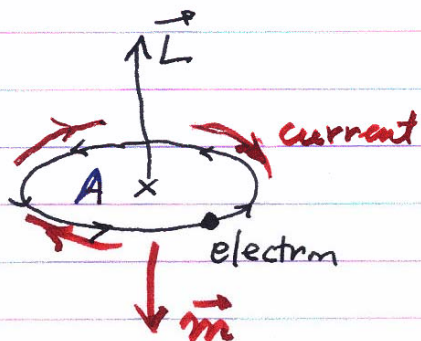


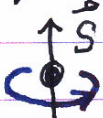
# Magnetization and Bound Current 9.2/1

Magnetic moment ( $\vec{m}$ ) and magnetization ( $\vec{M}$ )



$$\vec{m} = IA \hat{n}$$

$$\vec{m} = \frac{1}{2} \int \vec{r} \times \vec{j}(\vec{r}) d^3x$$

It's a bit more complicated, because an electron also has spin  and a

Corresponding spin magnetic moment 

$$\vec{m} = \frac{-e}{2m} (\vec{L} + 2\vec{S})$$

Check units:  $C \text{ kg m/s m} / \text{kg} = A$

## The essential idea for magnetism in matter

The magnetic field produced by a magnetized object is the same as that of a volume current density  $\vec{J}_{\text{Bound}}(\vec{x})$  and surface current density  $\vec{K}_{\text{Bound}}(\vec{x})$  given by

$$\vec{J}_B(\vec{x}) = \nabla \times \vec{M} \quad \text{and} \quad \vec{K}_B(\vec{x}) = \vec{M}(\vec{x}) \times \hat{n}$$

( $\vec{x}$  on the surface)

# Magnetization and Bound Current 9.2/2

Theorem A sample of matter with magnetization  $\vec{M}(\vec{x})$  has an associated magnetic field  $\vec{B}_{\text{Matter}}(\vec{x})$  which is the same as  $\vec{B}$  for a volume current density  $\vec{J}_{\text{Bound}}(\vec{x})$  and surface current density  $\vec{K}_{\text{Bound}}(\vec{x})$ ,

$$\vec{J}_{\text{B}}(\vec{x}) = \nabla \times \vec{M} \quad \text{and} \quad \vec{K}_{\text{B}}(\vec{x}) = \vec{M}(\vec{x}) \times \hat{n}$$

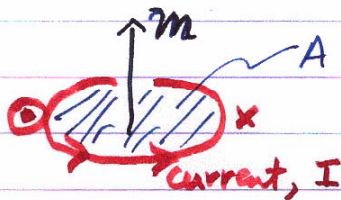
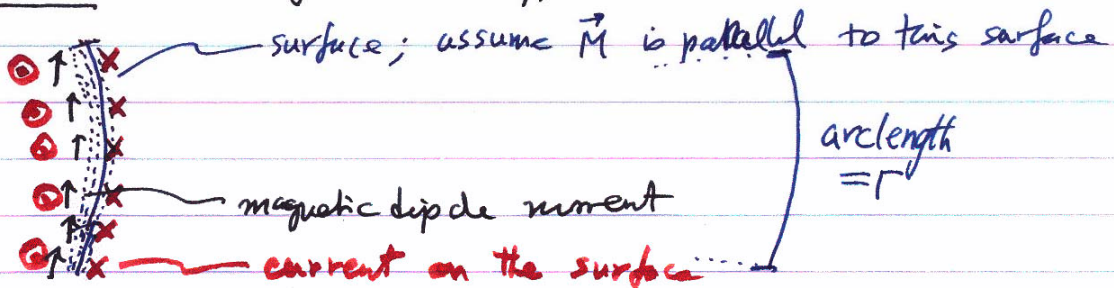
(in the volume) (on the surface)

Proof #1 See Equations 9-7 to 9-14.

