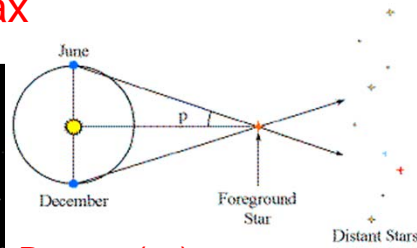


Measuring the Distances to the Stars: Parallax



Parsec (pc)

= distance to star for which
angle $p = 1$ arcsec
= 3.26 LY

Hwk 2 due Sept 23

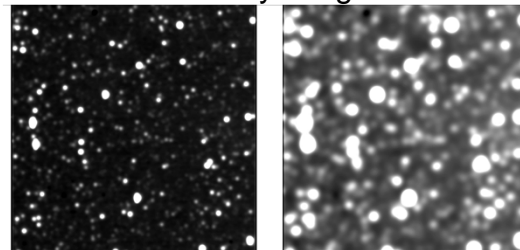
Hwk 3 not yet
assigned, but due
Sept 30

Midterm 1
Wed, Oct 2

<http://sci2.esa.int/interactive/media/html/sec23p1.htm>

What sets the parallax limit?

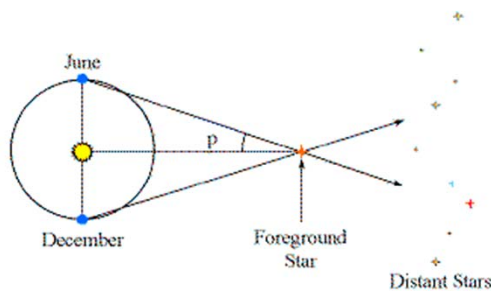
Fuzzy images



A field of stars

...blurred by Earth's atmosphere.

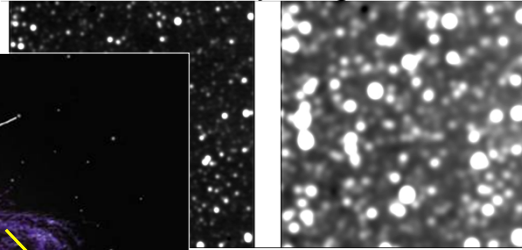
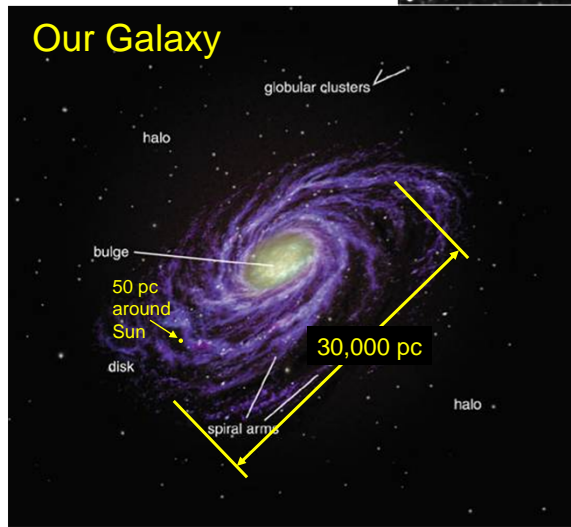
Old limit for parallax distances:
20-50 parsecs



What sets the parallax limit?

Fuzzy images

Our Galaxy

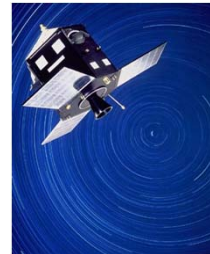
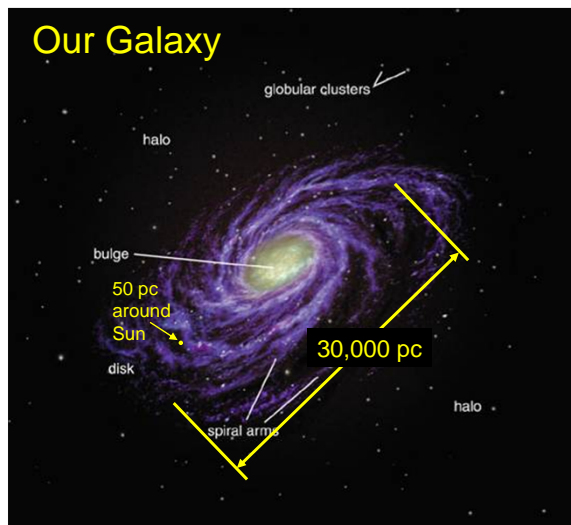


...blurred by Earth's atmosphere.

