Quiz Question (FRI 13)


$$
\begin{gathered}
x_{0}=-1 m \quad v=0 \\
v_{0}=2 m / s \quad v / m=3 s^{-2} \\
k / m=E=\text { constant } \\
\frac{1}{2} m v^{2}+\frac{1}{2} x^{2}=E \\
\frac{1}{2} m v_{0}^{2}+\frac{1}{2} k x_{0}^{2}=\frac{1}{2} k x_{f}^{2} \\
x_{f}=\sqrt{x_{0}^{2}+\frac{m_{1} v_{0}^{2}}{k}}=\sqrt{1+\frac{4}{3}} m=\sqrt{\frac{7}{3} m}
\end{gathered}
$$

Continuing Chapter 4 - Energy
Starting now, PHY 321 becomes more difficult.

| SECTION | TITLE |  |
| :---: | :--- | :--- |
| 4.6 | Energy in linear motion | MONDAY |
| 4.7 | Curvilinear motion, and 1D systems | MONDAY |
| 4.8 | Central forces | WEDNESDAY |
| 4.9 | Energy for a system of 2 particles | FRIDAY |
| 4.10 | Energy for many particles, and rigid bodies | FRIDAY |

Homework Assignment 8 includes four computer problems.

## Section 4.6. Energy in linear motion

 (Review)We went over this last time, and you have already read Section 4.6.

Here we are concerned with linear motion of a particle in a potential.


- the coordinate is $x$
- the equation of motion is

$$
m \ddot{x}=F(x)=-d U / d x
$$

- the energies are $T=1 / 2 \mathrm{~m} \dot{\mathrm{x}}^{2}$ and $\mathrm{U}(\mathrm{x})$

Solve the equation of motion
The first integral ( $\mathbf{x}^{\prime \prime} \Rightarrow \mathbf{x}^{\prime}$ ) comes from conservation of energy

$$
\dot{x}= \pm \sqrt{\frac{2}{m}[E-v(x)]}
$$

The second integral ( $\mathbf{x}^{\prime} \boldsymbol{\Rightarrow} \boldsymbol{x}$ ) comes from the time calculation

$$
d t^{\prime}=\frac{d x^{\prime}}{v}=\frac{d x^{\prime}}{\ddot{x}}
$$

Details - the sign and the turning points and the energy - need to be worked out ...
$\int_{t_{0}}^{t} d t^{\prime}=\int_{x_{0}}^{x} \frac{d x^{\prime}}{ \pm \sqrt{\frac{2}{2}\left[E-v\left(x^{\prime}\right)\right]}}$

$$
t-t_{0}= \pm \sqrt{\frac{m}{2}} \int_{x_{0}}^{x} \frac{d x^{\prime}}{\sqrt{E-U\left(x^{\prime}\right)}}
$$

A trivial example: the free fall problem; we did it last time.

A less trivial example: a mass on a spring; Taylor Problem 4.28. (assigned)

Here is a nontrivial example:
A particle moves on the $x$ axis, with $x>0$, and the potential energy is $\mathrm{U}(\mathrm{x})=\alpha / \mathrm{x}$. Assume $\alpha$ is positive, so the particles is repelled from the origin.

For example, consider a fixed charge at the origin and a moving charge (the "particle") on the positive $x$ axis.

Suppose the particle is released from rest at $\mathrm{x}_{0}$. Calculate $\mathrm{x}(\mathrm{t})$.

Sketch a picture.

$$
\mathrm{U}(\mathrm{x})=\alpha / \mathrm{x}
$$



The problem is to calculate $\mathrm{x}(\mathrm{t})$.

Equations

$$
\begin{gathered}
m \mathrm{x}=\alpha / \mathrm{x}^{2} \\
\mathrm{x}(0)=\mathrm{x}_{0} \text { and } \dot{\mathrm{x}}(0)=0
\end{gathered}
$$

We cant solve them directly.
We'll use the conservation of energy to get the first integral.

