

PHY 842 SUMMARY

The purpose of a course like PHY 842 is to learn...

- first principles
- math methods
- examples

The course had many subjects:

CH9 Electromagnetic waves & optics

+ Frequency dependence of polarization;
 $\epsilon(\omega)$

CH11 Radiation by systems; $\vec{J}(\mathbf{x},t) \rightarrow dP/d\Omega$

+ Radiation by charges; $\vec{r}(t) \rightarrow \vec{E}$ and $\vec{B} \rightarrow dP/d\Omega$

→ general results

→ Larmor's formula

→ synchrotron radiation

→ Cherenkov radiation

→ Bremsstrahlung

CH12 Scattering of light

+ Diffraction of light

FIRST PRINCIPLES FOR EVERYTHING

Microscopic Maxwell's equations

$$\nabla \cdot \vec{E} = 4\pi \rho \quad \text{and} \quad \nabla \times \vec{E} = - \frac{\partial \vec{B}}{c \partial t}$$

$$\nabla \cdot \vec{B} = 0 \quad \text{and} \quad \nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{\partial \vec{E}}{c \partial t}$$

gaussian units

First Principles for continuous media

Macroscopic Maxwell's equations

$$\nabla \cdot \vec{D} = 4\pi \rho_{\text{free}} \quad \text{and} \quad \nabla \times \vec{E} = - \frac{\partial \vec{B}}{c \partial t}$$

$$\nabla \cdot \vec{B} = 0 \quad \text{and} \quad \nabla \times \vec{H} = \frac{4\pi}{c} \vec{J}_{\text{free}} + \frac{\partial \vec{D}}{c \partial t}$$

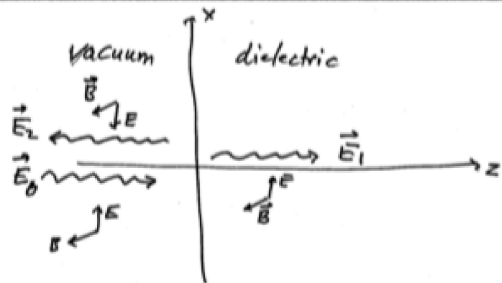
BUT what are \vec{D} and \vec{H} ?

*They come from \vec{P} and \vec{M} (how?)
& that leads to ϵ and μ (how?)*

You'll need to know this for the final exam.

Optics

Light is normally incident, from air, on the planar surface of a dielectric medium. What fraction of the energy reflects from the surface?



$$\vec{E}_0 = E_0 \hat{e}_x e^{i(kz - \omega t)}$$

$$\vec{H}_0 = \vec{B}_0 = E_0 \hat{e}_y e^{i(kz - \omega t)}$$

$$\vec{E}_1 = E_1 \hat{e}_x e^{i(k'z - \omega t)}$$

$$\vec{H}_1 = \vec{B}_1 = B_1 \hat{e}_y e^{i(k'z - \omega t)}$$

$$\vec{E}_2 = -E_2 \hat{e}_x e^{i(kz - \omega t)}$$

$$\vec{H}_2 = \vec{B}_2 = E_2 \hat{e}_y e^{i(kz - \omega t)}$$

Boundary Conditions:

$$E_0 - E_2 = E_1$$

$$E_0 + E_2 = B_1$$

Questions:

(1) What is B_1 ? (2) What is k' ?

The classical electron theory $\Rightarrow \epsilon(\omega)$

How did we do that?

recall...

$$m \ddot{\vec{x}} = -\gamma \dot{\vec{x}} - m \omega_0^2 \vec{x} - e \vec{E}_0 e^{-i\omega t}$$

Re implied

$$\vec{P}(t) = -e \dot{\vec{x}}(t) n_b$$

$$\vec{P} = \chi_e(\omega) \vec{E}_0 e^{-i\omega t} \text{ and } \epsilon = 1 + 4\pi \chi_e$$

\Rightarrow

$$\epsilon(\omega) = 1 + \frac{4\pi e^2}{m} \frac{n_b^2}{\omega_0^2 - \omega^2 - i\omega\gamma}$$

complex and frequency dependent, which has interesting consequences.

Plasma frequency?