Old limit for parallax distances:
20-50 parsecs

Sharp Images from Space

Our Galaxy



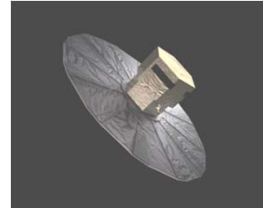
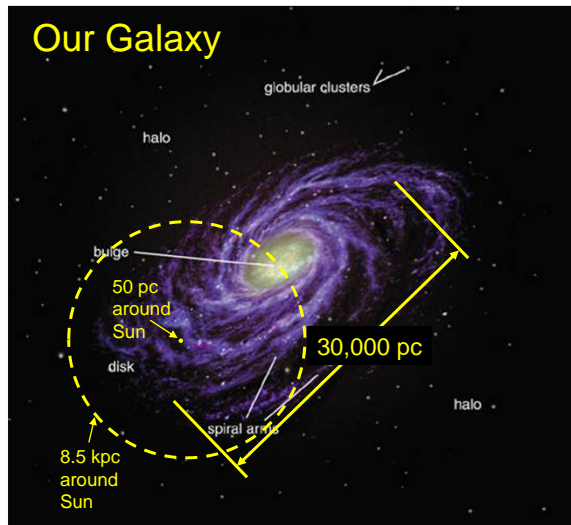
Hipparcos

Old limit for parallax distances:
20-50 parsecs

Hipparcos (1989-1993):
100-200 parsecs
($1\sigma = 1$ milliarcsec = 1kpc)

Coming Soon

Our Galaxy



GAIA spacecraft: Dec 2011 launch

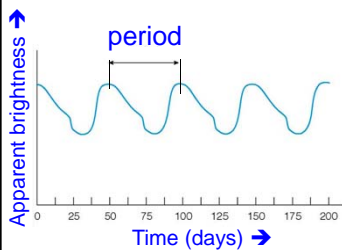
Old limit for parallax distances:
20-50 parsecs

Hipparcos (1989-1993):
100-200 parsecs
($1\sigma = 1$ milliarcsec = 1kpc)

GAIA: 8.5 kpc

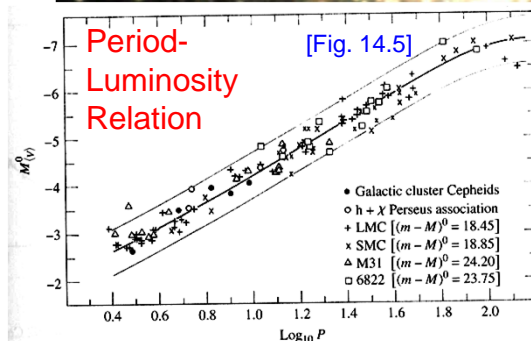
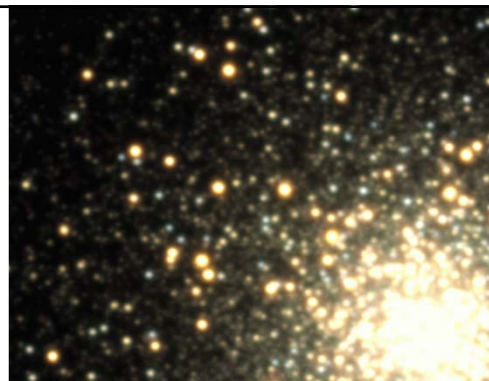
Pulsating Variable Stars

- These stars regularly expand & contract.
- Like a big spring.
- Change in size →
 - change in temperature
 - change in luminosity



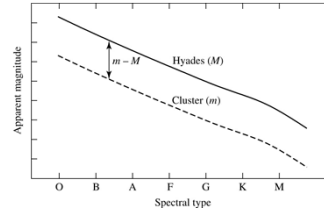
P-L relation

- discovered in Magellanic Clouds
- calibrated locally, using (statistical) parallaxes

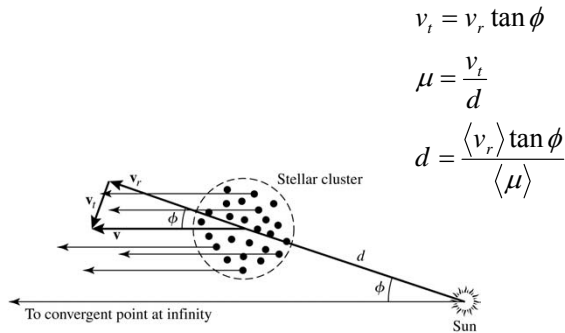
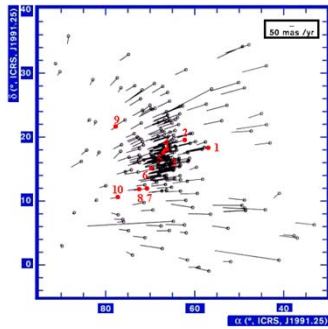


Measuring Distances inside the MW

- Parallax
- Pulsating variables
- Main sequence fitting for clusters



- Calibrate with Hyades (moving cluster method = pp. 919-922)



