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Chapter 9 : PLANE E. M. WAVES... ... AND PROPAGATION IN MATTER

★ Waves in matter ; reflection and refraction ; part 1

★ Sections 9.2 and 9.3

Reflection and Refraction from a Dielectric Interface

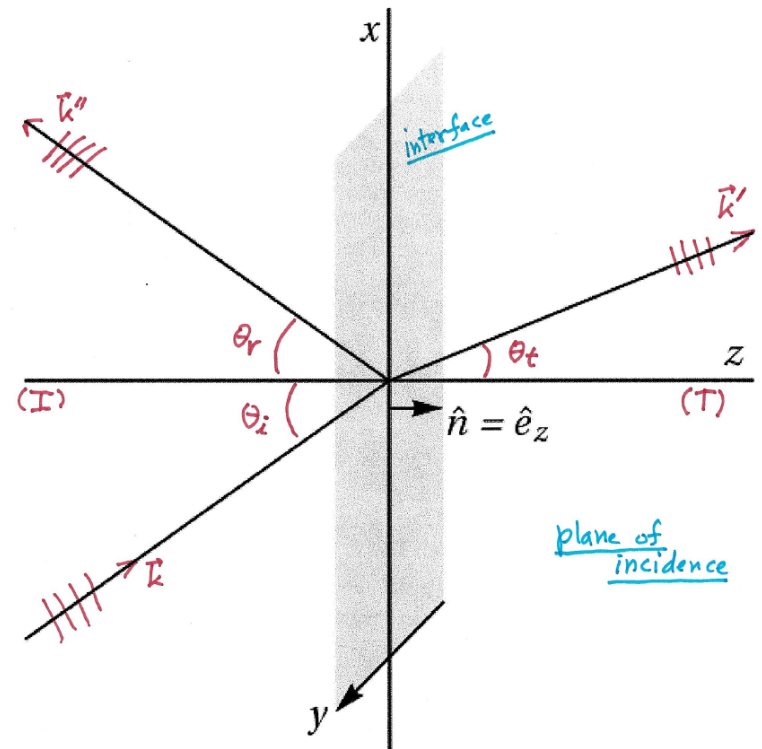
In either dielectric we have plane waves with frequency ω . Recall the properties of plane waves (Sec. 9.1.) *Now all we need to do is satisfy the four boundary conditions.*

WT use a different and more difficult method with Green's functions. (See page 444.)

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W1F1

W1F1



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Be careful when referring to Wilcox and Thron,
because they use a bad notation.
In Sections 9.2 ~ 9.4 there are two regions.

Region I = the region of the incident and
reflected waves; WT calls this "region 2".

Region T = the region of the transmitted wave;
WT calls this "region 1".

In the lecture, I'll use ϵ_I and ϵ_T .

To compare to WT,

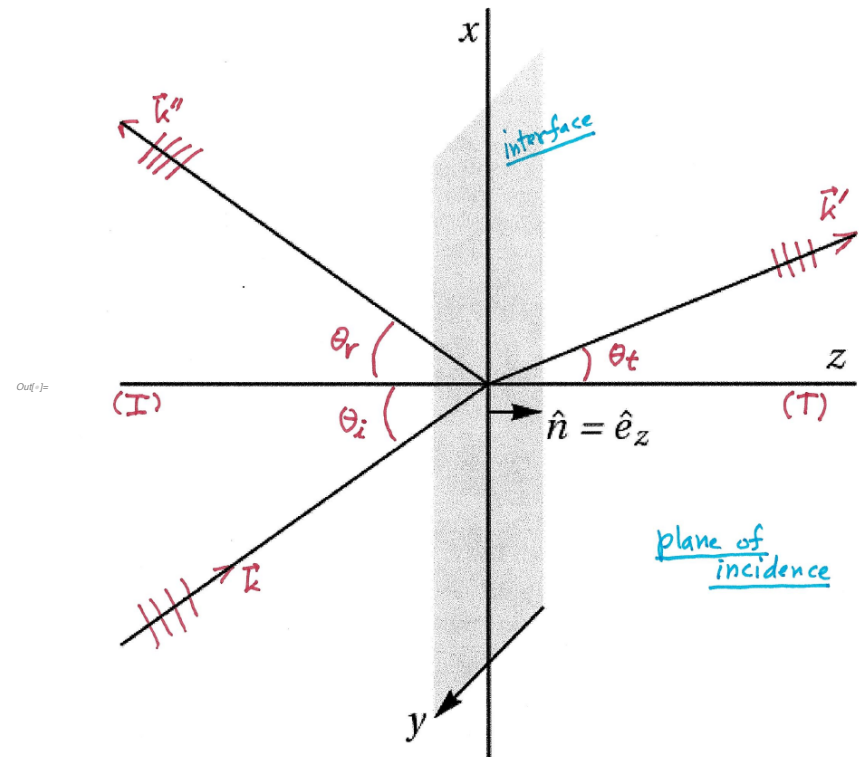
set $\epsilon_I = \epsilon_2$ and $\epsilon_T = \epsilon_1$.