1. Conservation of energy

$$
\begin{aligned}
& E=\frac{1}{2} m \dot{x}^{2}+\frac{\alpha}{x}=\frac{\alpha}{x_{0}} \\
& \dot{x}^{2}=\frac{2 \alpha}{m}\left(\frac{1}{x_{0}}-\frac{1}{x}\right) \\
& \dot{x}=+\sqrt{\frac{2 \alpha}{m x_{0}}}\left(1-\frac{x_{0}}{x}\right)^{\frac{1}{2}}
\end{aligned}
$$

2. The time calculation

$$
\begin{aligned}
& d t^{\prime}=\frac{d x^{\prime}}{v}=\frac{d x^{\prime}}{\dot{x}} \\
& \int_{0}^{t} d t^{\prime}=\int_{x_{0}}^{x} \frac{d x^{\prime}}{\sqrt{\frac{2 \alpha}{m+x}}\left(1-x / x^{\prime}\right)^{\frac{1}{2}}} \\
& t=\sqrt{\frac{m x}{2 \alpha}} \int_{x_{0}}^{x}\left(\frac{x^{\prime}}{x^{\prime}-x_{0}}\right)^{1 / 2} d x^{\prime} \\
& t=\sqrt{\frac{m x_{0}^{3}}{2 \alpha}} \int_{1}^{x\left(x_{0}\right.}\left(\frac{u}{u-1}\right)^{\frac{1}{2}} d x
\end{aligned}
$$

3. Final Result

$$
\begin{aligned}
& t=\sqrt{\frac{m x_{0}^{3}}{2 \alpha}}\left\{\left[\frac{x}{x_{0}}\left(\frac{x}{x_{0}}-1\right)^{12}+\arcsin h\left[\left(\frac{x}{x}-1\right)^{\frac{1}{2}}\right]\right\}\right. \\
& t=\sqrt{\frac{m x^{3}}{2 \alpha}}\left\{\frac{1}{2} \sinh 2 \psi+\psi\right\} \leftarrow \text { baratanchive } \\
& \text { whee } \sinh \psi=\left(\frac{x}{x_{0}}-1\right)^{1 / 2} ; x=x_{0} \cosh ^{2} \psi \psi
\end{aligned}
$$

## Section 4.7. Curvilinear motion and other examples of one-dimensional motion

A system is called "one-dimensional" if the configuration is determined by a single dynamical variable.

Linear motion is strictly one-dimensional; but it's not the only kind of "one-dimensional" motion in the generalized sense.

## Curvilinear Motion

Consider a particle that is constrained to move on a curve. That is
"one-dimensional" because only a single variable is required to specify the position of the particle.

$$
\mathrm{s}=\text { arclength }
$$



Figure 4.13 An object constrained to move on a curved track can be considered to be a one-dimensional system, with the position specified by the distance $s$ (measured

Even this train moving on a curve is an example of one-dimensional motion in the generalized sense;
the configuration of the moving object is determined by a single variable, $\mathrm{s}(\mathrm{t})$.


## Curvilinear Motion



Figure 4.13 An object constrained to move on a curved track can be considered to be a one-dimensional system, with the position specified by the distance $s$ (measured

I The variable is $s=$ arclength ; $s=s(\mathrm{t})$
I The force equation is $\mathrm{ms}=\mathrm{F}_{\text {tangential }}$
I The energies are $T=1 / 2 \mathrm{~m} \mathrm{~s}^{2}$

$$
\text { and } U(s) \text { where } F_{\text {tang. }}=-d U / d s
$$

I But what keeps the particle on the curve?

Picture $m$ as a bead threaded on a stiff wire.


Normal force
$\equiv$ force of constraint;
$\equiv$ keeps the particle on the curve ;
$\equiv$ but it does no work!!! ;
$\equiv \Delta \mathrm{T}$ comes only from $\mathrm{F}_{\text {tang. }}$.

The Simple Pendulum
(Taylor Problems 4.34 and 4.38)


The mass moves on a circular arc.
I The variable is $\theta=$ angle; $\quad \theta=\theta(\mathrm{t})$
I The force equation? requires torque
I The energies are $T=1 / 2 \mathrm{~mL}^{2}(\theta)^{2}$ and $\mathrm{U}(\theta)=\mathrm{mgy}=\mathrm{mg} \mathrm{L}(1-\cos \theta)$

I The equation of motion.
We could write $\mathrm{d} \boldsymbol{l} / \mathrm{dt}=$ torque;
or, $\mathrm{dE} / \mathrm{dt}=0$ implies

$$
\frac{1}{2} m L^{2} 2 \dot{\theta} \ddot{\theta}+m g L \sin \theta \dot{\theta}=0
$$

$$
\ddot{\theta}=-\frac{g}{L} \sin \theta
$$

the simple pendulum
This is "1-dimensional" in the
generalized sense:
$\exists$ a single dynamical variable.
Assigned in the homework.

