6.3. Wick's theorem Chanter 7 Feynman Diagrams and Rules in OFD.

Chapter 7. Feynman Diagrams and Rules in QED 7.1. Feynman diagrams in configuration space 7.2. Feynman diagrams in momentum space

7.3. Feynman rules for QED7.4. Leptons

8.2. Spin sums
8.3. Photon polarization sums

8.4-7. Examples 8.8-9. Bremsstrahlung Example. ""Pion nucleon scattering""
... but ignoring isospin.
(a purely academic exercise)

Fields: $\psi_{\alpha}(x)$ and $\varphi(x)$

Lagrangian density:

$$\mathcal{L} = \mathcal{L}_{free \, \phi} + \mathcal{L}_{free \, \psi} + \mathcal{L}_{interaction}$$

$$L_{free \, \varphi} = you \, know$$

$$\mathcal{L}_{free\,\psi} = you \, know$$

$$\mathcal{L}_{interaction} = g \overline{\psi} \psi \phi$$

Note that \mathcal{L} is relativistically invariant.

INTERACTING FIELDS

In Chapter 7 we will derive the methods for calculating transition matrix elements (Feynman diagrams) and reaction cross sections.

But today we'll do an example, to illustrate:

- computational techniques
- theoretical issues

Here is the example

$$\mathcal{L} = \mathcal{L}_{\gamma} + \mathcal{L}_{\rho} + \mathcal{L}_{int}$$

$$\mathcal{L}_{\gamma} = \overline{\gamma} (i8 - M) \gamma$$

$$\mathcal{L}_{\rho} = \frac{1}{2} \partial^{\mu} \partial_{\mu} \phi - \frac{M^{2}}{2} \phi$$

$$\mathcal{L}_{int} = g \phi \overline{\gamma} \gamma$$
No isospin: γ, M, γ''

$$\phi, M, \gamma'' \overline{\gamma}$$

Recall the transition amplitude

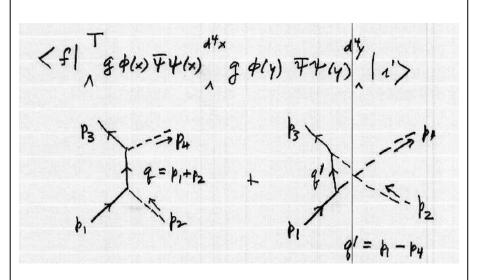
= the probability amplitude for evolution from initial state |i> at time −T to final state |f> at time +T; take the limit T → infinity.

The interaction Hamiltonian is

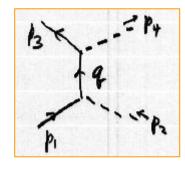
$$H_{I} = -\int dint d^{3}x$$

For the lowest order calculation, the relevant term in the perturbation expansion is $L_{\rm int}^{\ \ 2}$;

there are two Feynman diagrams.



The scattering amplitude



x(ofter factors)

S = g2 84 (P3+P4-P1-P2)

The cross section for polarized scuttering

$$d\sigma = (2\pi)^4 \delta^4 (P_4 - P_1) \sqrt{2m_2}$$

$$\frac{1}{4E_1E_2} \sqrt{164} \frac{1}{4} \frac{4^3 f_1}{(2\pi)^3 2E_4} |m|^2 ;$$
for un planized scattering,
$$d\sigma_M = \frac{1}{2} \sum_{s_3} \sum_{s_3} d\sigma$$

$$|M|^{2} = M^{*}M$$

$$= g^{4} u_{1}^{+}(\mathcal{Z}^{+}+A)(8^{0})u_{3}u_{3}(\mathcal{Z}+A)u_{1}$$

$$8^{0}\mathcal{Z}^{+} 8^{0} = 8^{0}(8^{0}\mathcal{Z}^{-}+8^{0}\mathcal{Z}^{-})8^{0}$$

$$= 8^{0}(8^{0}\mathcal{Z}^{0} + 8^{0}\mathcal{Z}^{0})8^{0}$$

$$= 8^{0}\mathcal{Z}^{0} - 8^{0}\mathcal{Z}^{0} = \mathcal{Z}^{0}$$

$$= 8^{0}\mathcal{Z}^{0} - 8^{0}\mathcal{Z}^{0} = \mathcal{Z}^{0} = \mathcal{Z}^{0}$$

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$$\sum_{S_{1}}^{S_{2}} u_{1} \overline{u_{1}} = \frac{K_{1} + M_{1}}{2M}$$

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(5) TRACE IDENTITIES (Appendix A)

(5) MANDELSTAM VARIABLES (see the handout)

$$|M|^{2} = \frac{g^{4}}{2M^{2}} \left[2Q_{0}p_{1}Q_{0}p_{3} - Q_{0}^{2}p_{1}p_{2} + Q_{0}^{2}M^{2} + AM\left(2Q_{0}p_{1} + 2Q_{0}p_{3}\right) + A^{2}\left(p_{1}p_{3} + M^{2}\right) \right]$$

Recall,
$$Q^{1} = \frac{p_{1}+p_{1}}{s-M^{2}} - \frac{p_{1}-p_{1}}{u-M^{2}} \text{ and } A = \frac{M}{s-M^{2}} - \frac{M}{u-M^{2}}$$

$$p_{1}+p_{2} = p_{3}^{2}+p_{4}^{2}$$

$$S = (p_{1}+p_{2})^{2} = M^{2}+m^{2}+2p_{1}p_{2}$$

$$t = (p_{1}-p_{3})^{2} = M^{2}+M^{2}-2p_{1}p_{3}$$

$$u = (p_{1}-p_{4})^{2} = M^{2}+m^{2}-2p_{1}p_{3}$$

Now we need Mathematica.

Some data on *low energy* π^+ - proton scattering (Homework Problem 6)

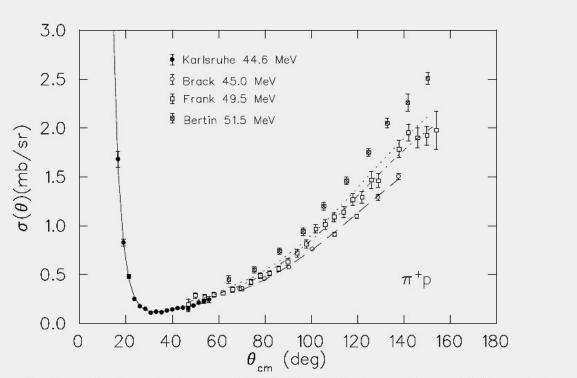


FIG. 1. Cross sections from π^+ proton scattering around 50 MeV. The Bertin data would agree with the prediction given by the dotted curve if they were consistent with the other data sets. The solid, long dash and dash-dot curves come from our fit and correspond to the Karlsruhe, Brack and Frank data sets.

π+ - proton
scattering from
the Particle
Data Group

The cross section has a peak at \sqrt{s} ~ 1200 MeV .

Why?