To compare to Jackson,

set $\epsilon_I = \epsilon_1$ and $\epsilon_T = \epsilon_2$.

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in[1]= W1F1



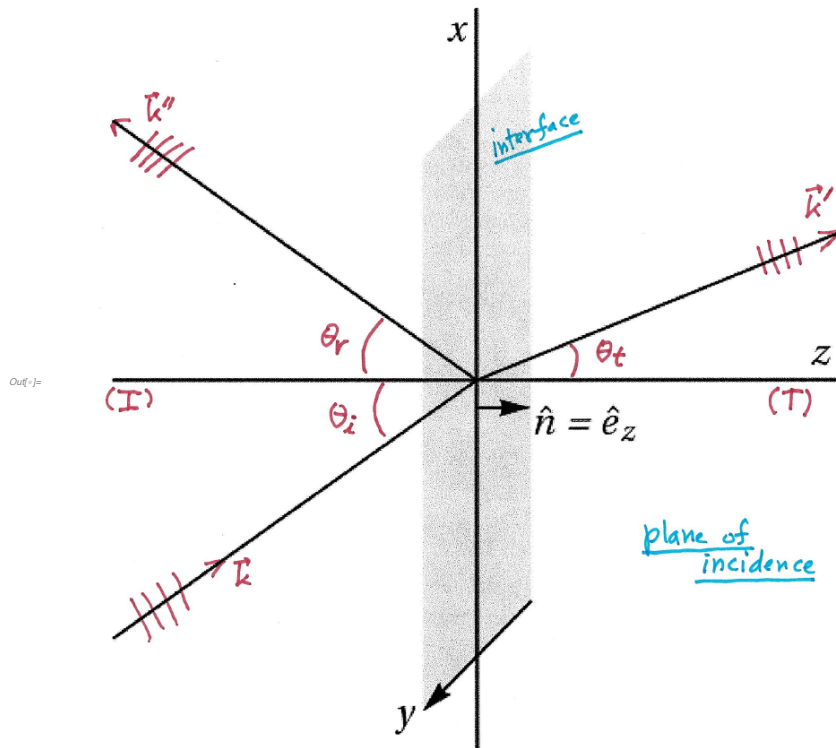
Interface = xy plane■

Plane of incidence = xz plane■

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Section 9.2: TE polarization

in[]:= W1F1



The TE waves are polarized with the electric fields transverse to the plane of incidence.

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in[]:= scTE1

$$\vec{E}(\vec{x}, t) = E_0 \hat{e}_y e^{i(\vec{k} \cdot \vec{x} - \omega t)}$$

$$\vec{k} = (k_x, 0, k_z)$$

$$\vec{E}'(\vec{x}, t) = E_0' \hat{e}_y e^{i(\vec{k}' \cdot \vec{x} - \omega t)}$$

$$\vec{k}' = (k_x, 0, k_z')$$

$$\vec{E}''(\vec{x}, t) = E_0'' \hat{e}_y e^{i(\vec{k}'' \cdot \vec{x} - \omega t)}$$

$$\vec{k}'' = (k_x, 0, k_z'')$$

Out[]:=

Magnetic fields.

