<u>Maxwell's Equations in Matter</u>
Start with the fundamental equations
(1) $\nabla \cdot \mathbf{E} = \rho / \varepsilon_{\theta}$
(2) $\nabla \cdot \mathbf{B} = 0$
(3) $\nabla \times \mathbf{E} = -\partial \mathbf{B}/\partial \mathbf{t}$
(4) $\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \varepsilon_0 \partial \mathbf{E} / \partial \mathbf{t}$
where $\rho(\mathbf{x}, \mathbf{t})$ and $\mathbf{J}(\mathbf{x}, \mathbf{t})$ satisfy the
continuity equation,
$\nabla \cdot J = -\partial \rho / \partial t$

(conservation of charge)

Now consider the effects of macroscopic matter; i.e., matter with many atoms (6 x 10^{23} per mole).

We'll treat the charge and current in atoms by averaging --- a good approximation because N is so large;

error ~ 1/ \sqrt{N} ~ 1/ $\sqrt{(6x10^{23})}$ ~ 1/ 10^{12}

Two of Maxwell's equations do not depend on matter: (2) $\nabla \cdot B = 0$ (primary magnetic monopoles do not exist) (3) $\nabla \times E = -\partial B/\partial t$ (EM induction is a field effect)

The other two depend on sources, i.e., matter, for which we'll make macroscopic approximations.

<u>Charge density</u>

$$\rho_{\text{micro}}(\mathbf{x}, \mathbf{t}) = \sum e_i \delta^3(\mathbf{x} - \mathbf{x}_i(\mathbf{t}))$$

i=1..N
or, in quantum theory,

$$\rho_{\text{micro}}(\mathbf{x},t) = \sum_{i=1..N} e_i |\psi_i(\mathbf{x},t)|^2$$

where N ~ 6 x 10^{23} Now define the macroscopic average,

 $\rho(\mathbf{x},t) = (1/\delta V) \int_{\delta V} \rho_{\text{micro}}(\mathbf{x}^{\prime},t) d^{3}x^{\prime}$

where δV is a volume around the point
x, both *Large* and *small* :
 -- *Large* compared to an atom
 -- *small* compared to the full system

Recall the methods for dielectric materials.

$$\rho(\mathbf{x}, \mathbf{t}) = \rho_{\text{free}}(\mathbf{x}, \mathbf{t}) + \rho_{\text{bound}}(\mathbf{x}, \mathbf{t})$$

free charge is charge that is not bound in atoms or molecules; bound charge is charge that is confined to an atom or molecule.

Review:

$$\rho_{\text{bound}}(\mathbf{x}, \mathbf{t}) = - \nabla \cdot \mathbf{P}$$

P = polarization = electric dipole
moment density

$$P(x,t) = n(x,t) \langle p \rangle$$

Gauss's Law for the electric field E(x,t)

$$\nabla \cdot \mathbf{E} = \rho_{\text{free}} / \varepsilon_0 + \rho_{\text{bound}} / \varepsilon_0 = \rho_{\text{free}} / \varepsilon_0 - \nabla \cdot \mathbf{P} / \varepsilon_0$$

 $\nabla \cdot (\varepsilon_0 \mathbf{E} + \mathbf{P}) = \rho_{\text{free}}$

Define the displacement field $\mathbf{D} = \varepsilon_0 \mathbf{E} + \mathbf{P}$; then $\nabla \cdot \mathbf{D} = \rho_{\text{free}}$.

But we will need a "constitutive equation" to relate E(x, t) and D(x,t).

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For linear dielectrics,
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 $D = \varepsilon E \quad or \quad P = \varepsilon_0 \chi_e E$ (permittivity or susceptibility)

All of this is familiar for static fields. It is a good enough approximation for time dependent fields, unless photon energy > ~keV. However, even for low photon energies, ε may depend on frequency.



Electric current density The fundamental field equation is $\nabla \times B = \mu_0 J + \mu_0 \varepsilon_0 \partial E / \partial t$ Now, what is current density J(x,t)?

Area δA



Any process that moves charge contributes to the current density.

Again, we separate free current and bound current.

In Chapter 9 we considered bound current that is related to the magnetization M(x,t),

 $J_{magn.} = \nabla \times M$; $M(x,t) = n(x,t) \langle m \rangle$

This is current inside an atom (or molecule) which gives the atom a magnetic dipole moment m.

But now, for time-dependent systems, there is another kind of bound current related to the polarization **P(x,t)**,

 $\mathbf{J}_{\mathbf{pol.}} = \partial \mathbf{P} / \partial t; \quad \mathbf{P}(\mathbf{x},t) = \mathbf{n}(\mathbf{x},t) \langle \mathbf{p} \rangle$

Result: $J = J_{free} + \nabla \times M + \partial P/\partial t$

Result: $J = J_{free} + \nabla \times M + \partial P/\partial t$ **Theorem.** $\nabla \cdot J = -\partial \rho / \partial t$ (continuity equation) **Proof.**

$$\nabla \cdot J = \nabla \cdot J_{\text{free}} + \nabla \cdot (\nabla \times M) + \nabla \cdot (\partial P / \partial t)$$
$$= - \partial \rho_{\text{free}} / \partial t + 0 + \partial (-\rho_{\text{bound}}) / \partial t$$
$$= - \partial \rho / \partial t$$

Q.E.D.

The Ampere-Maxwell equation in the presence of matter. $\nabla \times \mathbf{B} = \mu_{a}\mathbf{J} + \mu_{a}\varepsilon_{a} \partial \mathbf{E}/\partial \mathbf{t}$ = $\mu_{a}(J_{free} + \nabla \times M + \partial P/\partial t) + \mu_{a}\varepsilon_{a} \partial E/\partial t$ $\nabla \times (B/\mu_{a}-M) = J_{free} + \partial(\varepsilon_{a}E + P) /\partial t$ (4) $\nabla \times H = J_{\text{free}} + \partial D / \partial t$ where $H = B/\mu_a - M$ and $D = \varepsilon_a E + P$ displacement current = $\partial \mathbf{D} / \partial t$ Constitutive equations for linear magnetic materials (diamagnetic or paramagnetic):

$$\mathbf{B} = \mu \mathbf{H}$$
 or $\mathbf{M} = \mu_{\varrho} \chi_{m} \mathbf{B}$

<u>Maxwell's Equations</u>, including the atomic effects of magnetization and polarization in macroscopic matter...

 $\nabla \cdot \mathbf{D} = \rho_{\text{free}}$ (1) $\nabla \cdot \mathbf{B} = 0$ (2) (3) $\nabla \times \mathbf{E} = -\partial \mathbf{B}/\partial t$ $\nabla \times H = J_{\text{free}} + \partial D / \partial t$ (4) $\mathbf{D} = \varepsilon_{\alpha}\mathbf{E} + \mathbf{P} = \varepsilon \mathbf{E}$ $\mathbf{B} = \mu_{\alpha} (\mathbf{H} + \mathbf{M}) = \mu \mathbf{H}$

<u>Consequences for energy density and</u> <u>energy flux (see Section 11.X)</u>

The energy density, which includes both field energy and material energies, is u = ½ (E·D + B·H)

The energy flux is S = E × H

Exercise 11.X. The continuity equation for energy conservation is $\nabla \cdot S = -\partial B / \partial t - J_{free} \cdot E$