

## Some Universes

**Open vs. Closed:**

$$\left(H^2 - \frac{8}{3}\pi G\rho\right)R^2 = -kc^2$$

$$H^2(1 - \Omega)R^2 = -kc^2$$

$$H_0^2(1 - \Omega_0) = -kc^2$$

$k = 0 \rightarrow \Omega_0 = \Omega_{m,0} + \overset{\sim 0}{\Omega_{rel,0}} + \Omega_{\Lambda,0} = 1$

**Accelerating vs. Decelerating:**

*Deceleration parameter:*

$$q(t) = -\frac{R(t)\left[\frac{d^2R(t)}{dt^2}\right]}{\left[\frac{dR(t)}{dt}\right]^2}$$

+ accel. eqn:

$$\frac{d^2R}{dt^2} = -\frac{H^2}{2\rho_c} [\rho_m + 2\rho_{rel} - 2\rho_\Lambda] R$$

(Camouflaged accel. eqn.)

$$q(t) = \frac{1}{2}\Omega_m(t) + \overset{\sim 0}{\Omega_{rel}(t)} - \Omega_\Lambda(t)$$

**Expands Forever vs. Recollapses:**

Does  $dR/dt$  ever = 0? [See \[29.135\]](#)

**Definitions, results, etc.**

- \*  $r = R(t)$  or
- \*  $H = \frac{1}{R} \frac{dR}{dt}$

**Densities:**

- \* **Matter:**  $\rho_m = \rho_{o,m} R^{-3}$
- \* **Radiation:**  $\rho_r = \rho_{o,r} R^{-4}$
- \* **Dark energy:**  $\rho_\Lambda = \rho_{o,\Lambda} R^0$

$$\rho_c(t) = \frac{3H^2(t)}{8\pi G}$$

- \*  $\Omega(t) = \frac{\rho(t)}{\rho_c(t)}$
- \*  $\Omega \equiv \Omega_m + \Omega_{rel} + \Omega_\Lambda$

**Physics**

*Per unit mass:*  
K.E. + potential E. = Total Energy

$$\left(\left(\frac{1}{R} \frac{dR}{dt}\right)^2 - \frac{8}{3}\pi G\rho\right)R^2 = -kc^2$$

\*  $\rho = \frac{u}{c^2}$  \*

\* Temp. of radiation field:  $T_0 = RT(R)$  \*

$$(ds)^2 = (cdt)^2 - R^2(t) \left[ \left(\frac{d\varpi}{\sqrt{1-k\varpi^2}}\right)^2 + (\varpi d\theta)^2 + (\varpi \sin\theta d\phi)^2 \right]$$

\*  $\left[\left(\frac{1}{R} \frac{dR}{dt}\right)^2 - \frac{8}{3}\pi G\rho - \frac{1}{3}\Lambda c^2\right]R^2 = -kc^2$  \*

Cosmological Constant (a.k.a. Dark Energy)

Curvature  $k = \frac{1}{R^2} \times \begin{matrix} +1 \\ 0 \\ -1 \end{matrix}$

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$$\frac{d(R^3\rho)}{dt} = -\frac{P}{c^2} \frac{d(R^3)}{dt}$$

$$q(t) = -\frac{R(t)\left[\frac{d^2R(t)}{dt^2}\right]}{\left[\frac{dR(t)}{dt}\right]^2}$$

\*  $P = wu = w\rho c^2$

$$dU = -PdV$$

$$\frac{d^2R}{dt^2} = \left\{ -\frac{4}{3}\pi G \left[ \rho_m + \rho_{rel} + \rho_\Lambda + \frac{3(P_m + P_{rel} + P_\Lambda)}{c^2} \right] \right\} R$$

\* = you should be able to write these down from memory.

