## Geometrical Optics

## Refraction at Spherical Surface



Optical Path Length (OPL) from S to P is

$$
O P L=n_{1} l_{0}+n_{2} l_{i}
$$

where

$$
\begin{aligned}
& l_{0}=\sqrt{R^{2}+\left(s_{o}+R\right)^{2}-2 R\left(S_{o}+R\right) \cos \phi} \\
& l_{i}=\sqrt{R^{2}+\left(s_{i}-R\right)^{2}+2 R\left(S_{i}-R\right) \cos \phi},
\end{aligned}
$$

By Fermat's Principle

$$
\frac{d(O P L)}{d \phi}=0 \Rightarrow \frac{n_{1} R\left(s_{o}+R\right) \sin \phi}{2 l_{o}}-\frac{n_{2} R\left(s_{i}-R\right) \sin \phi}{2 l_{i}}=0
$$

or

$$
\frac{n_{1}}{l_{o}}+\frac{n_{2}}{l_{i}}=\frac{1}{R}\left(\frac{n_{2} s_{i}}{l_{i}}-\frac{n_{1} s_{o}}{l_{o}}\right)
$$

which means that not all rays from $S$ pass $P$.

## Paraxial Approximation

For small $\phi, l_{i} \approx s_{i}, l_{o} \approx s_{o}$, the equation becomes

$$
\frac{n_{1}}{s_{o}}+\frac{n_{2}}{s_{i}}=\frac{n_{2}-n_{1}}{R}
$$

which is independent of the OPL, so all rays pass through P.
Define object focal length $f_{o}=\frac{n_{1}}{n_{2}-n_{1}} R \equiv s_{o}$ when $s_{i}=\infty$.
Define image focal length $f_{i}=\frac{n_{2}}{n_{2}-n_{1}} R \equiv s_{i}$ when $s_{o}=\infty$


## Sign conventions

TABLE 5.1 Sign Convention for Spherical
Refracting Surfaces and Thin Lenses*
(Light Entering from the Left)

| $s_{o}, f_{o}$ | + left of $V$ |
| :--- | :--- |
| $x_{o}$ | + left of $F_{o}$ |
| $s_{i}, f_{i}$ | + right of $V$ |
| $x_{i}$ | + right of $F_{i}$ |
| $R$ | + if $C$ is right of $V$ |
| $y_{o}, y_{i}$ | + above optical axis |

*This table anticipates the imminent introduction of a few quantities not yet spoken of

## Thin Lenses



From the first surface,

$$
\frac{n_{m}}{s_{o l}}+\frac{n_{l}}{s_{i l}}=\frac{n_{l}-n_{m}}{R_{1}} .
$$

For simplicity, assume $s_{i 1}$ is negative, i.e., a virtual image.
Since $\left|s_{o 2}\right|=\left|s_{i l}\right|+d$, by the sign convection, $s_{o 2}=-s_{i 1}+d$.
Thus, from the second surface,

$$
\frac{n_{p}}{-s_{i 1}+d}+\frac{n_{m}}{s_{i 2}}=\frac{n_{m}-n_{l}}{R_{2}}
$$

Again, by sign convention, $R_{2}<0$. Adding the two equations, we have

$$
\frac{n_{m}}{s_{o l}}+\frac{n_{m}}{s i_{2}}=\left(n_{l}-n_{m}\right)\left(\frac{1}{R_{1}}-\frac{1}{R_{2}}\right)+\frac{n_{l} d}{\left(s_{i l}-d\right) s_{i l}}
$$

By thin lens approximation ( $d \rightarrow 0$ ), we have

$$
\frac{1}{s_{o}}+\frac{1}{s_{i}}=\left(\frac{n_{l}}{n_{m}}-1\right)\left(\frac{1}{R_{1}}-\frac{1}{R_{2}}\right)
$$

This is Thin-Lens Equation, or Lensmaker's Formula.
Since $\lim s_{i}=f=\lim s_{o}$, we can write the equation as

$$
\frac{1}{s_{o}}+\frac{1}{s_{i}}=\frac{1}{f}
$$

where

$$
\frac{1}{f}=\left(\frac{n_{l}}{n_{m}}-1\right)\left(\frac{1}{R_{1}}-\frac{1}{R_{2}}\right) .
$$

$f$ is the focal point.


## Focal Points and Planes



Figure 5.22 Object and image location for a thin lens.

