

## Chapter 3: Momentum and Angular Momentum

- 3.1 Conservation of Momentum
- 3.2 Rockets
- 3.3 The center of mass
- 3.4 Angular momentum for a single particle
- 3.5 Angular momentum for several particles

**Read Chapter 3 during the next two weeks.**

Homework Assignment #5  
due in class Wednesday, October 4

[21] Problem 3.4 \*\*

[22] Problem 3.5 \*\*

[23] Problem 3.6 \*

[24] Problem 3.10 \*

[25] Problem 3.12 \*\*

[26] Problem 3.13 \*\*

*Use the cover sheet!*

## Section 3.1 Conservation of Momentum

First, define momentum for a single particle,

$$\mathbf{p} = m \mathbf{v} .$$

Now consider a system containing N particles. For each particle there is momentum,

$$\mathbf{p}_\alpha = m_\alpha \mathbf{v}_\alpha .$$

(  $\alpha = 1 \ 2 \ 3 \ \dots \ N$  )

The total momentum of the system is  $\mathbf{P}$

$$\mathbf{P} = \mathbf{p}_1 + \mathbf{p}_2 + \mathbf{p}_3 + \dots + \mathbf{p}_N .$$

$$\mathbf{P} = \sum_{\alpha=1}^N \mathbf{p}_\alpha = \sum_{\alpha=1}^N m_\alpha \mathbf{v}_\alpha$$

***This is the crucial result:***

□ **Theorem.**  $\dot{\mathbf{P}} = \mathbf{F}^{\text{ext}}$

where  $\mathbf{F}^{\text{ext}}$  is the sum of all external forces acting on the particles.

( dot over a letter means d/dt )

## Proof

$$\vec{p} = \sum_{\alpha} \vec{p}_{\alpha} = \sum_{\alpha} m_{\alpha} \vec{v}_{\alpha}$$

$$\frac{d\vec{p}}{dt} = \sum_{\alpha} m_{\alpha} \frac{d\vec{v}_{\alpha}}{dt} = \sum_{\alpha} \vec{F}_{\alpha}$$

$$= \sum_{\alpha} \vec{F}_{\alpha}^{\text{ext}} + \sum_{\alpha} \sum_{\substack{\beta \\ (\beta \neq \alpha)}} \vec{F}_{\alpha\beta}$$

All the internal forces cancel in pairs  
by Newton's third law;

$$\vec{F}_{12} + \vec{F}_{21} = 0 \quad \text{or} \quad \vec{F}_{\alpha\beta} + \vec{F}_{\beta\alpha} = 0$$

$$\therefore \frac{d\vec{p}}{dt} = \sum_{\alpha} \vec{F}_{\alpha}^{\text{ext}} = \vec{F}^{\text{ext}}$$

## The principle of conservation of momentum

For an *isolated system* of  $N$  particles, the total momentum is a constant of the motion.

Proof: Because it is an isolated system, there are no external forces.

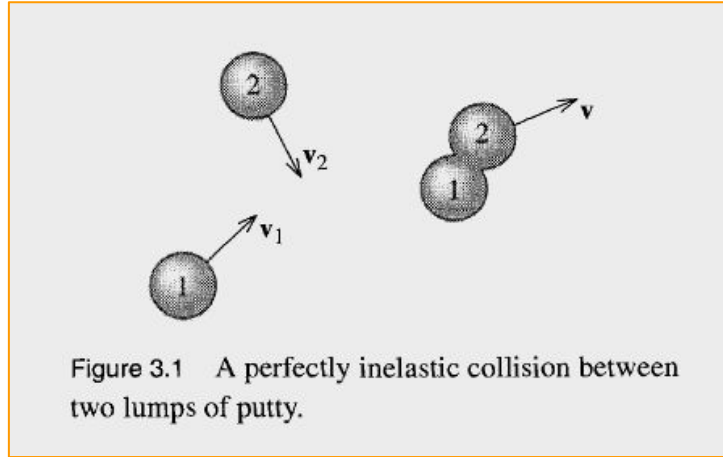
Thus the previous theorem states that

$$d\vec{p} / dt = 0.$$

Hence  $\vec{p}$  is constant in time.

