

Chapter 1

The Theory of the Photon

Sections 1.1 – 1.2

- quantum numbers $\{ \mathbf{k}, \sigma \}$
- operators $a_{\mathbf{k}\sigma}$ and $a_{\mathbf{k}\sigma}^\dagger$
- $H_{\text{rad}} = \sum \hbar \omega a_{\mathbf{k}\sigma}^\dagger a_{\mathbf{k}\sigma}$
- $\mathbf{P}_{\text{rad}} = \sum \hbar \mathbf{k} a_{\mathbf{k}\sigma}^\dagger a_{\mathbf{k}\sigma}$
- Field operators $\mathbf{A}(\mathbf{x}, t)$, $\mathbf{E}(\mathbf{x}, t)$, $\mathbf{B}(\mathbf{x}, t)$

Problem 1-1 (pages 24-25)

- The coherent state (Glauber)

Sections 1.3 – 1.4

- Interactions of radiation with matter

Interactions of Radiation and Matter

$$H = H_{\text{atom}} + H_{\text{rad}} + H_{\text{int}}$$

We know

$$H_{\text{atom}} = \frac{p^2}{2m} - \frac{Ze^2}{4\pi r} \quad (\text{one electron atom or ion})$$

$$H_{\text{rad}} = \sum_{\vec{k}\sigma} \hbar\omega \, a_{\vec{k}\sigma}^\dagger a_{\vec{k}\sigma} \quad \text{where } \omega = ck.$$

Now, what is H_{int} ?

Start with the Hamiltonian for an electron in an electromagnetic field,

$$H = \frac{(\vec{p} + \frac{e}{c} \vec{A})^2}{2m} - e\phi$$

(electron charge = $-e$)

Why?

Check Hamilton's equations,

Now write $H = p^2 / 2m - e \phi(\mathbf{x}) + H_{\text{int}}$

$$\dot{\vec{x}} = \frac{\partial H}{\partial \vec{p}} = \frac{1}{m} (\vec{p} + \frac{e}{c} \vec{A})$$

$$\begin{aligned} \dot{\vec{p}} &= - \frac{\partial H}{\partial \vec{x}} = - \frac{1}{m} (\vec{p} + \frac{e}{c} \vec{A})_j \frac{e}{c} \frac{\partial A_j}{\partial \vec{x}} \\ &\quad + e \frac{\partial \phi}{\partial \vec{x}} \\ &= - \dot{x}_j \frac{e}{c} \frac{\partial A_j}{\partial \vec{x}} + e \nabla \phi \quad (1) \end{aligned}$$

$$\begin{aligned} \text{Also,} \\ \dot{\vec{p}} &= m \ddot{\vec{x}} - \frac{e}{c} \dot{\vec{A}} \\ &= \vec{F} - \frac{e}{c} \left[\frac{\partial \vec{A}}{\partial t} + (\vec{v} \cdot \nabla) \vec{A} \right] \quad (2) \end{aligned}$$

$$\odot \quad \vec{F} = -e \left[\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right] \quad \text{as required}$$

$$\begin{aligned} \vec{v} \times \vec{B} &= \vec{v} \times (\nabla \times \vec{A}) \\ &= \nabla(\vec{v} \cdot \vec{A}) - (\vec{v} \cdot \nabla) \vec{A} \end{aligned}$$

So,

$$H_{\text{int}} = \frac{e}{mc} \vec{A}(\vec{x}) \cdot \vec{p} + \frac{e^2}{2mc^2} A^2(\vec{x})$$

Recall,

$$\vec{A}(\vec{x}) = \sum_{\vec{k}\sigma} N_k \vec{\epsilon}_{k\sigma} \left[a_{\vec{k}\sigma} e^{i\vec{k} \cdot \vec{x}} + a_{\vec{k}\sigma}^\dagger e^{-i\vec{k} \cdot \vec{x}} \right]$$

(in the Schroedinger picture !)

The first term of H_{int} will annihilate or create a photon.
The second term could annihilate or create 2 photons.

$$N_k = \text{SQRT} [\hbar c^2 / (2\omega\Omega)]$$

Radiative decays in atomic physics

$$a \rightarrow b + \gamma \quad \text{where} \quad \varepsilon_a > \varepsilon_b$$

We'll calculate the transition rate by Fermi's Golden Rule, from perturbation theory, treating H_{int} as a perturbation.

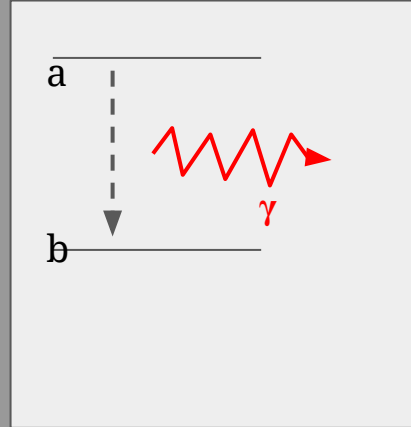
$$\text{Rate} = \frac{2\pi}{\hbar} \sum_f \left| \langle f | H_{\text{int}} | i \rangle \right|^2 \delta(\varepsilon_f - \varepsilon_i)$$

(first order approximation in time-dependent perturbation theory).