Transverse Magnification $M_{T}=\frac{y_{i}}{y_{o}}=-\frac{s_{i}}{s_{o}}=-\frac{x_{i}}{f}=-\frac{f}{x_{o}}$

| Quantity | Sign |  |
| :---: | :---: | :---: |
|  | + | - |
| $s$ | Real object | Virtual object |
| $s_{i}$ | Real image | Virtual image |
| $f$ | Converging lens | Diverging lens |
| $y_{0}$ | Erect object | Inverted object |
| $y_{i}$ | Erect image | Inverted image |
| $M_{T}$ | Erect image | Inverted image |

## TABLE 5.3 Images of Real Objects Formed by <br> Thin Lenses

| Convex |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Object |  |  | Image |  |
| Location | Type | Location | Orientation | Relative Size |
| $\infty>s_{o}>2 f$ | Real | $f<s_{i}<2 f$ | Inverted | Minified |
| $s_{o}=2 f$ | Real | $s_{i}=2 f$ | Inverted | Same size |
| $f<s_{o}<2 f$ | Real | $\infty>s_{i}>2 f$ | Inverted | Magnified |
| $s_{o}=f$ |  | $\pm \infty$ |  |  |
| $s_{o}<f$ | Virtual | $\left\|s_{i}\right\|>s_{o}$ | Erect | Magnified |
| Concave |  |  |  |  |
| Object |  |  | Image |  |
| Location | Type | Location | Orientation | Relative Size |
| Anywhere | Virtual | $\begin{aligned} \left\|s_{i}\right\| & <\|f\|, \\ s_{o} & >\left\|s_{i}\right\| \end{aligned}$ | Erect | Minified |

Longitudinal Magnification

$$
M_{L} \equiv \frac{d x_{i}}{d x_{o}}=-\frac{f^{2}}{x_{o}^{2}}=-M_{T}^{2}
$$



## Thin-Lens Combinations

For $L_{1}$,

$$
\frac{1}{s_{i l}}=\frac{1}{f_{1}}-\frac{1}{s_{o l}}
$$

Let $s_{o l}>f_{1}$ and $f_{1}>0$.
For $L_{2}$

$$
\begin{aligned}
& s_{o 2}=d-s_{i 1} \\
& \frac{1}{s_{i 2}}=\frac{1}{f_{2}}-\frac{1}{s_{o 2}}
\end{aligned}
$$

Thus,

$$
s_{i 2}=\frac{f_{2} d-f_{2} s_{o o} f_{1} /\left(s_{o l}-f_{1}\right)}{d-f_{2}-s_{o f} f_{1} /\left(s_{o l}-f_{1}\right)}
$$

and

$$
M_{T}=M_{T 1} M_{T 2}=\frac{f_{1} s_{i 2}}{d\left(s_{o l}-f_{1}\right)-s_{o o} f_{1}}
$$

Front focal length:

$$
\left.\frac{1}{s_{o l}}\right|_{s_{i 2}=\infty}=\frac{1}{f_{1}}-\frac{1}{d-f_{2}}=\frac{d-\left(f_{1}+f_{2}\right)}{f_{1}\left(d-f_{2}\right)} \Rightarrow b f l=\frac{f_{1}\left(d-f_{2}\right)}{d-\left(f_{1}+f_{2}\right)}
$$

Back focal length:

$$
\frac{1}{\left.s_{i 2}\right|_{s_{o l}=\infty}}=\frac{1}{f_{2}}-\frac{1}{d-f_{1}}=\frac{d-\left(f_{1}+f_{2}\right)}{f_{2}\left(d-f_{1}\right)} \Rightarrow b f l=\frac{f_{2}\left(d-f_{1}\right)}{d-\left(f_{1}+f_{2}\right)}
$$

If $d \rightarrow 0, f=b f l=f f l=\frac{f_{2} f_{1}}{f_{2}+f_{1}} \Rightarrow \frac{1}{f}=\frac{1}{f_{1}}+\frac{1}{f_{2}}$

## Stops (5.3)

Aperture stop: any element that determines the amount of light reaching the image.
Field stop: the element limiting the size or angular breadth of the object that can be imaged $y$ the system.

Field Stop: The size of the image plane.


Figure 5.33 Aperture stop and field stop.

Pupils: for determining the light of cone entering the image plane.
Entrance Pupil: the image of the aperture stop as seen from an axial point on the object side through those elements before the stop. Determine the cone of the light entering the image plane with respect to the object.

Exit Pupil: the image of the aperture stop as seen from an axial point on the image side through those elements before the stop. Determine the cone of the light entering the image plane with respect to the image plane.
$E_{x p}$ : the point at center of the exit pupil.
$E_{n p}$ : the point at center of the entrance pupil.

Chief Ray: any ray from an off-axis object point that passes through the center of the aperture stop. Either this ray or its extending line passes $E_{x p}$ or $E_{n p}$.