By Faraday's Law, $\vec{k} \times \vec{E} = (\omega/c) \vec{B}$;

also, $\omega/k = v_{\text{phase}} = c/n$;

$$\begin{aligned} \therefore \vec{B} &= n \hat{k} \times \vec{E} \\ &= n \{ -\cos\theta_i E_y, 0, \sin\theta_i E_y \} \end{aligned}$$

4. Boundary conditions at the interface ($z=0$)

- $E_{\text{tang.}} : E_0 + E_0'' = E_0'$
- $D_{\text{normal}} : 0 + 0 = 0$
- B_{normal} (z comp.) :

$$n_I \sin \theta_i (E_0 + E_0'') = n_T \sin \theta_t E_0'$$

i.e., $E_0 + E_0'' = E_0'$
- $H_{\text{tang.}}$ (x comp.) :

$$(1/\mu_I) n_I \cos \theta_i (E_0 - E_0'')$$

$$= (1/\mu_T) n_T \cos \theta_t E_0'$$

scTE2; scTE3; scTE4; scTE5

Two equations for two unknowns
(E_0' and E_0'')

$$E_0 + E_0'' = E_0'$$

$$\left(\frac{n}{\mu}\right)_I (E_0 - E_0'') \cos \theta_i = \left(\frac{n}{\mu}\right)_T E_0' \cos \theta_t$$

Out- \rightarrow

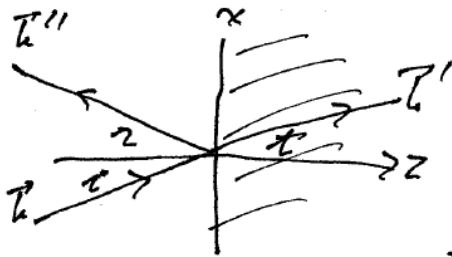
a bit of algebra \rightarrow

$$\frac{E_0'}{E_0} = \frac{2 \left(\frac{n}{\mu}\right)_I \cos \theta_i}{\left(\frac{n}{\mu}\right)_I \cos \theta_i + \left(\frac{n}{\mu}\right)_T \cos \theta_t}$$

$$\frac{E_0''}{E_0} = \frac{\left(\frac{n}{\mu}\right)_I \cos \theta_i - \left(\frac{n}{\mu}\right)_T \cos \theta_t}{\left(\frac{n}{\mu}\right)_I \cos \theta_i + \left(\frac{n}{\mu}\right)_T \cos \theta_t}$$

Section 9.3: TM polarization

[1] sCTM1



TM-polarization. Start with

$$\vec{H} = H_0 \hat{e}_y e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\vec{k} = k_x \hat{e}_x + k_z \hat{e}_z$$

$$\vec{H}' = H_0' \hat{e}_y e^{i(k_x' x - \omega t)}$$

$$\vec{k}' = k_x' \hat{e}_x + k_z' \hat{e}_z$$

$$\vec{H}'' = H_0'' \hat{e}_y e^{i(k_x'' x - \omega t)}$$

$$\vec{k}'' = k_x'' \hat{e}_x + k_z'' \hat{e}_z$$

The boundary conditions for H_{tang} and $B_{\text{normal}} \rightarrow$

$$H_0 + H_0'' = H_0' \quad (1)$$

$$0 + 0 = 0 \quad (\text{automatic})$$

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[1] sCTM2

Bound \vec{E}

$$\nabla \times \vec{H} = \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \text{ so } \vec{k} \times \vec{H} = -\frac{\omega}{c} \vec{D}$$

$$\begin{aligned} \vec{D} &= -\frac{1}{\omega} c \vec{k} \times \vec{H} \\ &= -\frac{c}{\omega} (k_x \hat{e}_x + k_z \hat{e}_z) \times H_y \hat{e}_y \\ &= -\frac{c}{\omega} H_y (-k_z \hat{e}_x + k_x \hat{e}_z) \\ &= \frac{c}{\omega} H_y (k_z \hat{e}_x - k_x \hat{e}_z) \end{aligned}$$