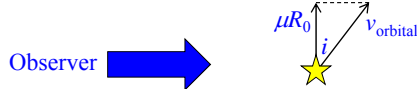
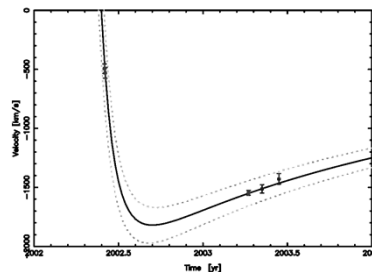
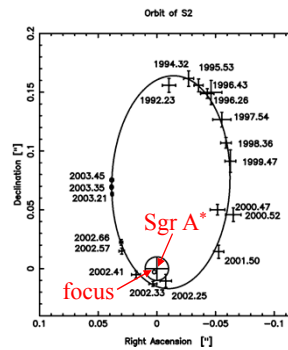
$$v_t = v_r \tan \phi$$

$$\mu = \frac{v_t}{d}$$

$$d = \frac{\langle v_r \rangle \tan \phi}{\langle \mu \rangle}$$

Distance R_0 to the Galactic Center

- Massive ($> 10^6 M_{\text{sun}}$) black hole at G.C.
- Can follow orbits of stars around it.
- Measured radial velocity (v_r) of one star at several points in its orbit, and proper motion μ (arcsec yr^{-1}).
- Orbit fit gives i (inclination), and predicts orbital velocity at each point.
- In concept, solve for $R_0 = 7.94 \pm 0.38 \text{ kpc}$
 - (actually, it comes out of orbit fit)



Eisenhauer et al. 2003, ApJ, 597, L121

Kinematics of the Milky Way

- From [CO] 24.3, especially pp 901-14.

- Coordinate systems

- Galactic latitude (b_{Π}), longitude (l_{Π})
 - Spherical coordinates centered on Sun

- R, θ, z

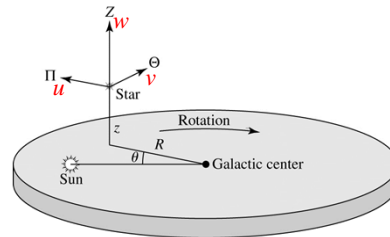
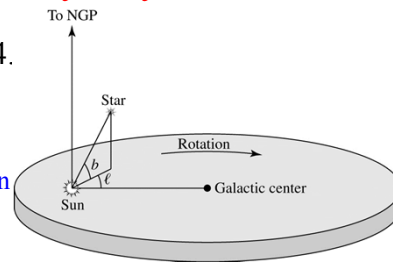
- Cylindrical coordinate system centered on Galactic Center

- Π, Θ, Z

- Velocity components in R, θ, z system.

- Peculiar velocities u, v, w

- Π, Θ, Z velocities but relative to *Local Standard of Rest*
- LSR is point instantaneously centered on Sun, but moving in a perfectly circular orbit.
- *Solar motion* = motion of sun relative to LSR



- Star density is higher towards GC.
 - Those stars are on orbits that Sun overtakes.

- Velocity ellipsoids and asymmetrical drift.

→ LSR's orbital velocity

$$\Theta_0(R_0) = 220 \text{ km s}^{-1}$$

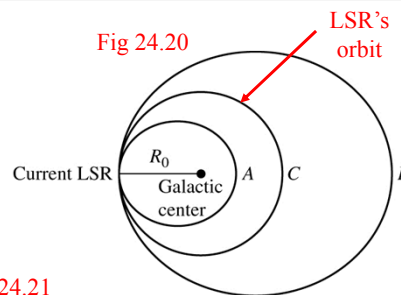
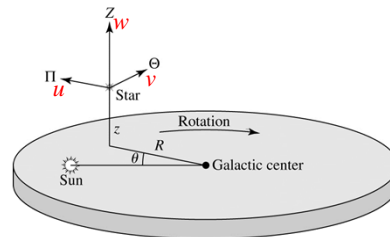
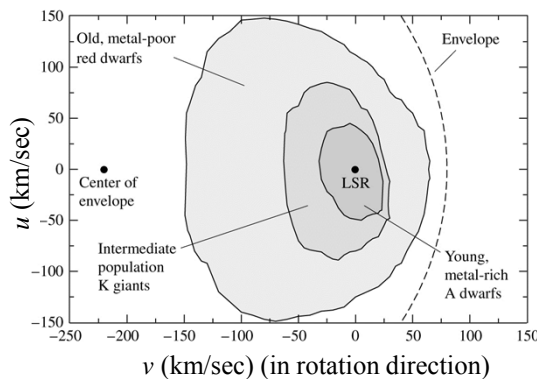


