

LECTURE # 29

Note Title

4/1/2009

- ~~Quantum Hall effect~~
- Graphene
- ~~Carbon Nanotubes~~
- Excitons and exciton-polaritons
- Mesoscopic physics and single electronics
- Spectroscopy of quantum dots and quantum wires
- ~~Spintronics in metals/semiconductors~~
- Neutron and electron scattering in solids
- Solid state devices for quantum information processing
- Magnetic properties of solids
- Superconductivity

INTRODUCTION
TO DENSITY
FUNCTIONAL THEORY

HARTREE - FOCK

$$E^{HF}[\{\psi_i, \psi_i^*\}] \rightarrow \text{FUNCTIONAL}$$

$$\frac{\delta E^{HF}}{\delta \psi_i^*} = \mu \psi_i \Rightarrow \text{EFFECTIVE 1 PARTICLE SCHR EQ}$$

HARTREE + EXCHANGE

DENSITY FUNCTIONAL :

$$E[n(\mathbf{r})] \quad \text{TOT ENERGY AS FUNCTIONAL OF}$$

$$n(\mathbf{r}) \quad (\text{ELECTRONIC DENSITY}) \quad \frac{\delta E[n(\mathbf{r})]}{\delta n} \rightarrow \text{VARIATION ON } n(\mathbf{r})$$

N ELECTRONS

$$H^{\text{TOT}} = \underbrace{\sum_i \frac{p_i^2}{2m}}_T + V^{\text{ION}} + \underbrace{\sum_{i,j} \frac{e^2}{|r_i - r_j|}}_{H^{(2)}}$$

ALWAYS THE SAME FOR ATOMS, MOLECULES, CRYSTALS

• ALL PROPERTIES OF N -ELECTRON SYSTEMS ARE UNIQUELY DETERMINED BY $V^{\text{ION}} = v(r)$

• ANY OBSERVABLE $O[v(r)]$ IS A FUNCTIONAL OF $v(r)$

• IN PARTICULAR, TOTAL DENSITY

$$\hat{\rho} = \sum_i \delta(r - r_i) \quad \langle \Psi^{\text{MANY-BODY}} | \hat{\rho} | \Psi^{\text{MANY-BODY}} \rangle = n(r)$$

$m[\psi(u)]$: DENSITY IS A FUNCTIONAL
OF $\psi(u)$

DENSITY FUNCTIONAL THEORY BASED ON

(I) Hohenberg - Kohn Theorem:

$\psi(u)$ IS UNIQUELY DETERMINED BY
 $m(u)$ i.e. $\psi[m(u)]$

REDUCTIO AD ABSURDUM:

IF STATEMENT OF THEOREM IS NOT TRUE THEN

$\exists \psi(u)$ AND $\psi'(u)$ THAT WILL GIVE THE
 \downarrow $\psi(u) \neq \psi'(u)$ \downarrow SAME $m(u)$
 ψ, E_0 ψ', E_0'

THEN

$$E_0 = \langle \psi | T + v + H^{(2)} | \psi \rangle = \langle \psi' | T + v + H^{(2)} | \psi' \rangle = \overset{+v' - v}{=}$$

$$\langle \psi' | T + v' + H^2 + v - v' | \psi' \rangle = E_0' + \langle \psi' | v - v' | \psi' \rangle$$

$$E_0 - E_0' < \int dr \, m(u) (v(u) - v'(u)) \quad (a)$$

$$E_0' = \langle \psi' | T + v' + H^{(2)} | \psi' \rangle < \langle \psi | T + v' + H^{(2)} | \psi \rangle \overset{+v' - v}$$

$$\Rightarrow E_0 - E_0' > \int du \, m(u) (v(u) - v'(u)) \quad (b)$$

(a) + (b) ARE TRUE ONLY IF $v(u) = v'(u)$

ABSURDUM

$$m(u) = \langle \psi' | \hat{p} | \psi' \rangle = \langle \psi | p | \psi \rangle$$

$$E[V(u)] \rightarrow E[V[m]] \Rightarrow$$

$$E[m(u)] \Rightarrow \begin{array}{l} \text{GROUND STATE} \\ \text{MANY BODY ENERGY} \end{array}$$