Normal is continuous at $z=0$, so

$$\frac{c}{\omega} k_x (H_0 + H_0'') = \frac{c}{\omega} k_x' H_0'$$

$$H_0 + H_0'' = H_0' \quad (\text{same as before}).$$

Tang. is continuous, so

$$\frac{c}{\omega} k_z (H_0 - H_0'') = \frac{c}{\omega} k_z' H_0'$$

$$\text{Note } \frac{c k_z}{\omega} = \frac{c k}{\omega} \cos \theta_2 = \eta_T \cos \theta_2$$

$$\text{and } \frac{c k_x}{\omega} = \eta_T \sin \theta_2$$

$$\therefore \left(\frac{\eta}{c}\right) \cos \theta_2 (H_0 - H_0'') = \left(\frac{\eta}{c}\right) \sin \theta_2 H_0'$$

Solve (1) and (2)

Note: $\frac{n}{\epsilon} = \frac{\mu}{\eta}$

$$\frac{H_0''}{H_0} = \frac{\left(\frac{\mu}{n}\right)_I \cos \theta_i - \left(\frac{\mu}{n}\right)_T \cos \theta_t}{\left(\frac{\mu}{n}\right)_I \cos \theta_i + \left(\frac{\mu}{n}\right)_T \cos \theta_t}$$

or

$$\frac{H_0'}{H_0} = \frac{2\left(\frac{\mu}{n}\right)_I \cos \theta_i}{\left(\frac{\mu}{n}\right)_I \cos \theta_i + \left(\frac{\mu}{n}\right)_T \cos \theta_t}$$

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$$\frac{H_0''}{H_0} = \frac{\left(\frac{\mu}{n}\right)_T \cos \theta_i - \left(\frac{\mu}{n}\right)_I \cos \theta_t}{\left(\frac{\mu}{n}\right)_T \cos \theta_i + \left(\frac{\mu}{n}\right)_I \cos \theta_t}$$

$$\frac{H_0'}{H_0} = \frac{2\left(\frac{\mu}{n}\right)_T \cos \theta_i}{\left(\frac{\mu}{n}\right)_T \cos \theta_i + \left(\frac{\mu}{n}\right)_I \cos \theta_t}$$

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FRESNEL'S EQUATIONS

For simplicity, assume $\mu_I = \mu_T = 1$.

(This is a very reasonable assumption, physically. Why?)

TE waves have (9.79)

$$E^{\text{trans}} / E^{\text{inc}} = 2 n_I \cos \theta_i / (n_I \cos \theta_i + n_T \cos \theta_t)$$

$$E^{\text{refl}} / E^{\text{inc}} = (n_I \cos \theta_i - n_T \cos \theta_t) / (n_I \cos \theta_i + n_T \cos \theta_t)$$

TM waves have (9.102)

$$B^{\text{trans}} / B^{\text{inc}} = 2 n_T \cos \theta_i / (n_T \cos \theta_i + n_I \cos \theta_t)$$

$$B^{\text{refl}} / B^{\text{inc}} = (n_T \cos \theta_i - n_I \cos \theta_t) / (n_T \cos \theta_i + n_I \cos \theta_t)$$