Fig 24.21

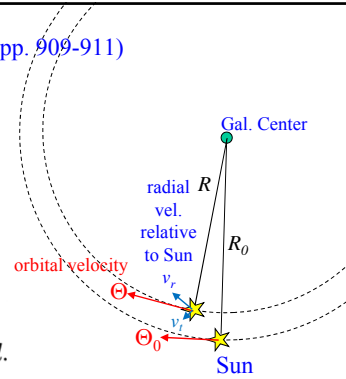


“Envelope” = halo stars

- Net rotation = 0

Differential Rotation (see [CO pp. 909-911])

- From the figure: $v_r = \Theta \cos \alpha - \Theta_0 \sin \ell$,
 $v_t = \Theta \sin \alpha - \Theta_0 \cos \ell$,
- Angular rotation velocity: $\Omega(R) \equiv \frac{\Theta(R)}{R}$
- + some geometry \rightarrow $v_r = (\Omega - \Omega_0) R_0 \sin \ell$,
 $v_t = (\Omega - \Omega_0) R_0 \cos \ell - \Omega d$.
- Taylor expansion:



$$\Omega(R) = \Omega_0(R_0) + \left. \frac{d\Omega}{dR} \right|_{R_0} (R - R_0) + \dots$$

$$v_r \simeq Ad \sin 2\ell,$$

$$v_t \simeq Ad \cos 2\ell + Bd$$

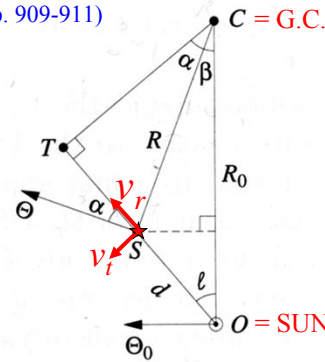
Oort's Constants:

$$A \equiv -\frac{1}{2} \left[\left. \frac{d\Theta}{dR} \right|_{R_0} - \frac{\Theta_0}{R_0} \right]$$

$$B \equiv -\frac{1}{2} \left[\left. \frac{d\Theta}{dR} \right|_{R_0} + \frac{\Theta_0}{R_0} \right]$$

Differential Rotation (see [CO pp. 909-911])

- From the figure: $v_r = \Theta \cos \alpha - \Theta_0 \sin \ell$,
 $v_t = \Theta \sin \alpha - \Theta_0 \cos \ell$,
- Angular rotation velocity: $\Omega(R) \equiv \frac{\Theta(R)}{R}$
- + some geometry \rightarrow $v_r = (\Omega - \Omega_0) R_0 \sin \ell$,
 $v_t = (\Omega - \Omega_0) R_0 \cos \ell - \Omega d$.
- Taylor expansion:



[Fig 24.22]

$$\Omega(R) = \Omega_0(R_0) + \left. \frac{d\Omega}{dR} \right|_{R_0} (R - R_0) + \dots$$

$$v_r \simeq Ad \sin 2\ell,$$

$$v_t \simeq Ad \cos 2\ell + Bd$$

Oort's Constants:

$$A \equiv -\frac{1}{2} \left[\left. \frac{d\Theta}{dR} \right|_{R_0} - \frac{\Theta_0}{R_0} \right]$$

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Oort's Constants

$$A \equiv -\frac{1}{2} \left[\frac{d\Theta}{dR} \Big|_{R_0} - \frac{\Theta_0}{R_0} \right]$$

$$B \equiv -\frac{1}{2} \left[\frac{d\Theta}{dR} \Big|_{R_0} + \frac{\Theta_0}{R_0} \right]$$

$$v_r \simeq Ad \sin 2\ell$$

$$v_t \simeq Ad \cos 2\ell + Bd$$

$$\Omega(R) = \frac{\Theta(R)}{R}$$

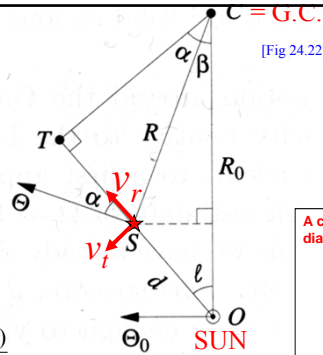
- Evaluate A, B from observations:
- Contain important information about Galactic rotation curve.

- Angular velocity for circular motion at R_0 :

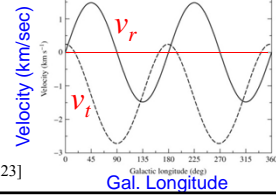
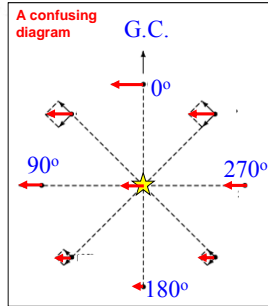
$$\Omega_0 = A - B = 27.2 \text{ km s}^{-1} \text{ kpc}^{-1}$$

- Gradient of rotation curve at R_0 :

$$\frac{d\Theta}{dR} \Big|_{R_0} = -(A + B) = -2.4 \text{ km s}^{-1} \text{ kpc}^{-1}$$



[Fig 24.22]