## [29.4] Observational Cosmology

$$(ds)^2 = (c dt)^2 - R^2(t) \left[ \left( \frac{d\varpi}{\sqrt{1-k\varpi^2}} \right)^2 + (\varpi d\theta)^2 + (\varpi \sin\theta d\phi)^2 \right]$$

$$\left[ \left( \frac{1}{R} \frac{dR}{dt} \right)^2 - \frac{8}{3}\pi G\rho - \frac{1}{3}\Lambda c^2 \right] R^2 = -kc^2$$

$$H_0^2(1 - \Omega_0) = -kc^2$$

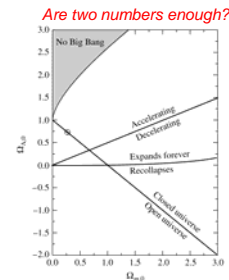
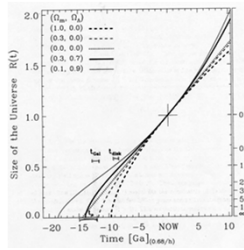
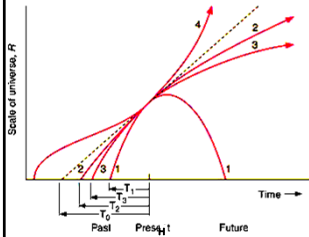
- Some theoretical parameter sets:

- $R(t)$  vs.  $t$
- $\Omega_{\Lambda,0}$  vs.  $\Omega_{m,0}$
- Curvature  $k$ ,  $dR/dt$ ,  $d^2R/dt^2$

- But what can we actually *measure* that will tell us which universe we live in?

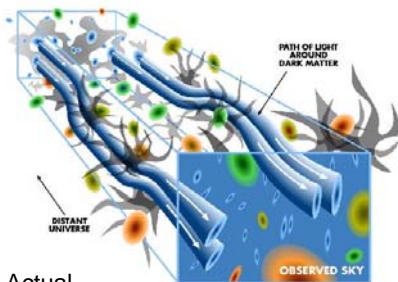
As a function of  $z$ :

- Apparent mag. of standard candles.
- Angular sizes.
- Space density of galaxies.



Are two numbers enough?

## The paths of photons through space-time



Actual  
(sketched in 3D space)

$$(ds)^2 = (c dt)^2 - R^2(t) \left[ \left( \frac{d\varpi}{\sqrt{1-k\varpi^2}} \right)^2 + (\varpi d\theta)^2 + (\varpi \sin\theta d\phi)^2 \right]$$

$$\int_{r_1}^{r_2} \frac{cdt}{R(t)} = - \int_{\varpi_1}^{\varpi_2} \frac{d\varpi}{\sqrt{1-k\varpi^2}} = \int_{\varpi_2}^{\varpi_1} \frac{d\varpi}{\sqrt{1-k\varpi^2}}$$

Take neg. square root so that  $\varpi$  will decrease with increasing  $t$

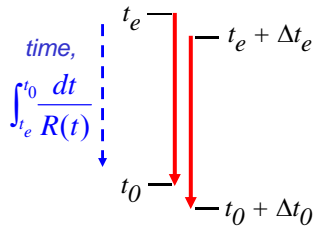
## Redshift and Cosmological Time Dilation

(See pg. 1200)

$$(ds)^2 = 0 = (c dt)^2 - R^2(t) \left[ \left( \frac{d\varpi}{\sqrt{1 - k\varpi^2}} \right)^2 + \dots \right]$$

For two radially travelling wave crests:

$$\int_{t_e}^{t_0} \frac{dt}{R(t)} = \frac{1}{c} \int_0^{\varpi_e} \frac{d\varpi}{\sqrt{1 - k\varpi^2}} = \int_{t_e + \Delta t_e}^{t_0 + \Delta t_0} \frac{dt}{R(t)}$$



Both red paths are comoving distance  $\varpi_e$  long.

$$\int_{t_0}^{t_0 + \Delta t_0} \frac{dt}{R(t)} = \int_{t_e}^{t_e + \Delta t_e} \frac{dt}{R(t)}$$

$$\frac{\Delta t_e}{R(t_e)} = \frac{\Delta t_0}{R(t_0)}$$

$$\frac{R(t_0)}{R(t_e)} = \frac{\Delta t_0}{\Delta t_e}$$

$$\frac{1}{R(t_e)} = \frac{\lambda_0}{\lambda_e} = 1 + z \quad [\text{CO 29.142}]$$

