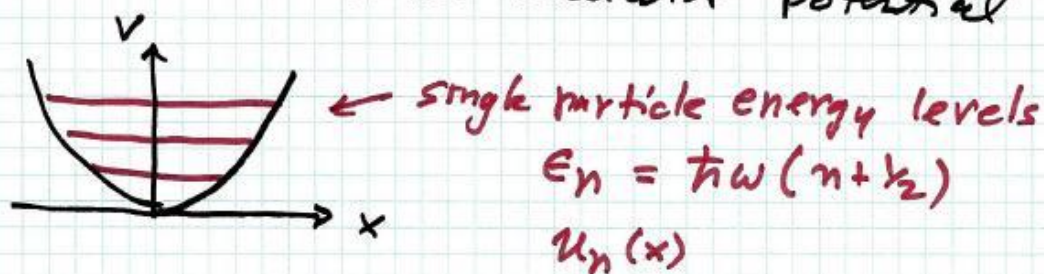


Homework Set 3

Problem 11 3 spin 0 particles in a harmonic oscillator potential



The three particle energy $= E = \hbar\omega(n_1 + \frac{1}{2} + n_2 + \frac{1}{2} + n_3 + \frac{1}{2})$
 $= \hbar\omega(n_1 + n_2 + n_3 + \frac{3}{2}) = \frac{9}{2} \hbar\omega.$

Therefore, $n_1 + n_2 + n_3 = 3.$

The 3 particle wave function must be

$$\Phi(x_1, x_2, x_3) = \alpha u_1(x_1) u_1(x_2) u_1(x_3)$$

n_1	n_2	n_3
1	1	1

$$+ \beta \frac{1}{\sqrt{6}} \left\{ u_2(x_1) u_1(x_2) u_0(x_3) + u_1(x_1) u_2(x_2) u_0(x_3) \right. \\
+ u_1(x_1) u_0(x_2) u_2(x_3) + u_2(x_1) u_0(x_2) u_1(x_3) \\
\left. + u_0(x_1) u_2(x_2) u_1(x_3) + u_0(x_1) u_1(x_2) u_2(x_3) \right\}$$

$$2 \ 1 \ 0$$

$$3 \ 0 \ 0$$

$$+ \gamma \frac{1}{\sqrt{3}} \left\{ u_3(x_1) u_0(x_2) u_0(x_3) \right. \\
+ u_0(x_1) u_0(x_2) u_3(x_3) \\
\left. + u_0(x_1) u_3(x_2) u_0(x_3) \right\}$$

where normalization requires $\alpha^2 + \beta^2 + \gamma^2 = 1.$

2 POINTS

Problem 12

$$\Phi(\vec{x}_1, \vec{x}_2) = \frac{1}{\sqrt{2}} [u_{p_1}(x_1) u_{p_2}(x_2) + u_{p_2}(x_1) u_{p_1}(x_2)]$$

where $u_p(\vec{x}) = \frac{1}{\sqrt{V}} e^{i\vec{p} \cdot \vec{x}}$ (finite volume normalization) assuming $\vec{p}_2 \neq \vec{p}_1$

(a) First, calculate $\langle \vec{p}'_1 \vec{p}'_2 | V_2 | \vec{p}_1 \vec{p}_2 \rangle$ ← b → ← a →

$$= \frac{1}{2V^2} \int d^3x_1 d^3x_2 \left\{ e^{-i\vec{p}'_1 \cdot \vec{x}_1} e^{-i\vec{p}'_2 \cdot \vec{x}_2} V(\vec{x}_1, \vec{x}_2) \left\{ e^{i\vec{p}_1 \cdot \vec{x}_1} e^{i\vec{p}_2 \cdot \vec{x}_2} + e^{i\vec{p}_2 \cdot \vec{x}_1} e^{i\vec{p}_1 \cdot \vec{x}_2} \right\} \right.$$

$$= \frac{1}{2V^2} \int d^3x_1 d^3x_2 \left\{ e^{i(\vec{p}_1 - \vec{p}'_1) \cdot \vec{x}_1} e^{i(\vec{p}_2 - \vec{p}'_2) \cdot \vec{x}_2} + e^{i(\vec{p}_2 - \vec{p}'_2) \cdot \vec{x}_1} e^{i(\vec{p}_1 - \vec{p}'_1) \cdot \vec{x}_2} + e^{i(\vec{p}_2 - \vec{p}'_1) \cdot \vec{x}_1} e^{i(\vec{p}_1 - \vec{p}'_2) \cdot \vec{x}_2} + e^{i(\vec{p}_1 - \vec{p}'_2) \cdot \vec{x}_1} e^{i(\vec{p}_2 - \vec{p}'_1) \cdot \vec{x}_2} \right\} V(\vec{x}_1, \vec{x}_2)$$

