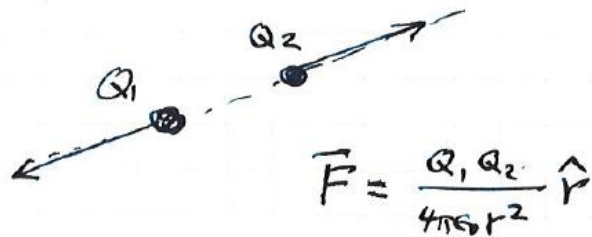


## Lecture 1 - Review



But this is not action at a distance.

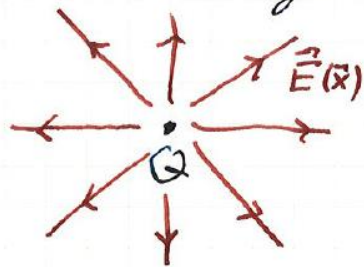
Every charge has an associated field, filling the space.

The field exerts forces on other charges.

## Field theory

$\vec{F} = q \vec{E}(\vec{r})$  ← the force on a test charge  $q$  at  $\vec{r}$

### Field theory



$$\vec{E}(\vec{r}) = \frac{Q \hat{r}}{4\pi\epsilon_0 r^2}$$

Units of  $\vec{E}$

N/C or V/m


Volt =  $\frac{\text{Joule}}{\text{Coulomb}}$

$$\frac{1}{4\pi\epsilon_0} = 8.99 \times 10^9 \frac{\text{Nm}^2}{\text{C}^2} \text{ (or, } \frac{\text{Vm}}{\text{C}})$$

The field equations of electrostatics

$$\nabla \cdot \vec{E} = \rho/\epsilon_0 \quad \text{and} \quad \nabla \times \vec{E} = 0.$$

Divergence and Curl



$$\vec{E}(\vec{x}) = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r} = E_r \hat{r}$$

$$\nabla \cdot \vec{E} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 E_r) = 0$$

But what about  $\nabla \cdot \vec{E} = \rho/\epsilon_0$ ?

There is a singularity at  $r=0$ .

We should have

$$\int_V \nabla \cdot \vec{E} d^3x = \oint_S \vec{E} \cdot d\vec{A} = \frac{Q}{4\pi\epsilon_0} \oint d\Omega = Q/\epsilon_0$$

Gauss's law

$$d\vec{A} = \hat{r} r^2 d\Omega$$

$$\text{Thus} \quad \nabla \cdot \vec{E} = \frac{Q}{\epsilon_0} \delta^3(\vec{x})$$

In general, then,

$$\nabla \cdot \vec{E} = \rho(\vec{x})/\epsilon_0 \quad (\text{superposition principle})$$

The curl

In spherical coordinates,

$$\nabla \times \vec{F} = \frac{1}{r^2 \sin \theta} \begin{vmatrix} \hat{r} & r\hat{\theta} & r\sin\theta\hat{\phi} \\ \partial/\partial r & \partial/\partial \theta & \partial/\partial \phi \\ F_r & rF_\theta & r\sin\theta F_\phi \end{vmatrix}$$

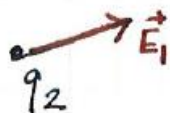
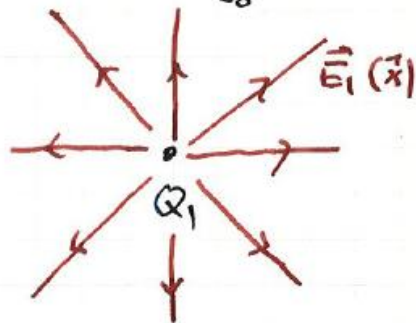
Now consider  $\vec{F} = \hat{r} F_r(r)$

Then

$$\nabla \times \vec{F} = \frac{1}{r^2 \sin \theta} \begin{vmatrix} \hat{r} & r\hat{\theta} & r\sin\theta\hat{\phi} \\ \partial/\partial r & \partial/\partial \theta & \partial/\partial \phi \\ F_r & 0 & 0 \end{vmatrix} = 0$$

## Electrostatics

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \text{and} \quad \nabla \times \vec{E} = 0$$



$$\vec{E}_1 = \frac{Q_1}{4\pi\epsilon_0 r^2} \hat{r}$$

$$\vec{F}_2 = q_2 \vec{E}_1 = \frac{Q_1 q_2}{4\pi\epsilon_0 r^2} \hat{r}$$

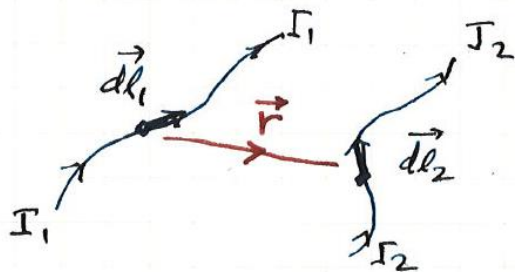
Field theory versus action at a distance