## Proper distance

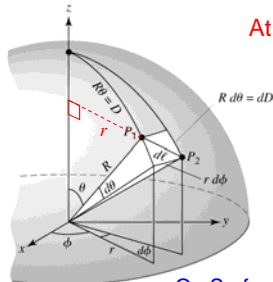
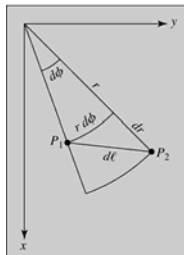
= the *current* distance to a distant object.

$$(ds)^2 = (c dt)^2 - R^2(t) \left[ \left( \frac{d\varpi}{\sqrt{1 - k\varpi^2}} \right)^2 + (\varpi d\theta)^2 + (\varpi \sin \theta d\phi)^2 \right]$$

$dt = 0$ , proper distance  $d_p(t) = \text{sqrt}(-ds^2)$

$$d_p(t) = R(t) \int_0^{\varpi} \frac{d\varpi'}{\sqrt{1 - k\varpi'^2}}$$

On Flat Surface



On Surface of Ball

At the current time (using  $R(t_0) = 1$ ):

Flat:  $d_{p,0} = \varpi$

Closed:  $d_{p,0} = \frac{1}{\sqrt{k}} \sin^{-1}(\varpi \sqrt{k})$

Open:  $d_{p,0} = \frac{1}{\sqrt{|k|}} \sinh^{-1}(\varpi \sqrt{|k|})$

### The particle horizon

Horizon distance = distance a photon has traveled since  $t = 0$ .

$$\int_{t_1}^{t_2} \frac{cdt}{R(t)} = - \int_{\varpi^2}^{\varpi^1} \frac{d\varpi}{\sqrt{1-k\varpi^2}} = \int_{\varpi^2}^{\varpi^1} \frac{d\varpi}{\sqrt{1-k\varpi^2}} \quad \leftarrow \text{For a photon}$$

$$d_p(t) = R(t) \int_0^{\varpi} \frac{d\varpi'}{\sqrt{1-k\varpi'^2}} = R(t) \int_{t_e}^{t_0} \frac{cdt'}{R(t')}$$

$$d_h(t) = R(t) \int_0^t \frac{cdt'}{R(t')}$$

Radiation dominated flat universe:  $R \propto t^{1/2} \rightarrow d_h(t) = 2ct$

Matter dominated flat universe:  $R \propto t^{2/3} \rightarrow d_h(t) = 3ct$

Matter dominated flat universe in terms of redshift  $\rightarrow d_h(z) = \frac{2c}{H_0 \sqrt{\Omega_{m,0}}} \frac{1}{(1+z)^{3/2}}$

Including  $\Omega_\Lambda \rightarrow d_h(t) = \left(\frac{\Omega_{m,0}}{\Omega_{\Lambda,0}}\right)^{1/3} \sinh^{2/3} \left(\frac{3}{2} H_0 t \sqrt{\Omega_{\Lambda,0}}\right) \int_0^t \frac{cdt'}{\left(\frac{\Omega_{m,0}}{\Omega_{\Lambda,0}}\right)^{1/3} \sinh^{2/3} \left(\frac{3}{2} H_0 t' \sqrt{\Omega_{\Lambda,0}}\right)}$   
 = 14.6 Gpc (WMAP) [29.158]

### The paths of photons in terms of proper distance.

$$(d\varpi)^2 = (cdt)^2 - R^2(t) \left[ \left(\frac{d\varpi}{\sqrt{1-k\varpi^2}}\right)^2 + (\varpi d\theta)^2 + (\varpi \sin\theta d\phi)^2 \right]$$