The initial state is $|i\rangle = |a\rangle_{\text{atom}} |\text{vacuum}\rangle_{\text{rad.}}$

The final state is $|f\rangle = |b\rangle_{\text{atom}} |\mathbf{k}, \varepsilon\rangle_{\text{rad.}}$

i.e., $a^\dagger_{\mathbf{k}\varepsilon} |\text{vacuum}\rangle_{\text{rad.}}$



Calculate the transition matrix element

$$\langle f | H_{\text{int}} | i \rangle =$$

$$\langle b | \langle \vec{k} \vec{\epsilon} | \frac{e}{mc} \vec{A} \cdot \vec{p} | \text{vacuum} \rangle | a \rangle$$

$$= \frac{e}{mc} \langle b | \underbrace{\langle \vec{k} \vec{\epsilon} | \vec{A} | \text{vac.} \rangle}_{= N_{\vec{k}} \hat{\epsilon}_{\vec{k}\sigma} e^{-i\vec{k}\cdot\vec{x}}}} \cdot \vec{p} | a \rangle$$

$$= N_{\vec{k}} \hat{\epsilon}_{\vec{k}\sigma} e^{-i\vec{k}\cdot\vec{x}} \\ (= \text{coefficient of } a_{\vec{k}\sigma}^+ \text{ in } \vec{A}(\vec{x}))$$

$$= \frac{e}{mc} N \hat{\epsilon} \cdot \langle b | e^{-i\vec{k}\cdot\vec{x}} \vec{p} | a \rangle$$

The size of the atom is \ll the wavelength of the photon ,

$$10^{-10} \text{ m} \ll 100 \text{ nm} = 10^{-7} \text{ m} ;$$

so we can approximate $e^{-i\vec{k}\cdot\vec{x}} \approx 1$. (*electric dipole approx.*)

$$\text{Rate} = \frac{2\pi}{\hbar} \sum_{\vec{k}\sigma} \left(\frac{e}{mc}\right)^2 N^2 \left| \hat{\epsilon} \cdot \langle b | \vec{p} | a \rangle \right|^2 \delta(\epsilon_b + \hbar\omega - \epsilon_a)$$

The rest is just math.

Example:

The decay $H(2p) \rightarrow H(1s) + \gamma$ for atomic hydrogen.

The atomic matrix element is

$$\langle 2p, \mathbf{m} | \mathbf{p} | 1s \rangle \quad \text{for } \mathbf{m} \in \{0, +1, -1\}$$

This is just a problem in ordinary quantum mechanics.

The calculation is not difficult.

Recall the wave functions,

$$\Phi_{1s}(\mathbf{x}) = \exp(-r/a_0) / (\pi a_0^3)^{1/2}$$

& let's take $m=0$,

(the total decay rate would be the same for $m = -1$ or $+1$)

$$\Phi_{2p}(\mathbf{x}) = (r/a_0) \exp(-r/2a_0) \cos\theta / (32 \pi a_0^3)^{1/2}$$

$$\begin{aligned} \langle 2p, 0 | \mathbf{p} | 1s \rangle &= -i\hbar \int r^2 dr \sin\theta d\theta d\varphi \Phi_{2p}^* \nabla \Phi_{1s} \\ &= \mathbf{e}_z (i\hbar/a_0) (2/3)^4 \sqrt{2} \end{aligned}$$

Sum over polarizations

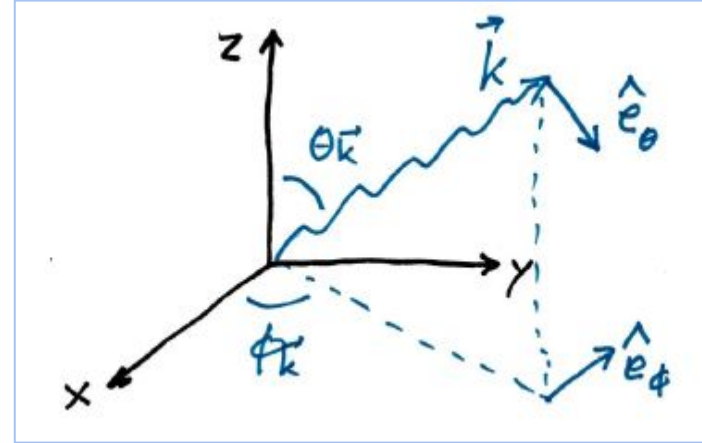
To get the *unpolarized decay rate* we must sum over the final two polarizations, ϵ_1 and ϵ_2 . These are orthogonal and perpendicular to \mathbf{k} . We can let $\epsilon_1 = \mathbf{e}_\phi$ and $\epsilon_2 = \mathbf{e}_\theta$. Note that $\mathbf{e}_z \cdot \mathbf{e}_\phi = 0$ and $\mathbf{e}_z \cdot \mathbf{e}_\theta = -\sin \theta_k$.