↳ Faraday;  
Maxwell

↓  
obsolete after  
Faraday and Maxwell

## Magnetostatics

↳ Static magnetic systems



The force on  $d\vec{l}_2$  due to the current in  $d\vec{l}_1$  is

$$\delta \vec{F} = I_2 d\vec{l}_2 \times d\vec{B}_1 \quad (\text{field theory})$$

where

$$d\vec{B}_1 = \frac{\mu_0}{4\pi} \frac{I_1 d\vec{l}_1 \times \hat{r}}{r^2} \quad (\text{field theory})$$

So

$$\delta \vec{F} = \frac{\mu_0 I_1 I_2}{4\pi r^2} d\vec{l}_2 \times (d\vec{l}_1 \times \hat{r})$$

The force on the complete

wire is

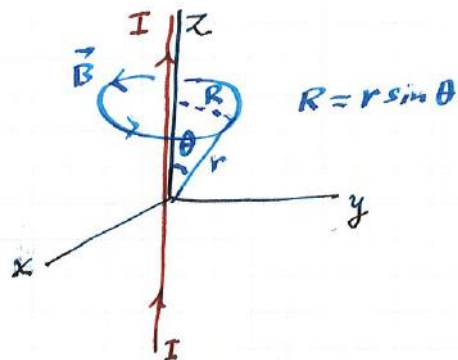
$$\vec{F} = \iint \frac{\mu_0 I_1 I_2}{4\pi r^2} d\vec{l}_2 \times (d\vec{l}_1 \times \hat{r})$$

This force is not action at a distance; it is a field effect.

$$\mu_0 = 4\pi \times 10^{-7} \frac{\text{Tm}}{\text{A}}$$

### Example

a long straight wire  
carrying current  $I$



$$\vec{B}(\vec{x}) = \frac{\mu_0 I}{2\pi R} \hat{\phi}$$

$$\nabla \times \vec{B} = \frac{1}{r^2 \sin \theta} \begin{vmatrix} \hat{r} & r\hat{\theta} & r\sin\theta\hat{\phi} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ 0 & 0 & r\sin\theta \frac{\mu_0 I}{2\pi r\sin\theta} \end{vmatrix}$$

$$\nabla \times \vec{B} = 0$$

Why not  $\nabla \times \vec{B} = \mu_0 \vec{J}$ ?

Again, there is a singularity,  
at  $R=0$ .

We should have Stokes thm  
 $d\vec{l} = \hat{\phi} d\ell$

$$\int_{\text{Disk}} (\nabla \times \vec{B}) \cdot d\vec{A} = \oint_{\text{circle}} \vec{B} \cdot d\vec{l}$$

$$= \frac{\mu_0 I}{2\pi R} \cdot 2\pi R = \mu_0 I$$

$$\text{Thus } \nabla \times \vec{B} = \mu_0 I \delta(x) \delta(y) \hat{k}$$

Or, in general,

$$\nabla \times \vec{B} = \mu_0 \vec{J}(\vec{x})$$

current density  
(current per unit area)

Check units:

$$\frac{1}{\mu} T = \frac{T_m}{A} \frac{A}{m^2} \quad \checkmark$$

Exercise: Prove  $\nabla \cdot \vec{B} = 0$ .



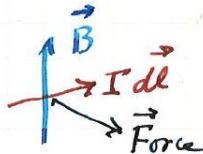
## Magnetostatics

- The field equations of magnetostatics are

$$\nabla \times \vec{B} = \mu_0 \vec{J} \quad \text{and} \quad \nabla \cdot \vec{B} = 0$$

- The force on a current element is

$$d\vec{F} = I d\vec{l} \times \vec{B}$$



- Units  $N = A m T$

$$T = \frac{N}{A m} = \frac{J}{A m^2} = \frac{V C}{C s m^2} = \frac{V s}{m^2}$$

$$\underline{1 T = 1 \frac{Vs}{m^2}}$$

- The force on a charged particle is

$$\vec{F} = q \vec{v} \times \vec{B}$$

$$(\text{Units: } C \frac{m}{s} T = A m T = N \checkmark)$$

END of the REVIEW