*/ by integration /*

Proof #2 (more geometrically)

Surface



Surface current density  $K_B$

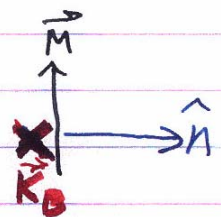
$$K_B \cdot \Gamma = \frac{\delta Q}{\delta t} = \frac{N \cdot I \delta t}{\delta t}$$

$$N = \# \text{ of atoms} = n \cdot (A\Gamma)$$

$$\text{Thus } K_B = \frac{n A \Gamma I}{\Gamma} = n I A = M$$

Direction of  $\vec{K}_B$  is same as  $\vec{M} \times \hat{n}$

$$\vec{K}_B = \vec{M} \times \hat{n} \quad \text{Q.E.D.}$$

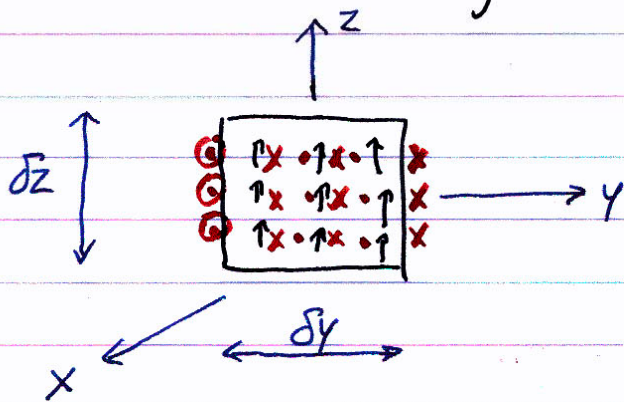


Check units:  $\frac{A}{m}$  and  $\frac{Am^2}{m^3}$  ✓

9.2/3

Volume

Consider a small subvolume inside the material. Set up a coordinate system with the z direction parallel to  $\vec{M}$



Current density  $\perp$  to the yz plane is  $J_{Bx} \hat{i}$

$$J_{Bx} \delta y \delta z = \frac{\delta Q}{\delta t}$$

The sign is tricky.

$$J_{Bx} \delta y \delta z = + \left( I_{\text{atom}} N \right)_{y + \frac{\delta y}{2}} - \left( I_{\text{atom}} N \right)_{y - \frac{\delta y}{2}}$$

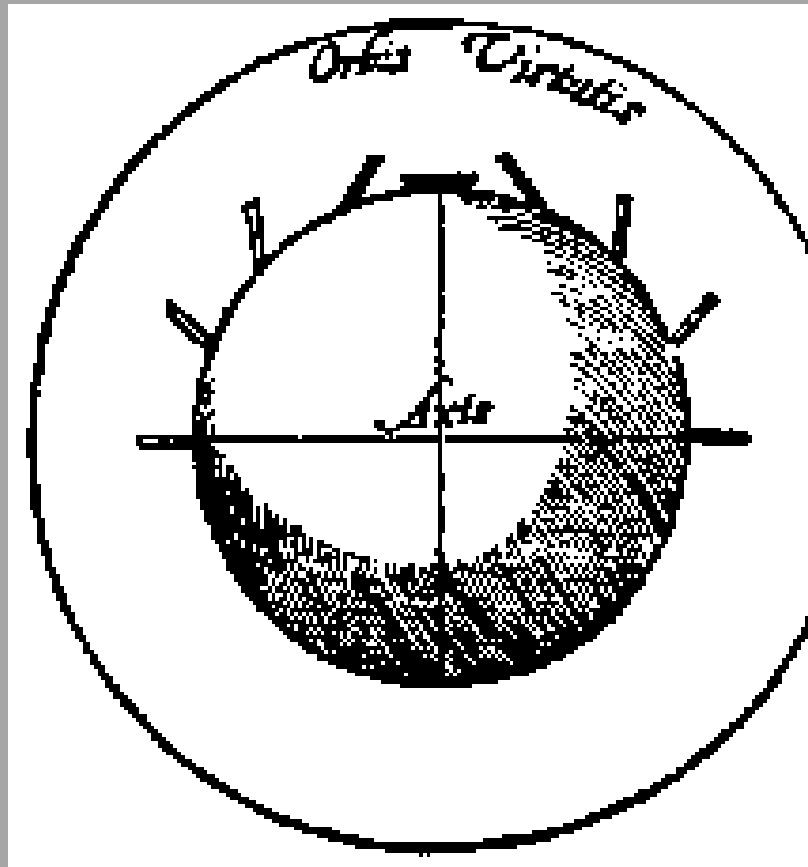
$$= (I n A \delta z)_{y + \frac{\delta y}{2}} - (I n A \delta z)_{y - \frac{\delta y}{2}}$$

$$= M_z \left( y + \frac{\delta y}{2} \right) \delta z - M_z \left( y - \frac{\delta y}{2} \right) \delta z$$

$$= \frac{\partial M_z}{\partial y} \delta y \delta z$$

$$J_{Bx} = \frac{\partial M_z}{\partial y} \quad (\text{assuming } \vec{M} \text{ is in the } z \text{ direction})$$