Introduce center of mass coordinate $\vec{R} = \frac{1}{2}(\vec{x}_1 + \vec{x}_2)$ and relative coordinate $\vec{r} = \vec{x}_1 - \vec{x}_2$;
or, $\vec{x}_1 = \vec{R} + \frac{1}{2}\vec{r}$ and $\vec{x}_2 = \vec{R} - \frac{1}{2}\vec{r}$.

$$\langle b | V_2 | a \rangle = \frac{1}{2V^2} \int d^3R d^3r V(\vec{x}_1, \vec{x}_2) \left\{ e^{i(\vec{p}_1 + \vec{p}_2 - \vec{p}'_1 - \vec{p}'_2) \cdot \vec{R}} \right\} \\ \left\{ e^{i(\vec{p}_1 - \vec{p}'_1 - \vec{p}_2 + \vec{p}'_2) \cdot \vec{r}/2} + e^{i(\vec{p}_2 - \vec{p}'_2 - \vec{p}_1 + \vec{p}'_1) \cdot \vec{r}/2} + e^{i(\vec{p}_2 - \vec{p}'_1 - \vec{p}_1 + \vec{p}'_2) \cdot \vec{r}/2} + e^{i(\vec{p}_1 - \vec{p}'_2 - \vec{p}_2 + \vec{p}'_1) \cdot \vec{r}/2} \right\}$$

Now let $b = a$; i.e., $\vec{p}'_1 = \vec{p}_1$ and $\vec{p}'_2 = \vec{p}_2$. Also,

Let $\vec{p}_1 - \vec{p}_2 = \vec{K}$. Then

$$\langle a | V_2 | a \rangle = \frac{1}{2V^2} \int d^3R d^3r V_2(\vec{x}_1, \vec{x}_2) \left\{ 2 + e^{-i\vec{K} \cdot \vec{r}} + e^{i\vec{K} \cdot \vec{r}} \right\}$$

(b) Now suppose $V_2(\vec{x}_1, \vec{x}_2) = V(\vec{x}_1 - \vec{x}_2) = V(\vec{r})$.

Then

$$\langle a | U | a \rangle = \frac{1}{2V^2} \int d^3R d^3r V(\vec{r}) [2 + e^{i\vec{k} \cdot \vec{r}} + e^{-i\vec{k} \cdot \vec{r}}]$$

$$= \frac{1}{2V} \{ 2 \hat{U}(\vec{0}) + \hat{U}(\vec{k}) + \hat{U}(-\vec{k}) \} \quad \vec{k} = \vec{p}_1 - \vec{p}_2$$

where $\hat{U}(\vec{k}) = \int d^3r e^{i\vec{k} \cdot \vec{r}} V(\vec{r}) = \text{Fourier transform}$

CHECK If $V(\vec{r}) = \text{constant}$ then $\langle a | U | a \rangle = C$ for $\vec{k} \neq 0$.
(C)

2 POINTS

Problem 13 To prove: $[\Omega, N] = 0$ for fermions.

Recall $N = \int \psi^\dagger(\vec{x}) \psi(\vec{x}) d^3x$

and $\Omega = \int \psi^\dagger(\vec{x}_1) \psi^\dagger(\vec{x}_2) V_2(\vec{x}_1, \vec{x}_2) \psi(\vec{x}_2) \psi(\vec{x}_1) d^3x_1 d^3x_2$,

where $\psi(\vec{x}) =$ the fermion field ; I suppress the spin index.

So $\{\psi(\vec{x}), \psi(\vec{y})\} = 0$ and $\{\psi(\vec{x}), \psi^\dagger(\vec{y})\} = \delta^3(\vec{x} - \vec{y})$.

