

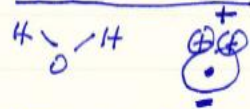
Electrostatics and Dielectrics

(Chapter 6) — a quick summary because you studied this in PHY 481.

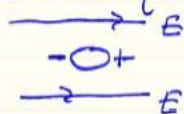
Polarization $\vec{P}(\vec{x})$

All matter is atomic (or molecular) and atoms and molecules contain charged particles. The net charge is 0, but there may be polarization.

Polar molecule



The dipole moment \vec{p} points from - to +;
in \vec{E} , \vec{p} experiences a torque toward alignment



Define "dipole moment"

- For 2 charges $\vec{p} = q \vec{\delta}$

- For a continuous distribution

$$\vec{p} = \int \vec{x} \rho(\vec{x}) d^3x$$

In the equilibrium state, \vec{p} aligns with the local field; $\langle \vec{p} \rangle = \alpha \vec{E}(\vec{x})$

average moment is a neighborhood of \vec{x} .

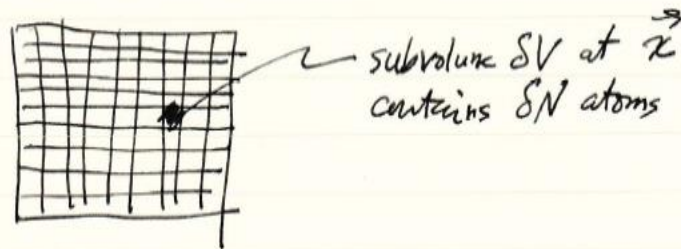
The Polarization Field $\vec{P}(\vec{x})$

Consider a sample of dielectric material.
Subdivide the sample into many small subvolumes.
But atoms are so tiny that each subvolume contains many atoms.

$$\text{Define } \vec{P}(\vec{x}) = \frac{1}{\delta V} \sum_{i=1}^{\delta N} \vec{p}_i ;$$

$$\text{or } \vec{P}(\vec{x}) = \frac{\delta N}{\delta V} \langle \vec{p} \rangle ;$$

i.e., $\vec{P}(\vec{x}) = \text{polarisation} = \text{dipole moment density}$



"Free charge" and "Bound charge"

$$\vec{\rho}(\vec{x}) = \rho_{\text{Free}}(\vec{x}) + \rho_{\text{Bound}}(\vec{x})$$

Charge that has
been added to the
dielectric

charge that belongs
to the atoms of the
dielectric

We can relate $\rho_{\text{Bound}}(\vec{x})$ to $\vec{P}(\vec{x})$.

Theorem

$$\rho_B(\vec{x}) = -\nabla \cdot \vec{P} \quad \text{for } \vec{x} \text{ inside the dielectric}$$

$$\sigma_B(\vec{x}) = \hat{n} \cdot \vec{P} \quad \text{for } \vec{x} \text{ on the surface of the dielectric}$$

Proof

Recall the electrostatic potential $V_a(\vec{x})$ of an electric dipole ($\vec{E} = -\nabla V_a$).

$$V(\vec{x}) = \int_{\text{sample}} \frac{d\vec{p} \cdot (\vec{x} - \vec{x}')}{4\pi\epsilon_0 |\vec{x} - \vec{x}'|^3} \quad \text{from the bound charges}$$

$$d\vec{p} = \vec{P}(\vec{x}') d^3x'$$

$$(\vec{x} - \vec{x}') / |\vec{x} - \vec{x}'|^3 = \nabla' \frac{1}{|\vec{x} - \vec{x}'|}$$

now integrate by parts

$$\vec{P} \cdot \nabla' \frac{1}{R} = \nabla' \cdot \left(\frac{\vec{P}}{R} \right) - \frac{\nabla' \cdot \vec{P}}{R}$$

$$V(\vec{x}) = \oint_{\text{surface}} \frac{\hat{n}' \cdot \vec{P}(\vec{x}')}{|\vec{x} - \vec{x}'|} dA' + \int_{\text{sample}} \frac{-\nabla' \cdot \vec{P}(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x'$$

$$= \oint_S \frac{\sigma_B dA'}{|\vec{x} - \vec{x}'|} + \int_V \frac{\rho_B(\vec{x}') d^3x'}{|\vec{x} - \vec{x}'|} \quad \text{4}\pi\epsilon_0$$

I.e., V is exactly the same as for surface charge density $\sigma_B = \hat{n}' \cdot \vec{P}(\vec{x}')$ and volume charge density $\rho_B = -\nabla' \cdot \vec{P}(\vec{x}')$

Q.E.D.

The Displacement Field $\vec{D}(\vec{r})$

Define $\vec{D} = \epsilon_0 \vec{E} + \vec{P}$,

Then

$$\nabla \cdot \vec{D} = \epsilon_0 \nabla \cdot \vec{E} + \nabla \cdot \vec{P}$$

$$= \rho - \rho_B$$

$$= \rho_F$$

$$\nabla \cdot \vec{D} = \rho_{Free} \quad \text{and} \quad \vec{D} = \epsilon_0 \vec{E} + \vec{P}.$$

Consequences

- Gauss's Law

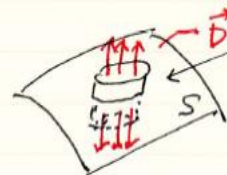
$$\oint_S \vec{D} \cdot d\vec{A} = Q_{Free} \quad / \text{by Gauss's theorem}/$$

- Boundary Conditions (important!)

D_n is continuous ~~across~~ a surface

if $\sigma_{Free} = 0$; otherwise

$$D_n^{(+)} - D_n^{(-)} = \sigma_{Free}.$$



Small cylindrical volume in limit height $\rightarrow 0$

$$\oint \vec{D} \cdot d\vec{A} = D_n^{(+)} A - D_n^{(-)} A = Q_{Free} \\ = \sigma_{Free} A$$

$$\therefore D_n^{(+)} - D_n^{(-)} = \sigma_{Free}$$

Isotropic Linear Dielectrics

Many isotropic insulators are "linear dielectrics"; i.e., there is a linear relationship between polarization and the electric field

$$\vec{P} = \epsilon_0 \chi_e \vec{E} \quad \langle \vec{P} \rangle \propto \vec{E}(\vec{x})$$

Glass, water, gases and liquids, etc.

Material parameters for linear dielectrics

- χ_e electric susceptibility
 $\vec{P} = \chi_e \epsilon_0 \vec{E}$ ← Note: $\chi_e > 0$
- ϵ permittivity
 $\vec{D} = \epsilon \vec{E}$ ← Note: $\epsilon > \epsilon_0$

Note that $\epsilon = \epsilon_0(1 + \chi_e)$ because
 $\vec{D} = \epsilon_0 \vec{E} + \vec{P} = (1 + \chi_e) \epsilon_0 \vec{E} = \epsilon \vec{E}$

- K dielectric constant
 $K = \epsilon / \epsilon_0$ ← Note: $K > 1$

$$K = \frac{\epsilon}{\epsilon_0} = 1 + \chi_e$$

Any of these material constants can be used to specify the dielectric properties of the material.