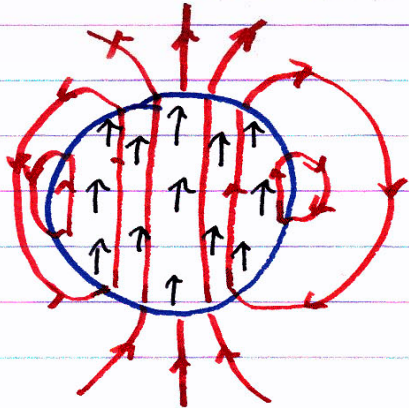
$$\vec{J}_B = \nabla \times \vec{M} \quad \underline{\underline{\text{Q.E.D.}}}$$



"Terrella" is Latin for "little Earth," the name given by William Gilbert to a magnetized sphere with which he demonstrated to Queen Elizabeth I his theory of the Earth's magnetism. By moving a small compass around the terrella and showing that it always pointed north-south, Gilbert argued that the same thing, on a vastly larger scale, was happening on Earth, and that the only reason why a compass pointed north-south

Later scientists such as Birkeland used the name "terrella" for magnetized spheres used inside vacuum chambers, together with electron beams, to study the motion of fast charged particles near the Earth. A sophisticated terrella experiment in a vacuum chamber is currently operated by Dr. Hafez u-Rahman at the University of California at Riverside.

Example Determine the magnetic field 9.2/84  
due to a uniformly magnetized sphere

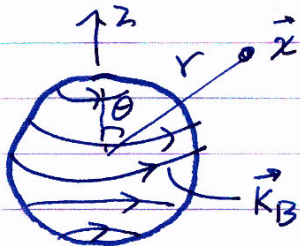


radius =  $a$

magnetization =  $M \hat{k}$

$$\vec{J}_B = \nabla \times \vec{M} = 0$$

$$\vec{K}_B = \vec{M} \times \hat{n} = M \hat{k} \times \hat{r} \\ = M \sin \theta \hat{\phi}$$



$$\vec{B} = \nabla \times \vec{A} \quad \text{where}$$

$$\vec{A}(\vec{x}) = \frac{\mu_0}{4\pi} \oint_S \frac{\vec{K}_B(\vec{x}')}{|\vec{x} - \vec{x}'|} dA'$$

$$\vec{A}(\vec{x}) = \frac{\mu_0 M}{4\pi} a^2 \hat{k} \times \oint_S \frac{\hat{r}'}{|\vec{x} - \vec{x}'|} d\Omega'$$