For a plasma, $\omega_0 = 0$.

$$v_x(t) = \dot{x}(t) = \frac{-eE_0}{m(\gamma - i\omega)} e^{-i\omega t}$$

$$J_x(t) = n_c (-e v_x(t)) = \sigma(\omega) E_x(t)$$

$$\therefore \sigma(\omega) = \frac{e^2}{m} \frac{n_c}{\gamma - i\omega}$$

\Rightarrow general dispersion relation is

$$c^2 k^2 = \omega^2 \mu \{ \epsilon(\omega) + 4\pi i \sigma(\omega)/\omega \}$$

exercise 9.5.5

For high frequencies,

$$\epsilon_{\text{eff}}(\omega) = 1 - \frac{\omega_p^2}{\omega^2} \text{ where } \omega_p = ???$$

The general theory of radiation

The general theory applies to radiation, and also to scattering of light by microscopic particles.

Exercise 12.1.1. Use the model of [Section 9.5](#) and the harmonic results in (11.112) to discuss the scattering from free or bound electrons in a medium in electric dipole approximation. Using a linearly polarized incoming wave and integrating over angles, show that one obtains the cross section

$$\sigma(\omega) \equiv \int d\Omega \left(\frac{dP}{d\Omega} \right)_{\text{avg}} / |\vec{S}_0| = \frac{8\pi}{3} r_e^2 f(\omega),$$

where $r_e = e^2/(mc^2)$ is called the *classical electron radius* (encountered again in [Chapter 14](#)), and where

$$f(\omega) \equiv \frac{\omega^4}{[(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2]},$$

which gives a resonance in the cross section at $\omega \approx \omega_0$. The low and high energy forms of this cross section are

$$\omega \ll \omega_0, \frac{\omega_0^2}{\gamma} : \quad \sigma(\omega) \approx \frac{8\pi}{3} r_e^2 \left(\frac{\omega}{\omega_0} \right)^2,$$

$$\omega \gg \omega_0, \gamma : \quad \sigma(\omega) \approx \frac{8\pi}{3} r_e^2.$$

The first result gives rise to so-called *Rayleigh scattering* and the second is called the *Thomson cross section*. In particular, Rayleigh scattering implies that sunlight is preferentially scattered at higher frequencies, resulting in our blue sky and red sunsets!

Lecture of October 26

The classical electron model of light scattering

- An electromagnetic wave impinges upon a molecule

The incoming E.M. wave is a plane wave, moving in the z direction, and linearly polarized in the x direction,

$$\vec{E}(\vec{x}, t) = \hat{e}_x E_0 e^{i(kz - \omega t)}$$

● Motion of the electron(s)

Using the classical electron model, the motion of the negative charge is described by

$$m \frac{d^2 \vec{r}}{dt^2} = -K \vec{r} - \gamma \frac{d\vec{r}}{dt} - e E_0 \hat{e}_x e^{-i\omega t}$$

i.e., a damped driven oscillator.

The steady-state solution: the electron oscillates in the x direction, with the driving frequency ω ;

$$\vec{r}(t) = \hat{e}_x e^{-i\omega t} A$$

where

$$A = \frac{-e E_0}{m(\omega_0^2 - \omega^2) - i\gamma \omega}$$

$$\omega_0 = \sqrt{K/m} = \text{the "natural frequency"}.$$

● The electron radiates

Larmor's formula; the average power integrated over all directions is

$$P_{\text{avg}} = \frac{2 e^2 a^2}{3 c^3} \text{ where } a^2 = \left\langle \left(\ddot{x} \right)^2 \right\rangle.$$

be careful about taking the real part of $x(t)$

We have $A = A_1 + i A_2$, and so

$$x(t) = \text{Re} \{ e^{-i\omega t} A \} = A_1 \cos(\omega t) + A_2 \sin(\omega$$