Matter dominated flat universe:

$$\int_0^t \frac{cdt'}{R(t')} = \int_{\varpi}^{\varpi_e} d\varpi'$$

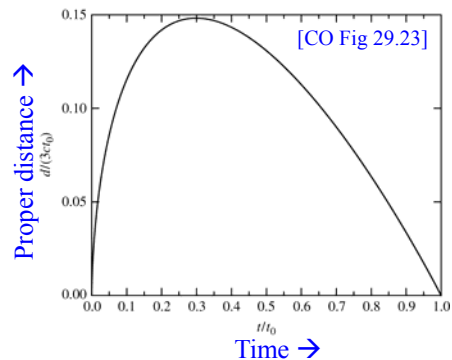
$$R(t) = \left(\frac{3}{2}\right)^{2/3} \left(\frac{t}{t_H}\right)^{2/3} = \left(\frac{t}{t_0}\right)^{2/3}$$

$$\varpi = \varpi_e - 3ct_0 \left(\frac{t}{t_0}\right)^{1/3}$$

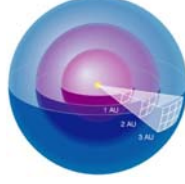
At  $t = t_0$ ,  $\varpi = 0 \rightarrow \varpi_e = 3ct_0$

Proper distance:

$$R(t)\varpi = d_p(t) = 3ct_0 \left[ \left(\frac{t}{t_0}\right)^{2/3} - \left(\frac{t}{t_0}\right) \right] \quad [29.165]$$



## Luminosity Distance



Redshift  $\rightarrow (1+z)$   
Time dilation  $\rightarrow (1+z)$

**In practice:**  
[CO pg. 1209]

$$F = \frac{L}{4\pi d^2}$$

$$F = \frac{L}{4\pi \varpi^2 (1+z)^2}$$

$$d_L = \varpi (1+z)$$

**We need a  $\varpi - z$  relation.**

In principle:

- R-W metric:

$$\int_{t_e}^{t_0} \frac{cdt}{R(t)} = \int_0^{\varpi_e} \frac{d\varpi}{\sqrt{1-k\varpi^2}}$$

$$\int_{R(t_e)}^{R(t_0)} \frac{c dR}{R \times dR/dt} = \int_0^{\varpi_e} \frac{d\varpi}{\sqrt{1-k\varpi^2}}$$

- Use Friedmann Eqn to find  $dR/dt$
- Then:  $R(t) = \frac{1}{1+z}$

**In practice:**

$$I(z) = H_0 \int_{\frac{1}{1+z}}^1 \frac{dR}{R(dR/dt)} = H_0 \int_0^z \frac{dz'}{H(z')}$$

$$I(z) = \int_0^z \frac{dz'}{\sqrt{\Omega_{m,0}(1+z')^3 + \Omega_{rel,0}(1+z')^4 + \Omega_{\Lambda,0} + (1-\Omega_0)(1+z')^2}}$$

$$d_{p,0}(z) = \frac{c}{H_0} I(z)$$

$$\varpi(z) = \frac{c}{H_0} S(z)$$

$$S(z) \equiv I(z) \quad (\Omega_0 = 1)$$

$$\equiv \frac{1}{\sqrt{\Omega_0 - 1}} \sin [I(z)\sqrt{\Omega_0 - 1}] \quad (\Omega_0 > 1)$$

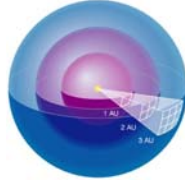
$$\equiv \frac{1}{\sqrt{1 - \Omega_0}} \sinh [I(z)\sqrt{1 - \Omega_0}] \quad (\Omega_0 < 1)$$

$$I(z) = \int_0^z \left\{ 1 - (1+q_0)z' + \left[ \frac{1}{2} + 2q_0 + \frac{3}{2}q_0^2 + \frac{1}{2}(1-\Omega_0) \right] z'^2 + \dots \right\} dz'$$

$$I(z) = z - \frac{1}{2}(1+q_0)z^2 + \left[ \frac{1}{6} + \frac{2}{3}q_0 + \frac{1}{2}q_0^2 + \frac{1}{6}(1-\Omega_0) \right] z^3 + \dots$$

$$\varpi \simeq \frac{cz}{H_0} \left[ 1 - \frac{1}{2}(1+q_0)z \right] \quad (\text{for } z \ll 1).$$

