
Problem 7-4 (Exercise 11.2.2) part (a)

Some preliminary calculations

Given this charge density,

$$\rho(\vec{x}', t') = \rho_0 a \delta(r' - a) \delta(\cos\theta') \sin(\phi' - \omega t')$$

The electric dipole moment is

$$\begin{aligned}\vec{p}(t') &= \int d^3x' \vec{x}' \rho(\vec{x}', t') \\ &= \rho_0 a^3 \int d\phi' \sin(\phi' - \omega t') \{ e_x a \cos\phi' + e_y a \sin\phi' \} \\ &= \rho_0 a^4 \pi \{ -e_x \sin(\omega t') + e_y \cos(\omega t') \}\end{aligned}$$

Thus

$$\begin{aligned}\vec{p}(t_0) &= -\rho_0 a^4 \pi \omega^2 \{ -e_x \sin(\omega t_0) + e_y \cos(\omega t_0) \} \\ |\vec{p}(t_0)| &= \rho_0 a^4 \pi \omega^2 \\ \vec{p}(t_0) \cdot \hat{n} &= -\rho_0 a^4 \pi \omega^2 \{ -\sin\theta \cos\phi \sin(\omega t_0) + \sin\theta \sin\phi \cos(\omega t_0) \} \\ &= -\rho_0 a^4 \pi \omega^2 \sin\theta \sin(\phi - \omega t_0) \\ \cos[\gamma(t_0)] &= -\sin\theta \sin(\phi - \omega t_0)\end{aligned}$$

The exact (i.e., time-dependent) power distribution

We'll calculate the power distribution from Eqs (11.55) and (11.41),

$$dP/d\Omega = r^2/(4\pi c) (\hat{n} \times \partial \vec{A} / \partial t)^2 \text{ and } \vec{A}(\vec{x}, t) \approx 1/(cr) \int d^3x' \vec{J}(\vec{x}', t_r)$$

$$\text{where } t_r = t - r/c + \hat{n} \cdot \vec{x}' / c = t_0 + \hat{n} \cdot \vec{x}' / c; \quad t_0 = t - r/c.$$

The current density is $\vec{J}(\vec{x}', t') = \rho(\vec{x}', t') \vec{v} = \rho(\vec{x}', t') \omega \hat{e}(\phi')$.

So,

$$\begin{aligned} \vec{A}(\vec{x}, t) &= 1/(cr) \int d^3x' \rho_0 a^2 \omega \delta(r'-a) \delta(\cos\theta') \sin(\phi' - \omega t_r) \hat{e}(\phi') \\ &= \Lambda \int d\phi' \sin(\phi' - \omega t_r) \{ -\hat{e}_x \sin\phi' + \hat{e}_y \cos\phi' \} \end{aligned}$$

where

$$\Lambda = (\rho_0 a^4 \omega) / (cr) \text{ and } t_r = t_0 + \hat{n} \cdot \vec{x}' / c; \quad t_0 = t - r/c.$$

$$\begin{aligned} \hat{n} \cdot \vec{x}' &= (\hat{e}_x \sin\theta \cos\phi + \hat{e}_y \sin\theta \sin\phi + \hat{e}_z \cos\theta) \cdot (\hat{e}_x a \cos\phi' + \hat{e}_y a \sin\phi') \\ &= a \sin\theta \cos(\phi - \phi') \end{aligned}$$

$$\vec{A}(\vec{x}, t) = \Lambda \int d\phi' \sin[\phi' - \omega t_0 + k a \sin\theta \cos(\phi - \phi')] \{ -\hat{e}_x \sin\phi' + \hat{e}_y \cos\phi' \}$$

Note $k = \omega/c$.

Let $k a \sin\theta = \xi$.

The time derivative —

$$\dot{\vec{A}} = -\omega \Lambda \int d\phi' \cos[\phi' - \omega t_0 + \xi \cos(\phi - \phi')] \{ -\hat{e}_x \sin\phi' + \hat{e}_y \cos\phi' \}$$

Important trick: Expand in Bessel functions

$$\text{Equation (4.11)} : \exp\{i \xi \cos(\phi - \phi')\} = \sum_m i^m e^{im(\phi - \phi')} J_m(\xi)$$

Write $\cos[\dots] = \operatorname{Re} \exp[i(\phi' - \omega t_0)] \exp[i \xi \cos(\phi - \phi')]$.

$$\dot{\vec{A}} = -\omega \Lambda \operatorname{Re} \int d\phi' \exp[i(\phi' - \omega t_0)] \exp[i \xi \cos(\phi - \phi')] \{ -\hat{e}_x \sin\phi' + \hat{e}_y \cos\phi' \}$$

$$\begin{aligned}
&= -\omega \wedge \operatorname{Re} \sum_m i^m J_m(\xi) \int d\phi' \exp[i(\phi' - \omega t_0)] \exp[i m (\phi - \phi')] \\
&\quad \{-\hat{e}_x \sin\phi' + \hat{e}_y \cos\phi'\} \\
&= -\omega \wedge \operatorname{Re} \sum_m i^m J_m(\xi) e^{im\phi} e^{-i\omega t_0} \\
&\quad \times \int d\phi' \exp[-i(m-1)\phi'] \{-\hat{e}_x \sin\phi' + \hat{e}_y \cos\phi'\}
\end{aligned}$$

Note that the integral over ϕ' is zero unless $m = 0$ or $m = 2$.