Use these identities

$$[AB, C] = A[B, C] + [A, C]B$$

$$[AB, C] = A\{B, C\} - \{A, C\}B$$

$$[N, \Omega] = \underbrace{\int d^3x d^3x_1 d^3x_2 V_2(\vec{x}_1, \vec{x}_2)}_{\text{denote this by } F} [\psi_x^\dagger \psi_x, \psi_1^\dagger \psi_2^\dagger \psi_2 \psi_1]$$

$$= F \left(\psi_x^\dagger [\psi_x, \psi_1^\dagger \psi_2^\dagger \psi_2 \psi_1] + [\psi_x^\dagger, \psi_1^\dagger \psi_2^\dagger \psi_2 \psi_1] \psi_x \right)$$

$$= F \left(\psi_x^\dagger \left([\psi_x, \psi_1^\dagger \psi_2^\dagger] \psi_2 \psi_1 + \psi_1^\dagger \psi_2^\dagger [\psi_x, \psi_2 \psi_1] \right) \right. \\ \left. + \left([\psi_x^\dagger, \psi_1^\dagger \psi_2^\dagger] \psi_2 \psi_1 + \psi_1^\dagger \psi_2^\dagger [\psi_x^\dagger, \psi_2 \psi_1] \right) \psi_x \right)$$

Note : $[\psi_x, \psi_2 \psi_1] = 0$ and $[\psi_x^\dagger, \psi_1^\dagger \psi_2^\dagger] = 0$.

Also : $[\psi_x, \psi_1^\dagger \psi_2^\dagger] = \{\psi_x, \psi_1^\dagger\} \psi_2^\dagger - \psi_1^\dagger \{\psi_x, \psi_2^\dagger\}$
 $= \delta^3(\vec{x} - \vec{x}_1) \psi_2^\dagger - \delta^3(\vec{x} - \vec{x}_2) \psi_1^\dagger$

and $[\psi_x^\dagger, \psi_2 \psi_1] = \delta^3(\vec{x} - \vec{x}_2) \psi_1 - \delta^3(\vec{x} - \vec{x}_1) \psi_2$

$$[N, \Omega] = F \left(\delta^3(\vec{x} - \vec{x}_1) \left(\psi_x^\dagger \psi_2^\dagger \psi_2 \psi_1 - \psi_1^\dagger \psi_2^\dagger \psi_2 \psi_x \right) \leftarrow \vec{x} = \vec{x}_1 \right. \\ \left. + \delta^3(\vec{x} - \vec{x}_2) \left(-\psi_x^\dagger \psi_1^\dagger \psi_2 \psi_1 + \psi_1^\dagger \psi_2^\dagger \psi_1 \psi_x \right) \leftarrow \vec{x} = \vec{x}_2 \right)$$

$$= 0$$

QED

2 POINTS

Problem 14 Consider $\psi(\vec{x}, t)$ fermionic; suppress the spin index

In the Heisenberg picture,

$$i\hbar \frac{\partial \psi}{\partial t} = - [H, \psi]$$

where $H = H_1 + H_2$

$$\leftarrow \equiv F[\psi]$$

$$\begin{cases} H_1 = \int \psi^\dagger_\xi t_\xi \psi_\xi d^3\xi \\ H_2 = \frac{1}{2} \int \psi_1^\dagger \psi_2^\dagger V_2(\vec{x}_1, \vec{x}_2) \psi_2 \psi_1 d^3x_1 d^3x_2 \end{cases}$$

First term

$$\begin{aligned} [H_1, \psi_x] &= \int d^3\xi [\psi^\dagger_\xi t_\xi \psi_\xi, \psi_x] \\ &= \int d^3\xi \left(\underbrace{\psi^\dagger_\xi \{t_\xi \psi_\xi, \psi_x\}}_0 - \underbrace{\{\psi^\dagger_\xi, \psi_x\} t_\xi \psi_\xi}_{\delta^3(\xi-x)} \right) \\ &= -t_x \psi_x \end{aligned}$$