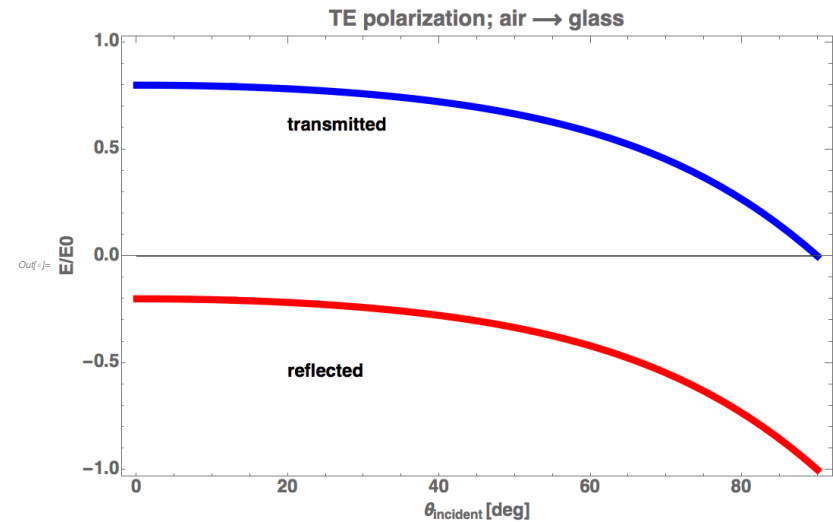
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Now plot graphs of some cases

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bs = {FontFamily → "Helvetica",
      FontWeight → Bold, FontSize → 24};
(* TE case; I = vacuum; T = glass *)
nI = 1; nT = 1.5;
β[α_] = ArcSin[(nI / nT) * Sin[α]]; (* Snell's law *)
trans[α_] =
  (2 * nI * Cos[α]) / (nI * Cos[α] + nT * Cos[β[α]]);
refl[α_] = (nI * Cos[α] - nT * Cos[β[α]]) /
  (nI * Cos[α] + nT * Cos[β[α]]);
epi = {Line[{{0, 0}, {90, 0}}],
       Text["transmitted", {20, 0.65}, {-1, 1}],
       Text["reflected", {20, -0.5}, {-1, 1}]};
plotTE =
  Plot[{trans[θi * Pi / 180], refl[θi * Pi / 180]},
       {θi, 0, 90},
       PlotRange → {{-2, 92}, {-1.03, 1.03}},
       PlotStyle → {{Blue, Thickness[0.01]},
                    {Red, Thickness[0.01]}},
       Frame → True, FrameLabel →
         {"θincident [deg]", "E/E0"},
       BaseStyle → bs, ImageSize → 2^8 * 4,
       PlotLabel → "TE polarization; air → glass",
       Epilog → epi] // Rasterize

```



When light is incident from air into glass, the reflected electric field wave undergoes a 180° phase change. (What about the magnetic field wave?) When light is incident from glass into air, the reflected electric field wave has no phase change. (Homework problem 3-7.)
(Recall Newton's Rings; the central spot is dark because of destructive interference.)

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nI:= (* TM case; I = vacuum; T = glass *)
nI = 1; nT = 1.5;
β[α_] = ArcSin[(nI / nT) * Sin[α]]; (* Snell's law *)
trans[α_] =
  (2 * nT * Cos[α]) / (nT * Cos[α] + nI * Cos[β[α]]);
refl[α_] = (nT * Cos[α] - nI * Cos[β[α]]) /
  (nT * Cos[α] + nI * Cos[β[α]]);
epi = {Line[{{0, 0}, {90, 0}]},
  Text["transmitted", {20, 1.0}, {-1, 1}],
  Text["reflected", {20, 0.}, {-1, 1}];
plotTM =
  Plot[{trans[θi * Pi / 180], refl[θi * Pi / 180]},
    {θi, 0, 90},
    PlotRange → {{-2, 92}, {-1.03, 1.53}},
    PlotStyle → {{Blue, Thickness[0.01]},
      {Red, Thickness[0.01]}},
    Frame → True, FrameLabel →
      {"θincident [deg]", "B/B0"},
    BaseStyle → bs, ImageSize → 2^8 * 4,
    PlotLabel → "TM polarization ; air → glass",
    Epilog → epi] // Rasterize

```