solid angle  
 $dA' = a^2 d\Omega'$

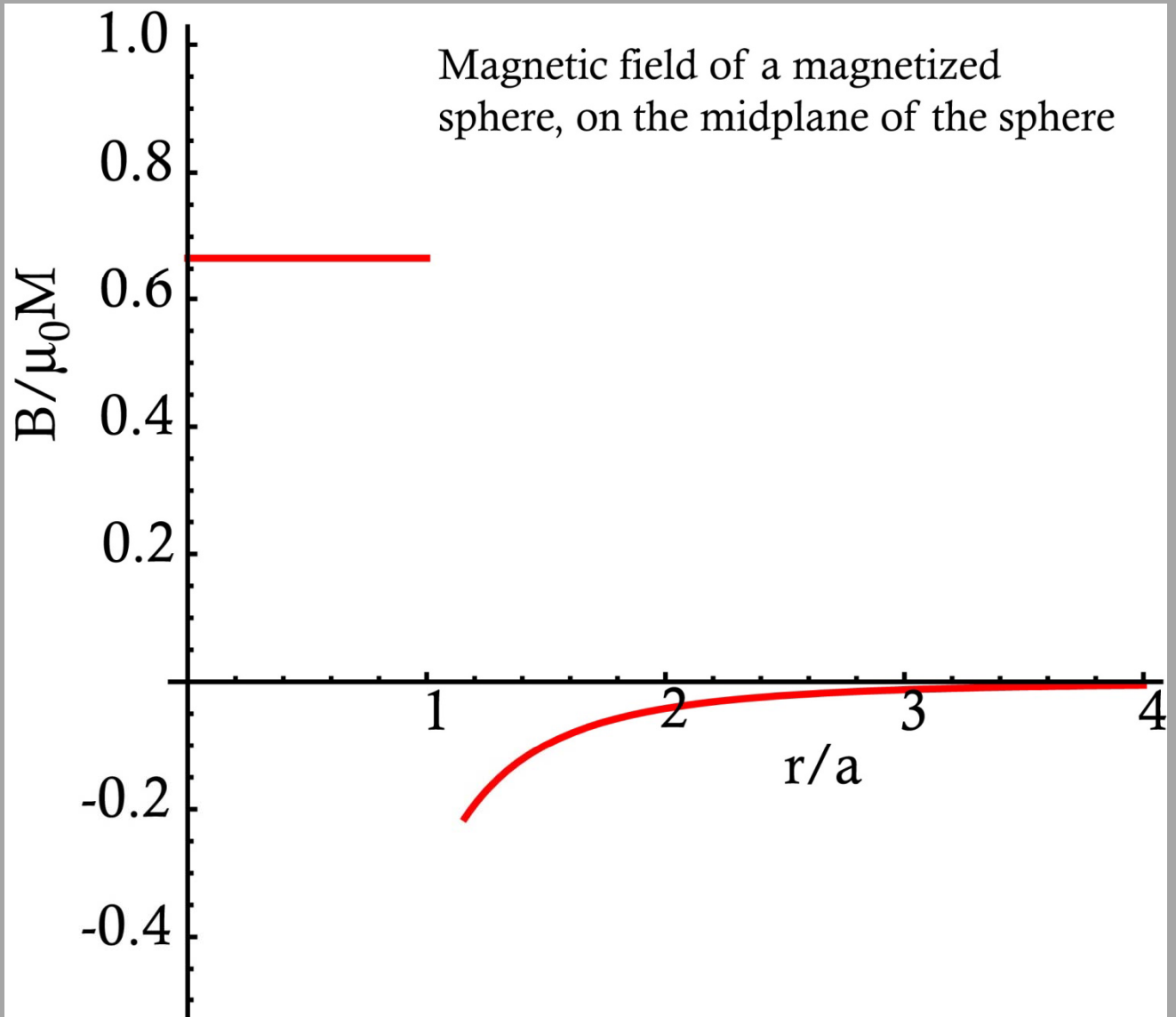
An interesting integral - see Eqs. 8.62 - 8.65

$$\int \frac{\hat{r}'}{|\vec{x} - \vec{x}'|} d\Omega' = \begin{cases} \frac{4\pi}{3} \frac{\vec{x}}{a^2} & \text{if } r < a \\ \frac{4\pi}{3} \frac{\vec{x}a}{r^3} & \text{if } r > a \end{cases}$$

$$\vec{A}(\vec{x}) = \begin{cases} \frac{1}{3} \mu_0 M \hat{k} \times \vec{x} & \text{if } r < a \text{ (inside)} \\ \frac{1}{3} \mu_0 M \frac{\hat{k} \times \vec{x} a^3}{r^3} & \text{if } r > a \text{ (outside)} \end{cases}$$

$$\vec{B}(\vec{x}) = \begin{cases} \frac{2}{3} \mu_0 M \hat{k} & \text{if } r < a \text{ (a uniform field)} \\ \frac{1}{3} \mu_0 M \frac{a^3}{r^3} [3\hat{r}(\hat{k} \cdot \hat{r}) - \hat{k}] & \text{if } r > a \end{cases}$$

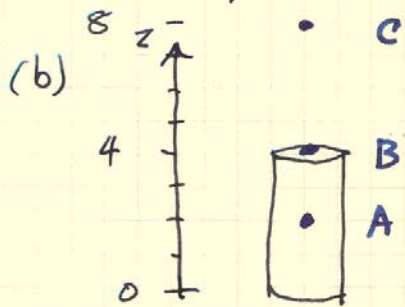
(a pure point dipole field)



## Homework

- (a) Determine the magnetic field of a uniformly magnetized cylinder (with radius  $a$  and height  $h$ ) for points on the axis of the cylinder.

Hint First, what is the bound surface current? Then use the Biot-Savart law to determine  $\vec{B}$ .



Suppose  $a=1$  and  $h=4$ .

(b) Evaluate  $\frac{B_z}{\mu_0 M}$  at point A

(c) Evaluate  $\frac{B_z}{\mu_0 M}$  at point B

(d) Evaluate  $\frac{B_z}{\mu_0 M}$  at point C

Show your work for all parts of the problem.

Demo: magnetometer  
and bar magnet