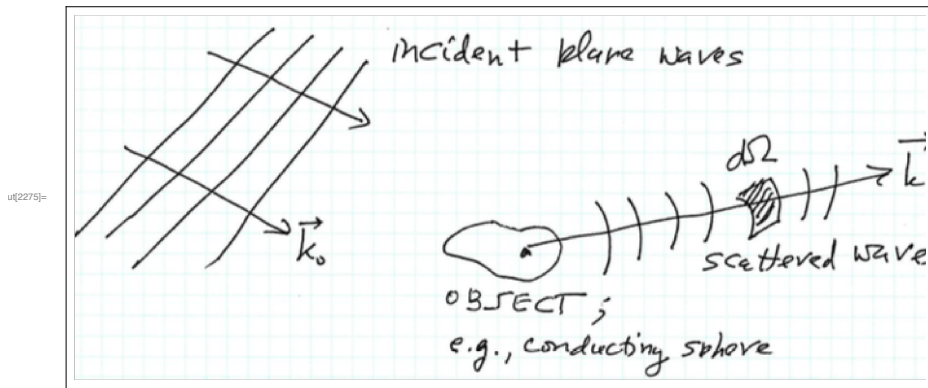


## Chapter 12 - Scattering and Diffraction

### Scattering

Plane waves incident on an object...

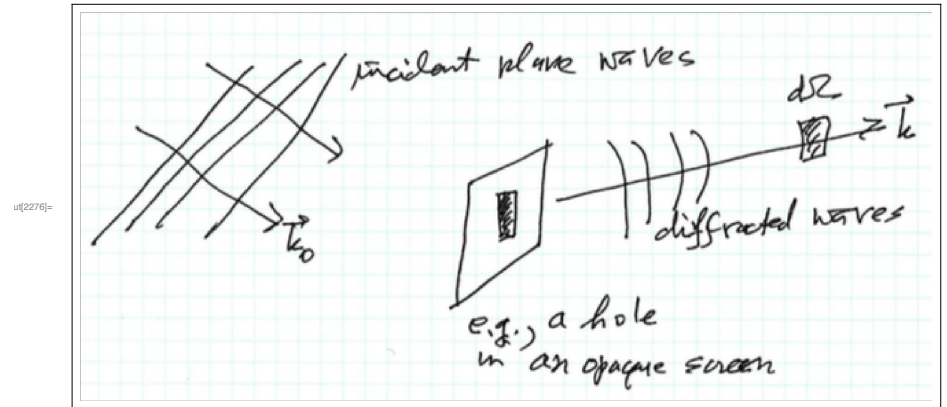
scan1



### Diffraction

Plane waves incident on a hole in a screen...

scan2



For both problems, the goal is to find the scattered fields in the far zone.

Wilcox and Thron: first scattering, then diffraction;

Jackson: first diffraction, then scattering

## Review Section 12.1

In[2278]:= scan3

$$\vec{E} = \vec{E}_{\text{inc}} + \vec{E}_{\text{sc}}$$

$$\vec{E}_{\text{inc}}(\vec{x}, t) = \hat{\epsilon}_i \cdot E_0 e^{i(\vec{k}_0 \cdot \vec{x} - \omega t)}$$

$$\vec{E}_{\text{sc}}(\vec{x}, t) = \vec{E}_{\text{sc}}(\vec{x}) e^{-i\omega t}$$

$$\vec{E}_{\text{sc}}(\vec{x}) = \frac{e^{ikr}}{r} E_0 \vec{f}(\vec{k}, \vec{k}_0) \quad (\text{far zone})$$

Out[2278]:=

Scattering amplitude

$$\vec{f}(\vec{k}, \vec{k}_0) = \frac{2k}{c\epsilon_0} \int d^3x' \vec{J}(\vec{x}') e^{-i\vec{k}' \cdot \vec{x}'}$$

$$\Rightarrow \left( \frac{d\sigma}{d\Omega} \right)_{f,i} = \left| \hat{\epsilon}_f \cdot \vec{f}_{(i)}(\vec{k}, \vec{k}_0) \right|^2$$

## Section 12.2

### The Kirchhoff identity for scattering

The equation for the scattering amplitude  $\vec{f}(\vec{k}, \vec{k}_0)$  in Section 12.1 depends on the current density  $\vec{J}(\vec{x}')$ .

Gustav Kirchhoff had a different idea — express  $\vec{f}(\vec{k}, \vec{k}_0)$  in terms of the fields on the boundary surfaces.