## Example 4.7 <br> STABILITY OF A CUBE BALANCED ON A CYLINDER



Figure 4.14 A cube, of side $2 b$ and center $C$, is placed on a fixed horizontal cylinder of radius $r$ and


The cylinder is fixed.
The cube is free to roll from side to side, not slipping on the cylinder. \{center directly above center for $\theta=0$ \}

Calculate $\mathrm{U}(\theta)$.

This is a " 1 d " problem; the variable is $\theta$. Analyze the potential energy. (See the Figure.)

Let $\mathrm{h}=$ the height of the center of mass of the cube. Then $U=m g h$. Now express $h$ in terms of the angle $\theta$.

Example 4.7
STABILITY OF A CUBE BALANCED ON A CYLINDER


Figure 4.14 A cube, of side $2 b$ and center $C$, is placed on a fixed horizontal cylinder of radius $r$ and

Geometrical analysis

$$
\begin{aligned}
& r \theta=d=C B \\
& h=(r+b) \cos \theta+d \sin \theta \\
& U(\theta)=m g h=m g[(r+b) \cos \theta+r \theta \sin (\theta)]
\end{aligned}
$$

Condition for equilibrium

$$
\frac{d U}{d \theta}=\pi g[-(r+b) \sin \theta+r \sin \theta+r \theta \cos \theta]
$$

$$
U^{\prime}(0)=0 \text { so } \theta=0 \text { is an excitiorium anffacotion }
$$

Condition for stability

$$
\begin{aligned}
& \frac{d^{2} U}{d \theta^{2}}=m g[-b \cos \theta+r \cos \theta-r \theta \sin \theta] \\
& U^{\prime \prime \prime}(0)=m g(r-b)
\end{aligned}
$$

The condition for stability in $z^{\prime \prime}(\%) \geqslant 0 ; b \leq r$

## STABILITY ANALYSIS

Plot $\mathrm{U}(\theta)$ for different values of $\mathrm{b} / \mathrm{r}$. (2r=diameter; 2b = width)

The cube balanced at $\theta=0$ is stable if $\mathrm{b} \leqq \mathrm{r}$;
i.e., $\theta=0$ is a stable equilibrium
if the width of the cube is smaller than the diameter of the cylinder.


Another example:

## The Atwood machine

Figure 4.15


## Atwood machine

The configuration depends on a single
variable ( $x$ ) because the length of the string is constant ( $L$ );
$\mathrm{L}=\mathrm{x}_{2}+\pi \mathrm{R}+x$; or $\quad \mathrm{x}_{2}=\mathrm{L}-\pi \mathrm{R}-x$
Analysis by energies:

$$
\begin{aligned}
& T=T_{1}+T_{2}=1 / 2\left(m_{1}+m_{2}\right) x^{\prime 2} \\
& \left.U_{1}=-m_{1} g x \quad \text { (xis downward }\right) \\
& U_{2}=-m_{2} g x_{2}=+m_{2} g x+\text { constant } \\
& \quad E=1 / 2\left(m_{1}+m_{2}\right) x^{\prime 2}+\left(m_{2}-m_{1}\right) g x
\end{aligned}
$$

Energy is constant, so
$d E / d t=0=\left(m_{1}+m_{2}\right) x^{\prime} x^{\prime \prime}+\left(m_{2}-m_{1}\right) g x^{\prime}$

$$
\left(m_{1}+m_{2}\right) x^{\prime \prime}=\left(m_{1}-m_{2}\right) g
$$

Result : constant acceleration,

$$
a=\left(m_{1}-m_{2}\right) /\left(m_{1}+m_{2}\right) g
$$

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Homework Assignment \#8
due in class Wednesday, October 25
[37] Problem 4.26 *
[38] Problem \(4.28^{* *}\) and Problem \(4.29^{* *}\) [Computer]
[39] Problem \(4.33^{* *}\) [Computer]
[40] Problem 4.34 **
[41] Problem \(4.37^{* * *}\) [Computer]
[42] Problem \(4.38^{* * *}\) [Computer]
Use the cover page.
- This is a long assignment, so start working on it now.
- Work together with 1 or 2 other students on the computer calculations.
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