[CO Fig 24.23]

- Distance to Galactic Center

$$R_0 \sim 8 \text{ kpc}$$

- Sun's orbital velocity

$$\Theta_0 = \Omega_0 \times R_0 \sim 220 \text{ km s}^{-1}$$

- Sun's orbital period

$$P = \frac{R_0}{\Theta_0} \sim 230 \text{ million yrs.}$$

- Approx. mass interior to Sun's orbit:

$$F_{\text{centrifugal}} = F_{\text{grav}}$$

$$\frac{m\Theta_0^2}{R_0} = \frac{GmM}{R_0^2}$$

$$M = \frac{\Theta_0^2 R_0}{G} \sim 9 \times 10^{10} M_{\odot}$$

Rotation Curves

- Spherical mass shell, uniform density

Inside shell: no effect

Outside shell:

acts as if all mass at center

$$\frac{dF_{gravity}}{d\Omega} \propto \frac{\rho \times d(vol)}{r^2}$$

$$d(vol) = d\Omega r^2 dr$$

- $F_{centripital} = F_{gravity}$

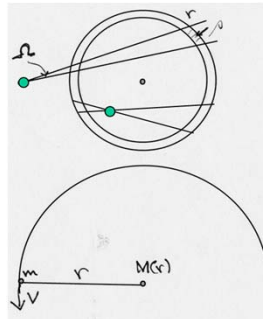
$$\frac{mv^2}{r} = \frac{GM(r)m}{r^2}$$

$$v = \left(\frac{GM(r)}{r} \right)^{1/2}$$

Where

v = space velocity

= Θ if circular orbit.



- Inside spherical mass distribution

$$M(r) = \int_0^r \rho(r') 4\pi r'^2 dr' = \bar{\rho} \frac{4}{3} \pi r^3$$

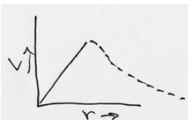
$$v \propto r$$



- Outside spherical mass distribution

$$M(r) = \text{const.}$$

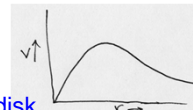
$$v \propto r^{-1/2}$$



- Spher. Distr. + exponential disk

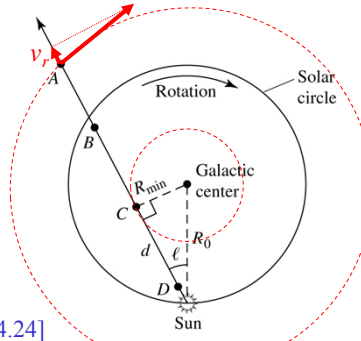
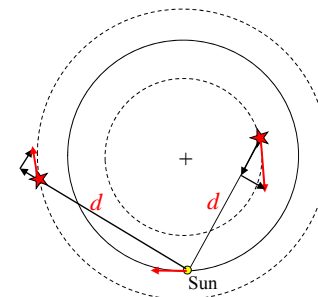
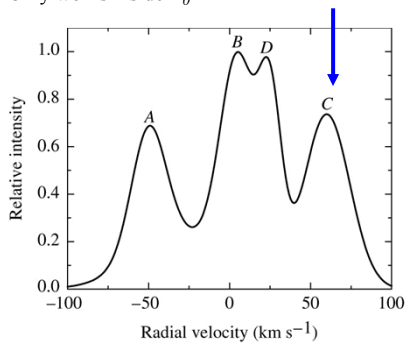
$$V = \left(\frac{G}{F} \right)^{1/2} \left\{ M(r) + 4m_d s^2 \left[I_0(s) K_0(s) - I_1(s) K_1(s) \right] \right\}^{1/2}$$

$$s = \frac{1}{2} r/h \quad I, K = \text{Bessel functions}$$



Measuring the MW Rotation Curve

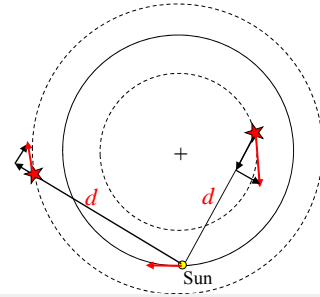
- In principle, for stars, clusters, etc:
 - measure distance d and v_r
 - v_r = radial velocity w.r.t. Sun
 - assume circular orbit
- For H I 21cm, CO, etc. radio emission:
 - Only can measure v_r
 - Use tangent point method
 - Only works inside R_0



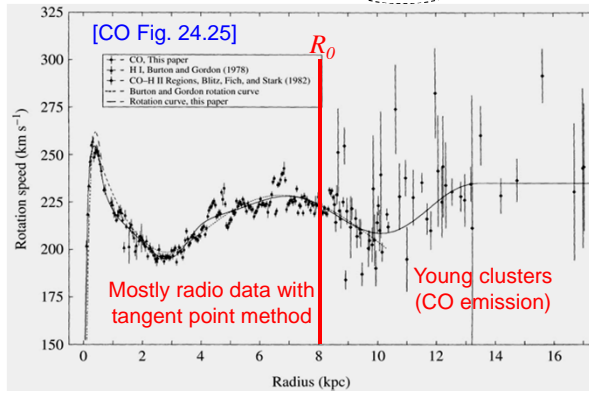
[CO fig 24.24]