## Luminosity Distance



Redshift  $\rightarrow (1+z)$   
Time dilation  $\rightarrow (1+z)$

**In practice**  
(because of that @#% cosmological constant)

$$d_L(z) = \frac{c}{H_0} (1+z) S(z)$$

$$S(z) \equiv I(z) \quad (\Omega_0 = 1)$$

$$\equiv \frac{1}{\sqrt{\Omega_0 - 1}} \sin [I(z)\sqrt{\Omega_0 - 1}] \quad (\Omega_0 > 1)$$

$$\equiv \frac{1}{\sqrt{1 - \Omega_0}} \sinh [I(z)\sqrt{1 - \Omega_0}] \quad (\Omega_0 < 1)$$

**About right...**

From previous slide:

$$\varpi \simeq \frac{cz}{H_0} \left[ 1 - \frac{1}{2}(1+q_0)z \right] \quad (\text{for } z \ll 1).$$

$$d_L(z) \simeq \frac{cz}{H_0} \left[ 1 + \frac{1}{2}(1-q_0)z \right] \quad (\text{for } z \ll 1).$$

$$m - M = 5 \log_{10}(d_L/10 \text{ pc})$$

$$m - M \simeq 5 \log_{10} \left[ \frac{c}{(100 \text{ km s}^{-1} \text{ Mpc}^{-1})(10 \text{ pc})} \right] - 5 \log_{10}(h)$$

$$+ 5 \log_{10}(z) + 5 \log_{10} \left[ 1 + \frac{1}{2}(1-q_0)z \right] \quad (\text{for } z \ll 1)$$

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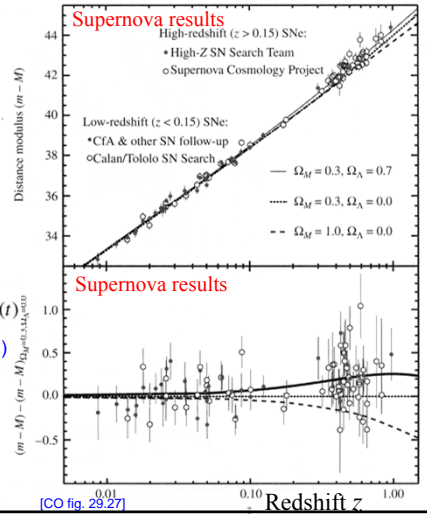
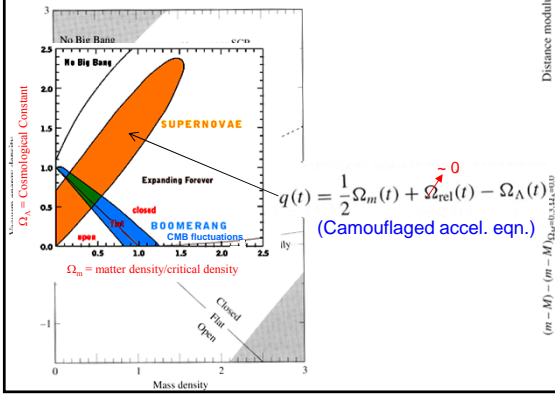
$$+ 5 \log_{10}(1+z) + 5 \log_{10}[S(z)]$$

$$= 42.38 - 5 \log_{10}(h) + 5 \log_{10}(1+z) + 5 \log_{10}[S(z)].$$

$$m - M \simeq 42.38 - 5 \log_{10}(h) + 5 \log_{10}(z) + 1.086(1-q_0)z \quad (\text{for } z \ll 1). \quad [\text{CO 29.188}]$$

## $q_0$ – the accelerating universe

- Type Ia supernovae are best standard candles.
  - Least scatter in luminosity
- 2 independent groups get same answer
  - High- $z$  Supernova Search
    - Garnavich et al. 1998, ApJ 509, 74
  - Supernova Cosmology Project
    - Perlmutter et al. 1999, ApJ 517, 565
- Found *acceleration*
  - Not deceleration as expected.



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