## Preliminaries

### Green's first identity

In[2279]= scan4

$$\int_V d^3x (\psi \nabla^2 \phi + \nabla \psi \cdot \nabla \phi) = \oint_S da \psi \hat{n} \cdot \nabla \phi$$

Proof  $\oint_S da \hat{n} \cdot (\psi \nabla \phi)$

$$= \int_V d^3x \nabla \cdot (\psi \nabla \phi) \quad (\text{Gauss theorem})$$

$$= \int_V d^3x \{ \psi \nabla^2 \phi + \nabla \psi \cdot \nabla \phi \} \quad \checkmark$$

Out[2279]=

### Green's second identity

In[2281]= scan5

$$\int_V d^3x [\psi \nabla^2 \phi - \phi \nabla^2 \psi] = \oint_S da (\psi \hat{n} \cdot \nabla \phi - \phi \hat{n} \cdot \nabla \psi)$$

Proof using (1)

$$\int_V d^3x [\psi \nabla^2 \phi - \phi \nabla^2 \psi + \nabla \psi \cdot \nabla \phi - \nabla \phi \cdot \nabla \psi]$$

$$= \oint_S da [\psi \hat{n} \cdot \nabla \phi - \phi \hat{n} \cdot \nabla \psi] \quad \checkmark$$

Out[2281]=

## Notations

in[2282]= WT12p1 (\* FIGURE 12.1 \*)

First we use Green's second identity to obtain a new expression for the electric field. The Helmholtz equations for  $\vec{E}_{sc}$  and  $G^{out}$  are:

$$-(\nabla'^2 + k^2)\vec{E}_{sc}(\vec{x}') = \frac{4\pi}{c}ik \left( \vec{J}(\vec{x}') + \frac{1}{k^2}\nabla'(\nabla' \cdot \vec{J}(\vec{x}')) \right), \quad (12.38)$$

$$(\nabla'^2 + k^2)G^{out}(\vec{x}, \vec{x}', k) = -4\pi\delta(\vec{x} - \vec{x}'). \quad (12.39)$$

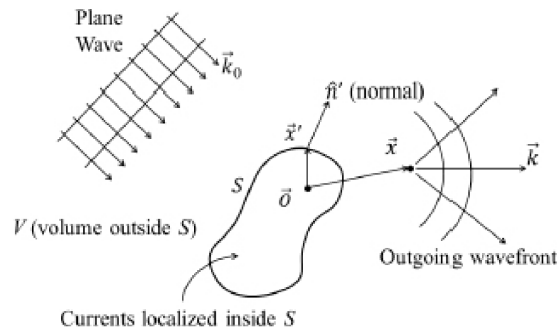


Fig. 12.1 Notation for derivation of Kirchhoff's identity.

After a long derivation, the cross section

is

$$\begin{aligned} \left( \frac{d\sigma}{d\Omega} \right)_{f,i} &= \left| \hat{\epsilon}_f \cdot \vec{f}_{(i)}(\vec{k}, \vec{k}_0) \right|^2 \\ &= |Q|^2 \end{aligned}$$

and

$$\begin{aligned} Q &= \frac{ik}{4\pi E_0} \oint_S da' e^{-i\vec{k} \cdot \vec{x}'} \\ &= \left\{ \hat{\epsilon}_f \cdot \left[ \hat{n}' \times \vec{B}_{sc} + \vec{k} \times (\hat{n}' \times \vec{E}_{sc}) \right] \right\} \end{aligned}$$

This determines the cross section, in terms of the scattered fields on the surface  $S$ . Recall,  $S$  is any surface that encloses the current.

## Section 12.3 : Short wavelength scattering from a conducting sphere

In[2251]= WT12p2

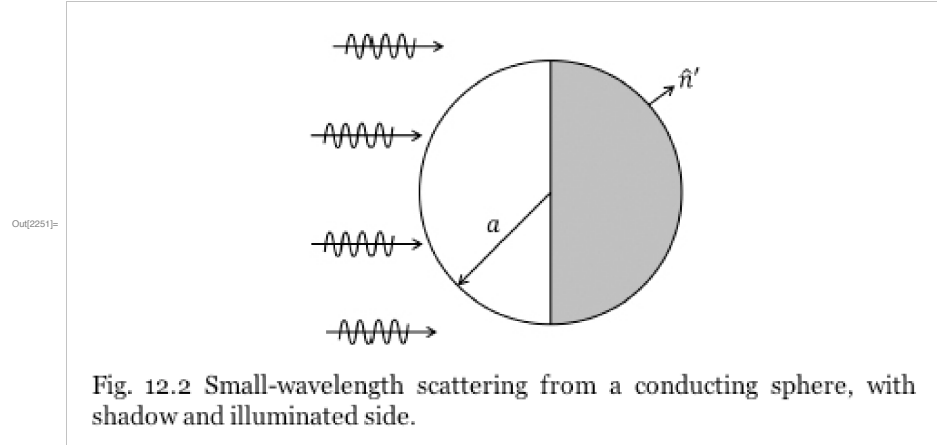


Fig. 12.2 Small-wavelength scattering from a conducting sphere, with shadow and illuminated side.