Measuring the MW Rotation Curve

- In principle, for stars, clusters, etc:
 - measure distance d and v_r
 - v_r = radial velocity w.r.t. Sun
 - assume circular orbit
- For H I 21cm, CO, etc. radio emission:
 - Only can measure v_r
 - Use tangent point method
 - Only works inside R_0



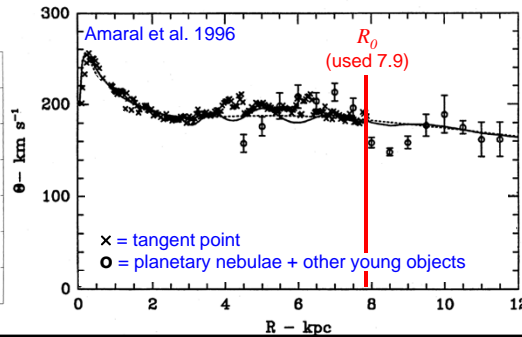
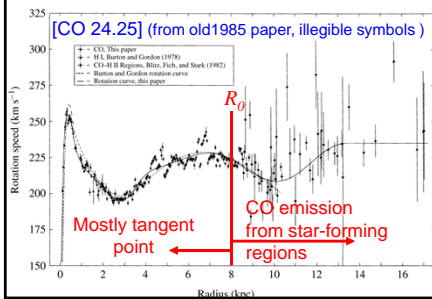
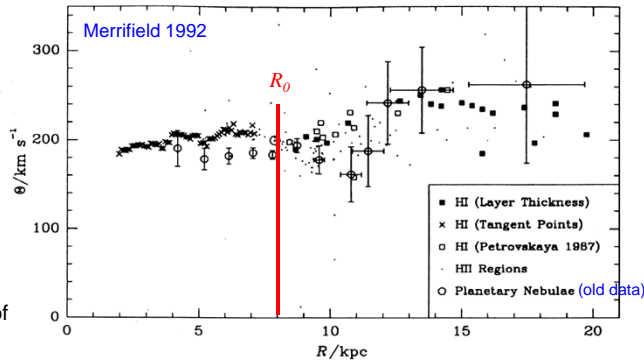
- Outside R_0**
 - Must use other methods.
 - Must know actual distance + velocity.

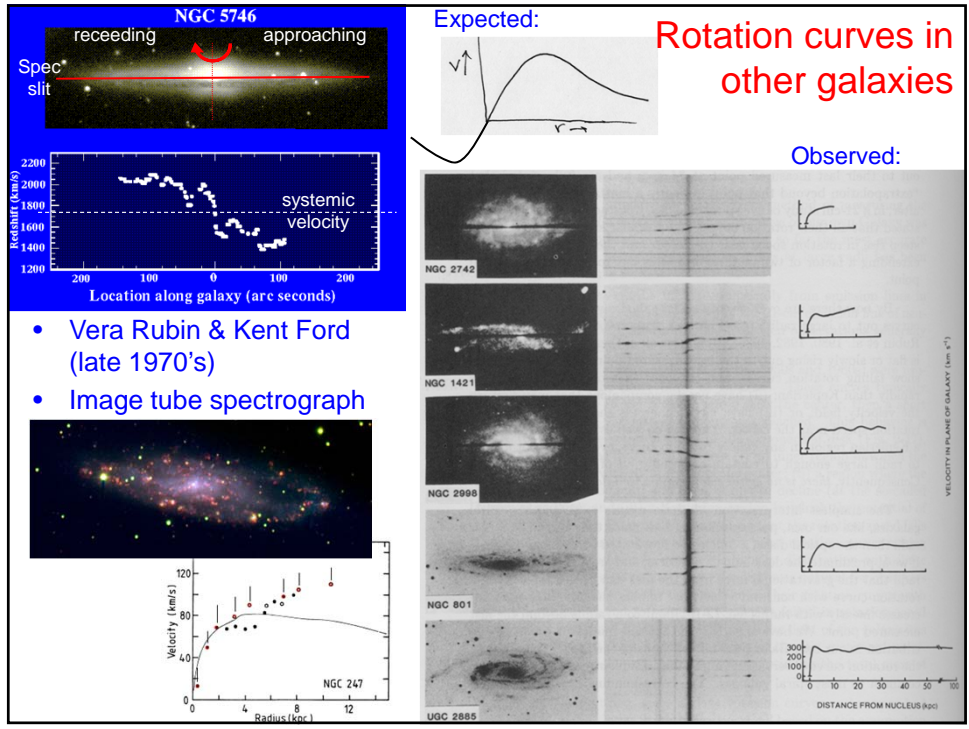


Galactic Rotation Curves

(take your pick)

- Most data points interior to R_0 are from tangent point method.
- Beyond R_0 most data points are young objects with known distances & motions.
- But Merrifield used variation of apparent thickness of H I disk with longitude (complicated).





