

Chapter 5. Photons: Covariant Theory

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Section 6.1

Natural dimensions and units

Set $\hbar = 1$ and $c = 1$.

I.e., omit all factors of $\hbar$ and c during calculations.

At the end of the calculation, restore the factors of $\hbar$ and c by unit analysis.
(There is always a unique way to do this.)

Read Section 6.1 for a detailed discussion of this method.

Section 6.2

The S-matrix expansion

We already know this from PHY 855.

Use the Interaction Picture to calculate the transition matrix elements, treating the interaction in perturbation theory.

(1) The Hamiltonian is $H = H_0 + H_I$.

(2) The S-matrix, $S = U(\infty, -\infty)$.

The goal of quantum theory is to calculate *time evolution* of the states of the system.

Let $|\psi, t_0\rangle$ be the state at time t_0 .
Then the state at time t is (in the Schrodinger picture)

$$|\psi, t\rangle = e^{-iH(t-t_0)} |\psi, t_0\rangle$$

Now suppose $t_0 \rightarrow -\infty$,
and $|\psi, t_0\rangle \rightarrow |i\rangle$,
where $|i\rangle$ consists of free particles
(i.e., very far apart).

Then the state at time t , for $t \rightarrow \infty$, is

$$\begin{aligned} |\psi, t\rangle &= e^{-iH(t-t_0)} |i\rangle \\ &= \sum_j |j\rangle \langle j| e^{-iH(t-t_0)} |i\rangle \\ &= \sum_j |j\rangle S_{ji} \end{aligned}$$

So the probability that the state $|f\rangle$ will be observed at time t is

$$P(i \rightarrow f) = |S_{fi}|^2.$$

Now, consider, in the Heisenberg picture,

$$\langle \Psi_0 | A(t) B(t_0) | \Psi_0 \rangle$$

$$= \langle \Psi_0 | A(0) e^{-iHt} e^{iHt_0} B(0) | \Psi_0 \rangle$$

$$= \langle \Psi_0 | A(0) U(t, t_0) B(0) | \Psi_0 \rangle ;$$

...which must be the same function in the interaction picture

$$= \langle 0 | A(0)$$

$$\int_{t_0}^t T \exp [-i H_I(t') t'] dt' B(0) | 0 \rangle_{IP}$$

Thus,

$$U(\infty, -\infty) = S =$$

$$\int_{t_0}^t T \exp [-i H_I(t') t'] dt'$$

(3)

$\Rightarrow$ MANDL&SHAW Eq. (6.22b)

$$S = \sum_{n=0}^{\infty} (-i)^n / n! \int dt_1 \dots \int dt_n T\{ H_I(t_1) \dots H_I(t_n) \} .$$

Now apply this result to Quantum Field Theory:

$$H_I(t) = \int d^3x \mathcal{H}_I(t, \mathbf{x})$$

$$S = \sum_{n=0}^{\infty} (-i)^n / n! \int d^4x_1 \dots \int d^4x_n T\{ \mathcal{H}_I(\mathbf{x}_1) \dots \mathcal{H}_I(\mathbf{x}_n) \} .$$

Eq. (6.23)

“This equation is the Dyson expansion of the S-matrix. It forms the starting point for the approach to perturbation theory used in this book.”

$$\text{Usually } \mathcal{H}_I(\mathbf{x}) = -\mathcal{L}_I(\mathbf{x})$$

Section 6.3.

Wick's theorem

We studied this in PHY 855.

① The time-ordered product of two fields is equal to the normal-ordered product plus a c-number contraction.

$$T\{A(x_1) B(x_2)\} =$$

$$N\{A(x_1) B(x_2)\} + \underbrace{A(x_1) B(x_2)}$$

② The c-number contraction is a propagator.

$$\underbrace{A(x_1) B(x_2)} \equiv \langle 0 | T \{ A(x_1) B(x_2) \} | 0 \rangle$$

⊗ Examples

Real scalar field (6.32a) :

$$\underbrace{\phi(x_1)\phi(x_2)} = i \Delta_F(x_1-x_2)$$

Complex scalar field (6.32b):

$$\underbrace{\phi(x_1)\phi^\dagger(x_2)} = \underbrace{\phi^\dagger(x_2)\phi(x_1)} = i \Delta_F(x_1-x_2)$$

Dirac field (6.32c):

$$\underbrace{\psi(x_1)\bar{\psi}(x_2)} = -\underbrace{\bar{\psi}(x_2)\psi(x_1)} = i S_F(x_1-x_2)$$

Electromagnetic field (6.32d):

$$\underbrace{A_\mu(x_1) A_\nu(x_2)} = (\text{a future lecture})$$

The contraction of two distinguishable fields is 0. *E.g.*, contraction of electron field and quark field = 0.

④ The time-ordered product of *any number of fields* can be written as the sum of all normal ordered products multiplied by c-number contractions.
(Wick's Theorem)

⑤ The vacuum expectation value of any normal-ordered product is 0.

⑥ The vacuum expectation value of any time-ordered product is the sum of all complete contractions.

 That's where we get Feynman diagrams.

Exercise: Read Chapter 6.

Homework Problem due April 1.

Problem 4.

Use anticommutation relations and definitions to prove

$$T\{\psi_{\alpha}(x) \bar{\psi}_{\beta}(y)\} - N\{\psi_{\alpha}(x) \bar{\psi}_{\beta}(y)\} = i S_F(x-y)_{\alpha\beta}$$