

LECTURE #25

Note Title

3/22/2010

HARTREE - FOCK MEAN FIELD

CORRELATION EFFECTS = EFFECTS BEYOND MEAN FIELD

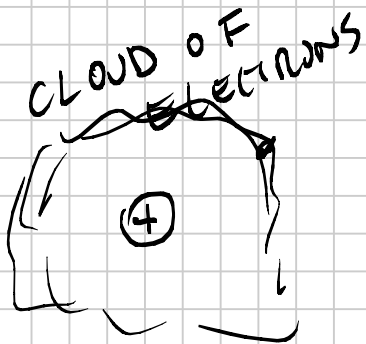
SCREENING

$\cdot e$ TEST CHARGE

$$\phi^{\text{ext}} = -\frac{e^2}{r}$$

⊕

IN THE PRESENCE OF ELECTRON GAS



TEST CHARGE

$$\phi(r) \neq \phi^{\text{ext}}$$

$$\rho(\vec{r}) = \rho^{\text{ext}}(\vec{r}) + \rho^{\text{INDUCED}}(r)$$

D DISPLACEMENT VECTOR

$$\nabla \cdot \mathbf{D} = 4\pi \rho^{\text{ext}}$$

$$\mathbf{D} = \overset{\text{ext}}{\epsilon} \cdot \overset{\text{ext}}{\mathbf{E}}$$

↓
 ϕ^{ext}

→ ϕ

→ ASSUME ISOTROPIC SYSTEM

$$\epsilon_{ij} = \epsilon \delta_{ij}$$

$$\phi^{\text{ext}}(\vec{r}) = \int d\mathbf{r}' \epsilon(\mathbf{r}-\mathbf{r}') \phi(\mathbf{r}')$$

→ NON LOCAL DEPENDENCE

↓ FT

$\epsilon(\mathbf{r}-\mathbf{r}') \rightarrow$ TRANSLATIONAL INVARIANCE

$$\phi^{\text{ext}}(q) = \epsilon(q) \phi(q)$$

$$\int f(x) g(y-x) dx$$

↓ FT

$$\tilde{f}(q) \tilde{g}(q)$$

$$\phi(q) = \frac{\phi^{\text{ext}}(q)}{\epsilon(q)}$$

$\epsilon(q) = \epsilon_0$ CONSTANT IN q -SPACE

↓
 $\epsilon(r-r') = \epsilon_0 \delta(r-r')$

$$\phi^{\text{ext}} = -\frac{e^2}{r} \rightarrow \phi = -\frac{e^2}{\epsilon_0 r}$$

$$\phi^{\text{ext}}(q) = \frac{4\pi e^2}{|q|^2} \rightarrow \frac{4\pi e^2}{\epsilon_0 |q|^2}$$

$$\vec{P} = \chi \vec{E}_0$$

$$\boxed{\epsilon(q) \leftrightarrow \chi(q)}$$

↓
 $\rho^{\text{INDUCED}}(q) = \chi(q) \phi(q)$

$$-\nabla^2 \phi(r) = 4\pi \rho(r)$$

$$\rho(r) = \rho^{\text{ext}} + \rho^{\text{IND}}$$

$$-\nabla^2 \phi^{\text{ext}}(r) = 4\pi \rho^{\text{ext}}(r)$$

↓ FT

$$q^2 \phi(q) = 4\pi \rho^{\text{ext}}(q) + 4\pi \rho^{\text{IND}}(q)$$

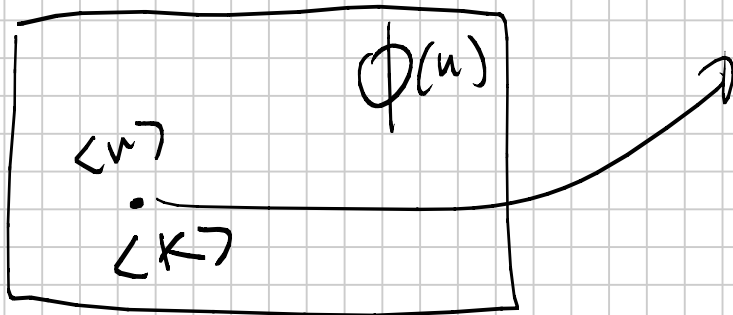
$$q^2 (\phi(q) - \phi^{\text{ext}}(q)) = 4\pi \chi(q) \phi(q)$$

$$q^2 (\cancel{\phi(q)} - \epsilon(q) \cancel{\phi(q)}) = 4\pi \chi(q) \cancel{\phi(q)}$$

$$\epsilon(q) = 1 - \frac{4\pi}{q^2} \chi(q) \rightarrow 1 - \frac{4\pi}{q^2} \frac{\rho^{\text{IND}}(q)}{\phi(q)}$$

THOMAS - FERMI APPROXIMATION

$\phi(u)$ WEAK DEPENDENCE ON r



ELECTRON AT $\langle u \rangle$
WITH AVERAGE $\langle k \rangle$
AND ENERGY

$$\epsilon(k) = \frac{\hbar^2 \langle k \rangle^2}{2m}$$

$$\phi(u) = 0$$

$$n(\mu) = 2 \int \frac{d^3 k}{(2\pi)^3} \frac{1}{e^{\beta(\epsilon(k) - \mu)} + 1}$$