## Example 3.1

*A perfectly inelastic collision*



The problem is to calculate  $\mathbf{v}$ .

**Principle:** *The total momentum is conserved.*

Before the collision,

$$\mathbf{P} = m_1 \mathbf{v}_1 + m_2 \mathbf{v}_2 \quad .$$

After the collision (stuck together)

$$\mathbf{P} = (m_1 + m_2) \mathbf{v} \quad .$$

The momentum is conserved (constant) so

$$(m_1 + m_2) \mathbf{v} = m_1 \mathbf{v}_1 + m_2 \mathbf{v}_2$$

$$\therefore \mathbf{v} = (m_1 \mathbf{v}_1 + m_2 \mathbf{v}_2) / (m_1 + m_2)$$

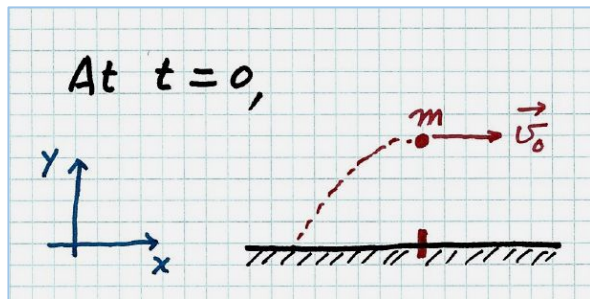
*Special cases:* If  $m_1 \gg m_2$  then  $\mathbf{v} \approx \mathbf{v}_1$  ;

if  $m_1 \ll m_2$  then  $\mathbf{v} \approx \mathbf{v}_2$  ;

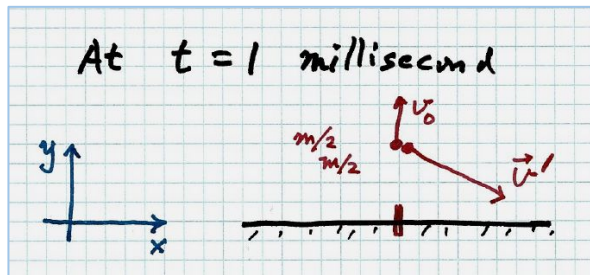
If  $m_1 = m_2$  then  $\mathbf{v} = \frac{1}{2} (\mathbf{v}_1 + \mathbf{v}_2)$  .

### Another example:

an exploding projectile (Taylor, Problem 3.2)



At  $t = 1$  ms, the shell explodes into 2 equal mass fragments, and one fragment goes straight up with speed  $v_0$



Calculate the velocity of the second fragment.

By conservation of momentum,

$$m v_0 \hat{e}_x = \frac{m}{2} v_0 \hat{e}_y + \frac{m}{2} \vec{v}'$$
$$\therefore \vec{v}' = 2v_0 \hat{e}_x - v_0 \hat{e}_y$$

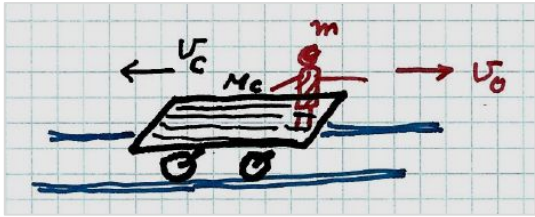
How much energy was released in the explosion?

$$\begin{aligned} T_f - T_i &= \frac{1}{2} \frac{m}{2} v_0^2 + \frac{1}{2} \frac{m}{2} v'^2 - \frac{1}{2} m v_0^2 \\ &= \frac{1}{4} m v_0^2 + \frac{1}{4} m (4v_0^2 + v_0^2) - \frac{1}{2} m v_0^2 \\ &= m v_0^2 \end{aligned}$$

← The explosion released that amount of energy, which was converted into kinetic energy of the fragments.