2nd term

$$\begin{aligned} [H_2, \psi(\vec{x})] &= \frac{1}{2} \int d^3x_1 d^3x_2 V_2(\vec{x}_1, \vec{x}_2) [\psi_1^\dagger \psi_2^\dagger \psi_2 \psi_1, \psi_x] \\ &= \frac{1}{2} \int d^3x_1 d^3x_2 V_2 \left(\underbrace{\psi_1^\dagger \psi_2^\dagger [\psi_2 \psi_1, \psi_x]}_0 + [\psi_1^\dagger \psi_2^\dagger, \psi_x] \psi_2 \psi_1 \right) \\ &= \frac{1}{2} \int d^3x_1 d^3x_2 V_2(\vec{x}_1, \vec{x}_2) \left(\underbrace{\psi_1^\dagger \{\psi_2^\dagger, \psi_x\}}_{\delta^3(x-x_2)} - \underbrace{\{\psi_1^\dagger, \psi_x\} \psi_2^\dagger}_{\delta^3(x-x_1)} \right) \psi_2 \psi_1 \\ &= \frac{1}{2} \int d^3y \left(V_2(\vec{y}, \vec{x}) \psi_y^\dagger \psi_x \psi_y - V_2(\vec{x}, \vec{y}) \psi_y^\dagger \psi_y \psi_x \right) \\ &= -\frac{1}{2} \int d^3y V(\vec{x}, \vec{y}) \psi_y^\dagger \left(\underbrace{\psi_y \psi_x - \psi_x \psi_y}_{= 2 \psi_y \psi_x} \right) \end{aligned}$$

Result

$$F[\psi] = t_x \psi(\vec{x}) + \int d^3y V(\vec{x}, \vec{y}) \psi^\dagger(y) (\psi(y) \psi(\vec{x}))$$

Note: $F[\psi]$ is NONLINEAR.

$$F[\psi] = t_x \psi(\vec{x}) + \int d^3x V(\vec{x}, \vec{y}) \psi^\dagger(y) \psi(y) \psi(\vec{x})$$

2 POINTS

Problem 15

For fermions, calculate

$$\langle 0 | c_\alpha c_\beta \psi^\dagger(\vec{r}_1) \psi^\dagger(\vec{r}_2) | 0 \rangle \equiv Q_{\alpha\beta}(\vec{r}_1, \vec{r}_2)$$

$$\psi(\vec{r}) = \sum_{\gamma} u_{\gamma}(\vec{r}) c_{\gamma}$$

$$\psi^\dagger(\vec{r}) = \sum_{\gamma} u_{\gamma}^*(\vec{r}) c_{\gamma}^\dagger$$

So

$$Q = \sum_{\gamma} u_{\gamma}^*(\vec{r}_1) \sum_{\delta} u_{\delta}^*(\vec{r}_2) \langle 0 | c_\alpha c_\beta \overset{\text{wavy}}{c_{\gamma}^\dagger} c_{\delta}^\dagger | 0 \rangle$$

$$= \sum_{\gamma\delta} u_{\gamma}^*(\vec{r}_1) u_{\delta}^*(\vec{r}_2) \left\{ \langle 0 | c_\alpha [-c_{\gamma}^\dagger c_\beta + \delta(\beta, \gamma)] c_{\delta}^\dagger | 0 \rangle \right\}$$

$$= - \sum_{\gamma\delta} u_{\gamma}^*(\vec{r}_1) u_{\delta}^*(\vec{r}_2) \langle 0 | c_\alpha c_{\gamma}^\dagger c_\beta c_{\delta}^\dagger | 0 \rangle$$

$$+ \sum_{\delta} u_{\beta}^*(\vec{r}_1) u_{\delta}^*(\vec{r}_2) \langle 0 | c_\alpha c_{\delta}^\dagger | 0 \rangle$$

$$\text{Note: } \langle 0 | c_{\gamma}^\dagger = 0 \text{ and } c_{\beta} | 0 \rangle = 0$$

$$\text{and } c_\alpha | 0 \rangle = 0 \implies$$

$$Q = - \sum_{\gamma\delta} u_{\gamma}^*(\vec{r}_1) u_{\delta}^*(\vec{r}_2) \delta_{\alpha\gamma} \delta_{\beta\delta} + \sum_{\delta} u_{\beta}^*(\vec{r}_1) u_{\delta}^*(\vec{r}_2) \delta_{\alpha\delta}$$

$$Q = u_{\beta}^*(\vec{r}_1) u_{\alpha}^*(\vec{r}_2) - u_{\alpha}^*(\vec{r}_1) u_{\beta}^*(\vec{r}_2)$$

2 POINTS

Note $Q_{\alpha\beta}(\vec{r}_1, \vec{r}_2)$ is antisymmetric w.r.t.
exchange of (α, β) , and exchange of $(\vec{r}_1, \vec{r}_2)$.