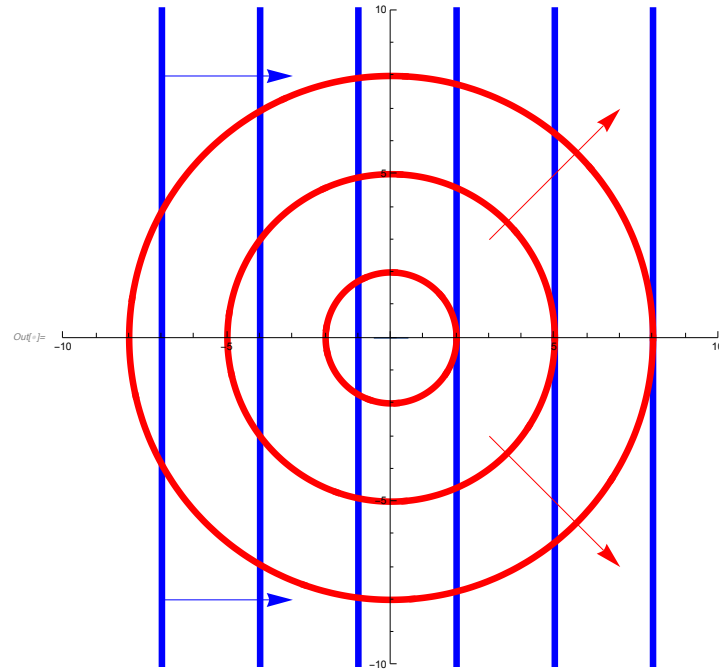
t)

$$\left\langle \ddot{x}^2 \right\rangle = \frac{1}{2} \omega^4 (A_1^2 + A_2^2)$$

$$\begin{aligned} P_{\text{avg}} &= \frac{e^2}{3 c^3} \omega^4 (A_1^2 + A_2^2) \\ &= \frac{e^4 E_0^2 \omega^4}{3 c^3 (\gamma^2 \omega^2 + m^2 (\omega^2 - \omega_0^2)^2)} \end{aligned}$$

● Scattering and the cross section

fig1513



The definition of the scattering cross section is

$$\sigma = \frac{P}{S_{\text{inc}}}$$

where P = outgoing power averaged over time;

and S_{inc} = incoming intensity

≡ incident power per unit area averaged over time

$$S_{\text{inc}} = \frac{c}{4\pi} \langle \vec{E} \times \vec{B} \rangle = \frac{c}{8\pi} E_0^2$$

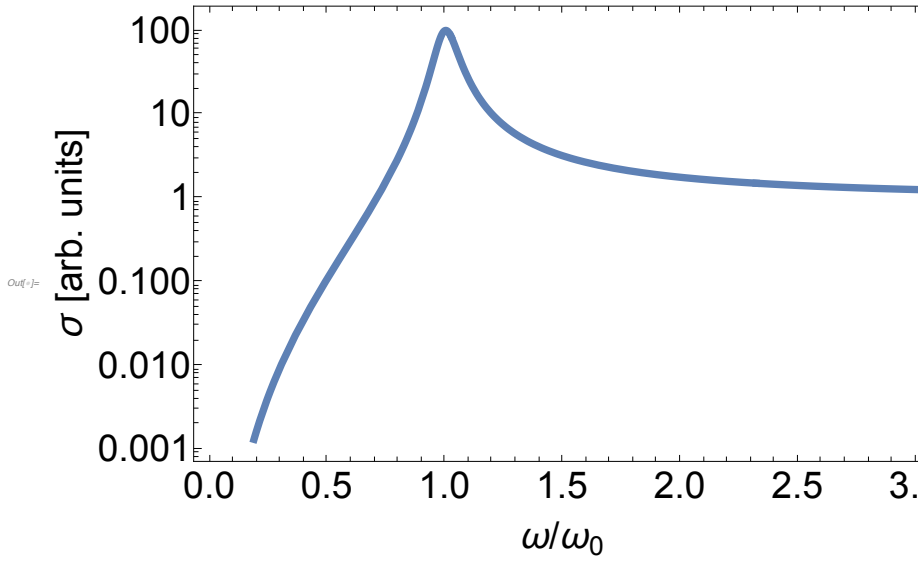
Thus,

$$\sigma = \frac{e^4 E_0^2 \omega^4}{3 c^3} \frac{1}{m^2 (\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \frac{8\pi}{c E_0^2}$$

$$\sigma = \frac{8\pi}{3} r_e^2 \frac{\omega^4}{(\omega_0^2 - \omega^2)^2 + (\gamma\omega/m)^2}$$

$$r_e = \frac{e^2}{mc^2} = \text{the "classical radius" of the electron}$$

in[]:= fig1514



← Rayleigh →

← resonant →

← Thomson →