UNIQUELY DETERMINED BY $m(u)$

$$E[m(u)] = \underbrace{\int du v(u) m(u)} + \underbrace{F[m(u)]}_{\substack{\text{KINETIC TERM} + \\ \text{COULOMB}}}$$

VARIATIONAL METHOD

$$\delta(E[m] - \mu \int du m(u)) = 0$$

$$\frac{\delta E[M]}{\delta n} = \mu$$

DENSITY FUNCTIONAL
EQUATION

$$\frac{\delta F[M]}{\delta n} + v(u) = \mu \quad \leftarrow$$

HOW DO WE WRITE $F[u]$?

$F[M]$ FOR A NON-INTERACTING SYSTEM

$H^{(2)} = 0$ NO $e-e$ COULOMB

$$E_S[M] = \int du v(u) n(u) + T_S[M(u)]$$

\downarrow NON-INTERACTING

$$\frac{\delta T_s [m]}{\delta n} + v(u) = \mu$$

(A) FROM DENSITY FUNCTIONAL THEORY

EQUIVALENT TO:

$$n(u) = \sum_j \theta(\mu - \epsilon_j) |\phi_j(u)|^2$$

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + v(u) \right] \phi_j(u) = \epsilon_j \phi_j(u)$$

(B) FROM SCHRÖD. EQ.

THIS LINK BETWEEN SE AND

DFT IS USED IN THE

(H) Kohn - Sham EQUATIONS

$$E[m(r)] = T_S^{(A)}[m] + \int dr v(r) m(r) + \frac{1}{2} \int dr dr' \frac{m(r) m(r')}{|r - r'|} + E_{xc}^{(C)}[m]$$

(A) $T_S[m]$ KINETIC ENERGY OF NON-INTERACTING SYSTEM WITH SAME DENSITY $m(r)$

(B) HARTREE

(C) EXCHANGE - CORRELATION

$$\frac{\delta E}{\delta n} = \mu$$

$$\frac{\delta T_s}{\delta n} + v(r) + \int dr' \frac{m(r')}{|n - n'|} + \frac{\delta E_{xc}[n]}{\delta n} = \mu$$

$v^\#(r)$ (pointing down to the integral term)
 $v^{xc}(r)$ (pointing up to the exchange-correlation term)

EQUIVALENT TO

$$\left[-\frac{\hbar^2 \nabla^2}{2m} + v^{\text{EFF}}(r) \right] \phi_J(r) = \epsilon_J \phi_J(r)$$

$$v^{\text{EFF}} = v(r) + v^\#(r) + v^{xc}(r)$$

$$n(r) = \sum_J \theta(\mu - \epsilon_J) |\phi_J(r)|^2$$

KOHN
SHAM
ORBITALS

• SOLVED SELF-CONSISTENTLY

• BEYOND HF → DEPENDS ON

WHAT IS IN $\sqrt{x_c}(u)$

$\sqrt{x_c}(u) \rightarrow E_{xc}[M]$

$E_{xc}[M]$ TAKEN OFTEN FROM

HOMOGENEOUS ELECTRON GAS

GREEN'S FUNCTIONS

QUANTUM MONTECARLO

GOOD GUESS

\bar{E}_{xc}

LOCAL

DENSITY

APPROX

LDA

DFT-LDA

$$E_{xc}[n] = \int d\mathbf{r} n(\mathbf{r}) \overbrace{E_{xc}(n(\mathbf{r}))}^{\text{LOCALITY}}$$

YOU COULD USE HOMOEGAS + HF FOR THE

EXCHANGE

↓ KINETIC

↗ EXCHANGE

$$\frac{E_{HF}}{N} = \left(\frac{2.21}{(r_s/a_0)^2} - \frac{0.916}{(r_s/a_0)} \right) \underbrace{\hspace{10em}}_{E_x(n)}$$

$$\frac{4\pi}{3} n_0^3 = \frac{1}{n}$$

$$\varepsilon_x[m] = -\alpha m^{1/3}$$

$$E^{xc}[m] = -\alpha \int du m(u) m^{1/3}(u)$$

$$v^x(u) \sim -\alpha \frac{4}{3} m^{1/3}(u)$$