So,

$$\begin{aligned}\sum_{\sigma} |\epsilon_{\sigma} \cdot \langle \mathbf{b} | \mathbf{p} | \mathbf{a} \rangle|^2 &= (\mathbf{e}_\theta \cdot \mathbf{e}_z)^2 (\hbar/a_0)^2 (2^9 / 3^8) \\ &= \sin^2 \theta_k (\hbar/a_0)^2 (2^9 / 3^8)\end{aligned}$$

Unpol. decay rate =

$$\frac{2\pi}{\hbar} \sum_{\mathbf{k}} \left(\frac{e}{mc}\right)^2 N^2 \sin^2 \theta_k \left(\frac{\hbar}{a_0}\right)^2 \frac{2^9}{3^8} \delta(\epsilon_b + \hbar\omega - \epsilon_a)$$



The differential decay rate (for $m = 0$)

We need to calculate the sum over \mathbf{k} .

Remember the infinite volume limit,

$$\sum_{\mathbf{k}} = \frac{\Omega}{(2\pi)^3} \int d^3k$$

(periodic boundary conditions)

$$\text{and } d^3k = k^2 dk d\Omega_k$$

To get the *total rate*, integrate over all directions of \mathbf{k} ;
to get the *differential rate*, don't integrate over
directions...

$$\frac{dR}{d\Omega_k} = \left\{ \frac{2\pi}{\hbar} \frac{\Omega}{(2\pi)^3} \left(\frac{e}{mc}\right)^2 \cancel{N^2} \left(\frac{1}{a_0}\right)^2 \frac{2^9}{3^8} \right\} [\sin^2 \theta_k] \frac{\hbar c^2}{2\omega\Omega} \\ \times \int k^2 dk \delta(\epsilon_b + \hbar ck - \epsilon_a)$$

important cancellation $\Omega/\Omega = 1$; ang. distribution $d\sin^2 \theta_k$
because $\Delta m = 0$

$$\text{and } \int \frac{k^2 dk}{a^3} \delta(\epsilon_b + \hbar ck - \epsilon_a) = \frac{1}{\hbar c^2} \left(\frac{\epsilon_b - \epsilon_a}{\hbar c} \right)$$

$$\frac{dR}{d\Omega_k} = \left\{ \frac{2\pi}{\hbar} \frac{\Omega}{(2\pi)^3} \left(\frac{e}{mc}\right)^2 \cancel{N^2} \left(\frac{\hbar}{a_0}\right)^2 \frac{2^9}{3^8} \right\} [\sin^2 \theta_k] \frac{\hbar c^2}{2\omega\Omega} \\ \times \int k^2 dk \delta(\epsilon_b + \hbar c k - \epsilon_a)$$

The k integration is evaluated with the delta function.

!!! Remember how to do delta function integrations:

$$\int_{-\infty}^{\infty} \delta(ax) f(x) dx = 1/a f(0)$$

(i.e., needs the Jacobian factor)

$$\int \frac{k^2 dk}{a} \delta(\epsilon_b + \hbar c k - \epsilon_a) = \frac{1}{\hbar c^2} \left(\frac{\epsilon_b - \epsilon_a}{\hbar c} \right)$$

A general formula is given in Mandl & Shaw:

Equation 1.50

$$\frac{dR}{d\Omega_k} = \frac{e^2}{8\pi^2} \frac{\omega^3}{\hbar c^3} |\langle b | \vec{x} | a \rangle|^2 \sin^2 \theta_k$$

in rationalized cgs units;

$$\alpha = \frac{e^2}{4\pi\hbar c} = \frac{1}{137}.$$

E.g., estimate

Order of magnitude for $H_{2p} \rightarrow H_{1s} + \gamma$,

$$\frac{\alpha}{c^2} \cdot \frac{1}{\hbar^3} \left(\frac{me^4}{2\hbar^2} \right)^3 \left(\frac{\hbar^2}{me^2} \right)^2 \propto m \alpha^5$$

Rydberg
en.

Bohr
rad.

Homework problems due Friday, January 27

Hint: For these calculations start with Equation (1.50).

Problem 7. Calculate accurately the mean lifetime of the 2p state of atomic hydrogen.

Problem 8. Estimate the mean lifetime for a nuclear gamma decay process. (You only need to calculate the order of magnitude.)