[CO pg. 917] Density as shown by flat rotation curves

Back to $F_{\text{CENTRIFUGAL}} = F_{\text{GRAVITATIONAL}}$

$$\frac{m v^2}{r} = \frac{G M(r) m}{r^2}$$

$$M(r) = \frac{v^2 r}{G}$$

$$\frac{dM(r)}{dr} = \frac{v^2}{G}$$

but also $\frac{dM(r)}{dr} = 4\pi r^2 \rho(r)$

$$\Rightarrow \rho(r) = \frac{v^2}{4\pi G r^2}$$

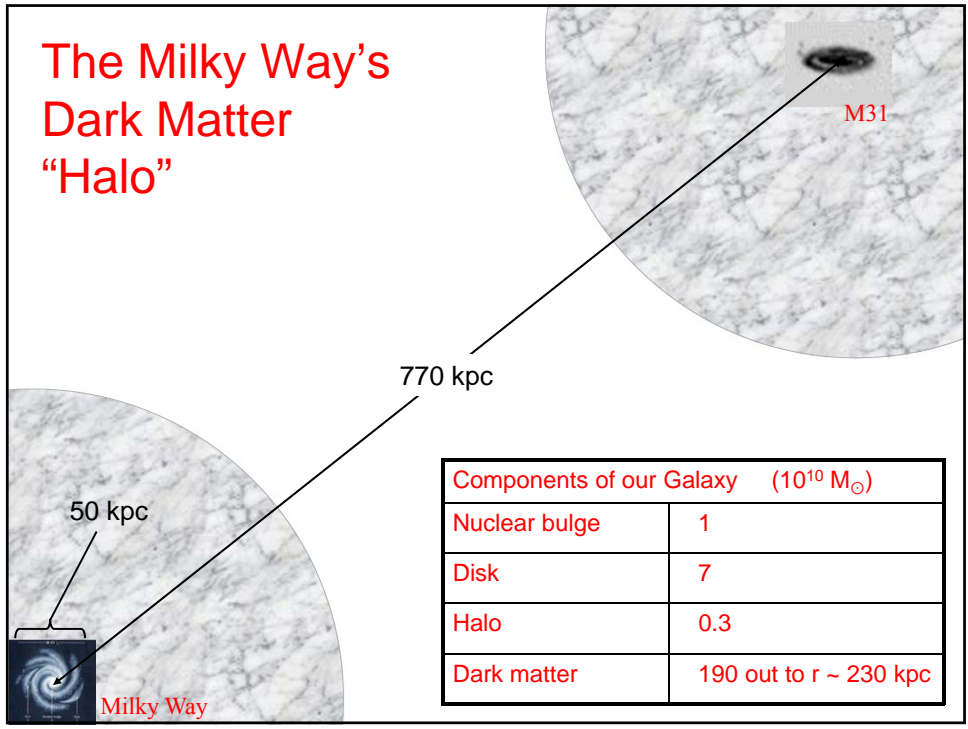
$v \sim \text{constant} \Rightarrow \rho(r) \propto \frac{1}{r^2}$

Use $\rho(r) = \frac{C_0}{(a^2 + r^2)}$

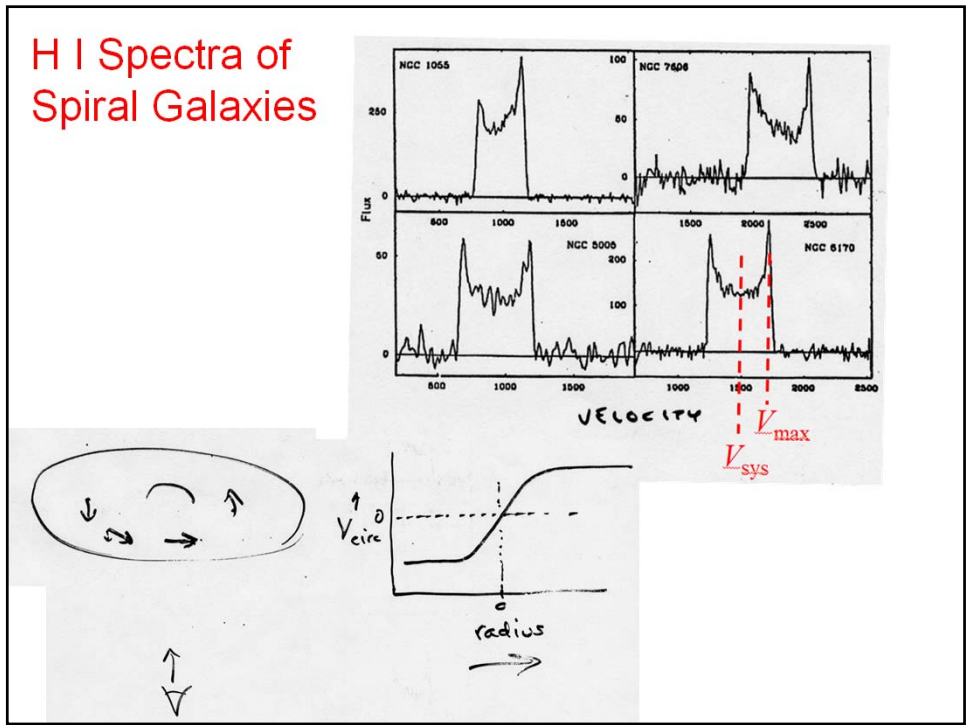
$$\rho(r) = \frac{\rho_0}{(r/a)(1+r/a)^2}$$