\downarrow
 $\epsilon(k) = \frac{\hbar^2 k^2}{2m}$

ADD $\phi(u)$

$$E(r, k) \rightarrow \frac{\hbar^2 k^2}{2m} - e\phi(u) \rightarrow \text{MAKES SENSE BECAUSE OF SEMICLASSICAL APPROX}$$

$$n'(r, \mu) = 2 \int \frac{d^3 k}{(2\pi)^3} \frac{1}{e^{\beta \left(\frac{\hbar^2 k^2}{2m} - \underbrace{e\phi(u)}_{\mu'(u) = \mu + e\phi(u)} - \mu \right)}}$$

$$n' = n(\mu + e\phi(u))$$

$$\rho^{\text{INDUCED}}(u) = -e [n(\mu + e\phi(u)) - n(\mu)]$$

$$e\phi(\mu) \ll \mu$$

$$\rho^{IND}(\mu) = -e \frac{dM}{d\mu} e\phi(\mu)$$

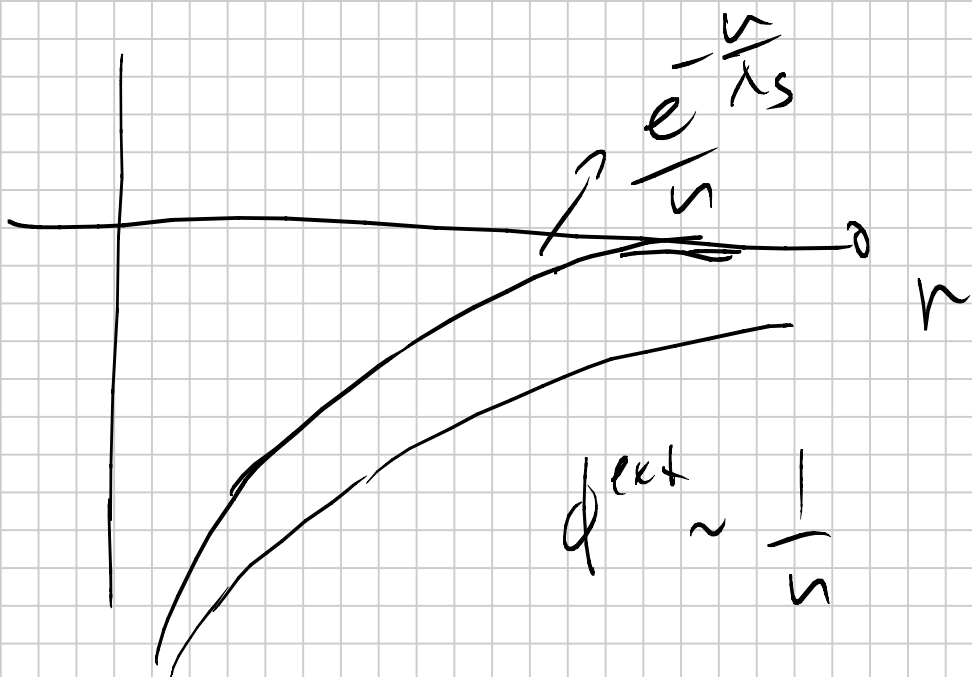
$$\epsilon^{TF}(q) = 1 - \frac{4\pi}{q^2} \frac{\rho^{IND}(q)}{\phi(q)} = 1 - \frac{4\pi}{q^2} \left(-e^2 \frac{dM}{d\mu} \right) =$$

$$1 + \frac{\left(4\pi e^2 \frac{dM}{d\mu} \right)}{q^2} = 1 + \frac{q_s^2}{q^2} \rightarrow q \text{ SCREENING}$$

$$d^{\text{ext}}(q) = e^2 \frac{4\pi}{q^2} \rightarrow \phi(q) = \frac{\phi^{\text{ext}}(q)}{\epsilon(q)} = \frac{4\pi e^2}{q^2 + q_s^2}$$

$$\phi^{\text{ext}}(u) = \frac{e^2}{u}$$

$$\phi^{\text{TF}}(u) = \frac{e^2}{u} e^{-q_s u} \rightarrow \text{YUKAWA POTENTIAL}$$



$$q_s = 4\pi e^2 \frac{dm}{d\mu}$$

$$T = 0$$



$$m = \int_0^{\infty} d\epsilon g(\epsilon) \vartheta(\mu - \epsilon)$$

$$\frac{dm}{d\mu} = \int d\epsilon g(\epsilon) \frac{d}{d\mu} \vartheta(\mu - \epsilon) = \int d\epsilon g(\epsilon) \delta(\epsilon - \mu)$$

$$= g(\mu)$$

$$q_s \propto K_F \left(\frac{r_s}{a_B} \right)^{1/2} \sim \left(\frac{a_B}{r_s} \right)^{1/2}$$

$a_B =$ BOHR
RADIUS