## Another example:

*a hobo jumps off a railroad flat car*  
(Taylor, Problem 3.4)



Assuming the car is initially at rest, calculate the increase of kinetic energy.

**Principle:** Momentum is conserved.

$$p = 0 = m v_0 + M_c v_c$$

$$v_c = - \frac{m v_0}{M_c}$$

$$\Delta K.E. = \frac{1}{2} m v_0^2 + \frac{1}{2} M_c v_c^2$$

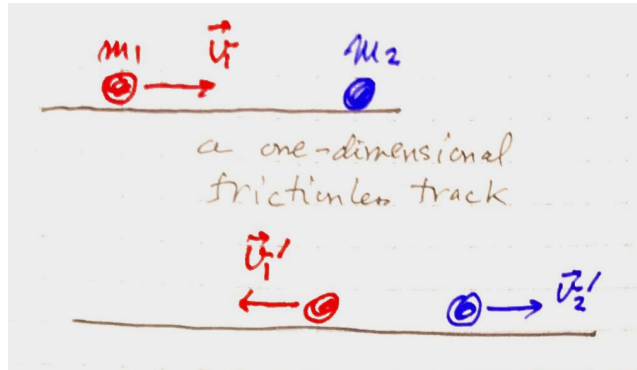
$$= \frac{1}{2} m v_0^2 + \frac{1}{2} M_c \left( -m v_0 / M_c \right)^2$$

$$= \frac{1}{2} m v_0^2 \left[ 1 + \frac{m}{M_c} \right] \leftarrow \begin{cases} \approx \frac{1}{2} m v_0^2 & \text{if } m \ll M_c \\ m v_0^2 & \text{if } m = M_c \end{cases}$$

**Comment.** In impulsive collisions, momentum is conserved because during the short time of the collision, external forces are negligible.

A head-on elastic collision in 1 dimension with one particle at rest

Before



After

$$m_1 u_1 = m_1 u_1' + m_2 u_2'$$

$$\frac{1}{2} m_1 u_1^2 = \frac{1}{2} m_1 u_1'^2 + \frac{1}{2} m_2 u_2'^2$$

(elastic)

Algebra to solve the equations  $\Rightarrow$

$$u_1' = \frac{m_1 - m_2}{m_1 + m_2} u_1 \quad \text{and} \quad u_2' = \frac{2m_1}{m_1 + m_2} u_1$$

Special cases

- $m_1 = m_2 \Rightarrow u_1' = 0, u_2' = u_1$
- $m_1 > m_2 \Rightarrow m_1$  continues forward
- $m_1 < m_2 \Rightarrow m_1$  bounces back

Conservation of momentum and Newton's third law

- *Is momentum always conserved?*

- We've already seen that *kinetic energy is not always conserved.*

Chapter 4 will introduce potential energy. But *mechanical energy (T+U) is not always conserved.*

**THERMODYNAMICS is necessary to understand that *total energy is always conserved; the first law of thermodynamics.*** ■

- *Is momentum always conserved?*

*Particle momentum* is not always conserved because there is field momentum. But *total momentum* is conserved.

- Example

Coulomb scattering:  $g + Q \rightarrow g + Q$

incoming outgoing

$m\vec{v}_i = m\vec{v}_f + M\vec{V}_f$

Bremsstrahlung:  $g + Q \rightarrow g + Q + \gamma$

or, E.M. Waves

The E.M. waves carry momentum

The image contains two diagrams. The top diagram shows Coulomb scattering: a particle 'g' (incoming) and a particle 'Q' (target) interact, resulting in 'g' and 'Q' (outgoing). The bottom diagram shows Bremsstrahlung: a particle 'g' and a particle 'Q' interact, resulting in 'g', 'Q', and a photon 'γ' (represented by wavy arrows). The text explains that the photon carries momentum.

**For PHY 321: In Newtonian mechanics, particle momentum is always conserved.**



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