- $dM(r)/dr \sim \text{constant}$
 - Unbounded mass distribution??
- NFW profile
 - Predicted for Cold dark matter (CDM)
 - Actual derived dark matter profiles often slightly different than this
 - What is CDM? Coming later in course.

The Milky Way's Dark Matter "Halo"

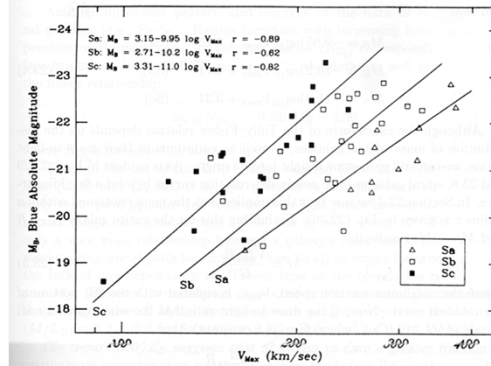
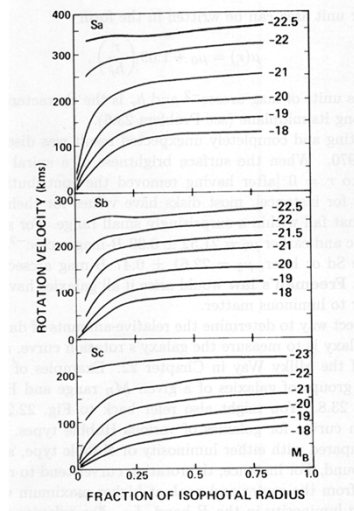


H I Spectra of Spiral Galaxies



Tully-Fisher Relation

Rotation Curves



- Maximum rotation velocity-Luminosity relation [FIG 25.10]
 - **Tully-Fisher** relation
 - $M_B = -9.95 \log_{10} V_{\text{Max}} + 3.15$ (Sa)
 - $M_B = -10.2 \log_{10} V_{\text{Max}} + 2.71$ (Sb)
 - $M_B = -11.0 \log_{10} V_{\text{Max}} + 3.31$ (Sc)

Semi-derivation of Tully-Fisher Relation:

$$M_B = -9.95 \log_{10} V_{\text{max}} + 3.15 \quad (\text{Sa})$$

$$M_B = -10.2 \log_{10} V_{\text{max}} + 2.71 \quad (\text{Sb})$$

$$M_B = -11.0 \log_{10} V_{\text{max}} + 3.31 \quad (\text{Sc})$$

- Mass interior to outermost R where rotation curve can be measured:

$$Mass = \frac{V_{\text{max}}^2 R}{G}$$

- Assume $L = Mass / \text{const.}$
- “Freeman Law” (observed fact ---maybe):

$$\text{Surf. Bright.} = \frac{L}{4\pi R^2} = \text{const.}$$

$$L = \text{const} \times V_{\text{max}}^4$$

- Convert to Absolute B-band magnitudes:

$$M_B = M_{\text{sun}} - 2.5 \log_{10} \left(\frac{L}{L_{\text{sun}}} \right) = -10 \log_{10} V_{\text{max}} + \text{const.}$$

Important as a
DISTANCE
calibrator!

SO FAR:

- Galaxy types
- Ancient history
- Milky Way and spiral galaxy morphology
 - Nuclear bulge
 - Disk
 - Stellar halo
 - Dark matter halo
- Star-forming regions
- Chemical enrichment

Hwk 3 Due Sept 29
CO 24.15 – just part (a).
CO 24.21
CO 24.36 (a),(b)
CO 25.14
CO 25.16

Distance measurements within MW

Kinematics of spiral galaxies

- Rotation curves → mass distribution
- **Spiral structure [CO 25.3]**
- General properties of S, E, Irr galaxies
- Midterm 1 (Tu. Oct 4)