m_{s}\right\rangle$, where $m_{s}= \pm \frac{1}{2}$, and $S_{z}\left|n_{x}, n_{y}, n_{z}, m_{s}\right\rangle=\hbar m_{s}\left|n_{x}, n_{y}, n_{z}, m_{s}\right\rangle$. (I am putting the space and spin quantum numbers together inside the same ket to avoid using the cumbersome tensor product notation.)
a) Express $\hat{L}_{z}$ in terms of the harmonic oscillator raising and lowering operators. Hint: you should get $\hat{L}_{z}=i \hbar\left(a_{x} a_{y}^{+}-a_{x}^{+} a_{y}\right)$.
b) The ground state of $H^{0}$ is 2-fold degenerate, due to spin. But since the two $\left|0,0,0, m_{s}\right\rangle$ states are already eigenstates of $H^{\prime}$, you can use standard first-order perturbation theory to calculate the energy shifts due to the magnetic field. Express your answers in terms of $\mu_{B} B_{\text {ext }}$ and $m_{s}$.
c) The first excited state of $H^{0}$ is six-fold degenerate ( 3 spatial states $\times 2$ spin states). Calculate the linear combinations of states that diagonalize $H^{\prime}$. To help you keep track of what you are doing, here are some suggestions. First, since all your states are eigenstates of $S_{z}$, leave spin out of the problem until the end; then you only have to diagonalize a $3 \times 3$ matrix rather than a $6 \times 6$ matrix. Since the original basis states $\left|n_{x}, n_{y}, n_{z}, m_{s}\right\rangle$ are not eigenstates of $\hat{L}_{z}$, you need to find linear combinations of them that are. Label the new states this way: $\left|n, l, m_{l}, m_{s}\right\rangle$, where $n=n_{x}+n_{y}+n_{z}$. You don't need to know $l$ to do this problem - you just need $m_{l}$. But you can probably guess what $l$ is once you know what $m_{l}$ is for the three states. Finally, when you have found the states that diagonalize $H^{\prime}$, calculate the Zeeman energy shifts of those states. How many distinct energies are there? Make a plot of energy vs. $\mu_{B} B_{\text {ext }}$ for all the states.
d) The second excited state of $H^{0}$ is twelve-fold degenerate ( 6 spatial states $\times 2$ spin states). Forget about spin altogether so you don't get lost. Construct the $6 \times 6$ matrix representation of $\hat{L}_{z}$. If you choose the order of your 6 states judiciously, your $6 \times 6$ matrix should break up into a $2 \times 2$ block, a $3 \times 3$ block, and a $1 \times 1$ block. Calculate the eigenvalues of $\hat{L}_{z}$ and their degeneracies. Guess what the values of $l$ are for this six-dimensional subspace. Don't bother to calculate the 12 Zeeman energies - I know you could do it if you had to!
3. [6] Griffiths problem 6.37. Follow the same strategy you used to solve Griffiths problem 6.36. Use symmetries to figure out which matrix elements of the form $\left\langle n, l^{\prime}, m_{l}{ }^{\prime}\right| z\left|n, l, m_{l}\right\rangle$ are zero. Rotational symmetry, $\left\{\hat{z}, \hat{L}_{z}\right\rfloor=0$, implies $\left(m_{l}{ }^{\prime}-m_{l}\right)\left\langle n, l^{\prime}, m_{l}{ }^{\prime}\right| z\left|n, l, m_{l}\right\rangle=0$. The parity transformation, $\hat{\Pi} \hat{z} \hat{\Pi}=-\hat{z}$, implies $(-1)^{l^{\prime}+l}\left\langle n, l^{\prime}, m_{l}^{\prime}\right| z\left|n, l, m_{l}\right\rangle=-\left\langle n, l^{\prime}, m_{l}^{\prime}\right| z\left|n, l, m_{l}\right\rangle$.

Use these same symmetries to do Griffiths problem 6.37. First, show which elements of the $9 \times 9$ matrix are zero. Then calculate the first non-zero matrix element, $\langle 3,0,0| z|3,1,0\rangle$, using Tables 4.3 and 4.7 in Griffiths. Use Mathematica to do the radial integration. You can take the values of the other nonzero matrix elements from Griffiths. Construct the $9 \times 9$ matrix representation of $\hat{z}$. If you choose the order of the 9 states carefully, then the matrix should break into a $3 \times 3$ block, two $2 \times 2$ blocks, and two trivial $1 \times 1$ blocks. Calculate the eigenvalues and their degeneracies. Don't forget to multiply the eigenvalues by $e E_{\text {ext }}$ to get the energies.