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In[1]:= m = 0;
A0 = Integrate[Exp[-I (m - 1) \phi'] Sin[\phi'], {\phi', 0, 2 Pi}],
B0 = Integrate[Exp[-I (m - 1) \phi'] Cos[\phi'], {\phi', 0, 2 Pi}]];
m = 2;
A2 = Integrate[Exp[-I (m - 1) \phi'] Sin[\phi'], {\phi', 0, 2 Pi}],
B2 = Integrate[Exp[-I (m - 1) \phi'] Cos[\phi'], {\phi', 0, 2 Pi}];

Out[1]= {1/2 \pi, \pi}
Out[2]= {-1/2 \pi, \pi}

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$$\begin{aligned}
\vec{A} &= -\omega \wedge \operatorname{Re} \{ J_0(\xi) e^{-i\omega t_0} [-e_x A0 + e_y B0] \\
&\quad - J_2(\xi) e^{i(2\phi - \omega t_0)} [-e_x A2 + e_y B2] \} \\
\vec{A} &= -\omega \wedge \operatorname{Re} \{ J_0(\xi) e^{-i\omega t_0} [-e_x i\pi + e_y \pi] \\
&\quad - J_2(\xi) e^{i(2\phi - \omega t_0)} [-e_x (-i\pi) + e_y \pi] \} \\
\vec{A} &= -\omega \wedge \pi \{ J_0(\xi) [-e_x \sin(\omega t_0) + e_y \cos(\omega t_0)] \\
&\quad - J_2(\xi) [-e_x \sin(2\phi - \omega t_0) + e_y \cos(2\phi - \omega t_0)] \}
\end{aligned}$$

Next, calculate $\hat{n} \times \vec{A}$

$$\hat{n} = e_x \sin\theta \cos\phi + e_y \sin\theta \sin\phi + e_z \cos\theta$$

$$\hat{n} \times e_x = -e_z \sin\theta \sin\phi + e_y \cos\theta$$

$$\hat{n} \times e_y = +e_z \sin\theta \cos\phi - e_x \cos\theta$$

$$\begin{aligned}
\hat{n} \times \vec{A} &= -\omega \wedge \pi \{ \\
&\quad J_0(\xi) [-(-e_z \sin\theta \sin\phi + e_y \cos\theta) \sin(\omega t_0) \\
&\quad + (e_z \sin\theta \cos\phi - e_x \cos\theta) \cos(\omega t_0)]
\end{aligned}$$

$$\begin{aligned}
& -J_2(\xi) [-(-e_z \sin \theta \sin \phi + e_y \cos \theta) \sin(2\phi - \omega t_0) \\
& \quad + (e_z \sin \theta \cos \phi - e_x \cos \theta) \cos(2\phi - \omega t_0)] \} \\
& \hat{n} \times \vec{A} = -\omega \wedge \pi \{ \\
& \quad J_0(\xi) [e_z \sin \theta \cos(\phi - \omega t_0) + \cos \theta (-e_x \cos(\omega t_0) - e_y \sin(\omega t_0))] \\
& \quad - J_2(\xi) [e_z \sin \theta \cos(\phi - \omega t_0) + \cos \theta (-e_x \cos(2\phi - \omega t_0) - e_y \sin(2\phi - \omega t_0))] \}
\end{aligned}$$

Simplifications

$$\text{term\#1} = \hat{e}_z (-\omega \wedge \pi) (J_0 - J_2) \sin \theta \cos(\phi - \omega t_0)$$

$$\text{term\#2} = \hat{e}_x \cos \theta (-\omega \wedge \pi) \{ -J_0 \cos(\omega t_0) + J_2 \cos(2\phi - \omega t_0) \}$$

$$\text{term\#3} = \hat{e}_y \cos \theta (-\omega \wedge \pi) \{ -J_0 \sin(\omega t_0) + J_2 \sin(2\phi - \omega t_0) \}$$

Now calculate $(\hat{n} \times \vec{A})^2$

$$(\hat{n} \times \vec{A})^2 = (\mathbf{t1})^2 + (\mathbf{t2})^2 + (\mathbf{t3})^2$$

The rest is just algebra.

It is not too difficult to calculate $\mathbf{t1}^2 + \mathbf{t2}^2 + \mathbf{t3}^2$.

It makes it easier to use Mathematica.

The final result is

$$(\hat{n} \times \vec{A})^2 = (\omega \wedge \pi)^2 \{ (J_0(\xi) + J_2(\xi))^2 \sin^2 \gamma - 4 J_0(\xi) J_2(\xi) \cos^2(\phi - \omega t_0) \}$$

Recall $\Lambda = (\rho_0 a^4 \omega) / (c r)$.

Thus,

$$\begin{aligned}
dP/d\Omega &= r^2/(4\pi c) (\hat{n} \times \vec{A})^2 \\
&= \pi \rho_0^2 \omega^4 a^8 / (4 c^3) \{ (J_0(\xi) + J_2(\xi))^2 \sin^2 \gamma - 4 J_0(\xi) J_2(\xi) \cos^2(\phi - \omega t_0) \}
\end{aligned}$$

Part (b)

Now calculate the time-averaged radiation rate.

We'll have $\langle \cos^2(\phi - \omega t_0) \rangle$ and $\langle \sin^2(\phi - \omega t_0) \rangle$, which are equal to 1/2.

Also, $\langle \sin^2 \gamma \rangle = 1 - \langle \cos^2 \gamma \rangle = 1 - (1/2) \sin^2 \theta = (1/2) (1 + \cos^2 \theta)$.

So the result is

$$\begin{aligned}\langle dP/d\Omega \rangle &= \pi \rho_0^2 \omega^4 a^8 / (4 c^3) \circ \{ (J_0 + J_2)^2 / 2 * (1 + \cos^2 \theta) - 2 J_0 J_2 \} \\ &= \pi \rho_0^2 \omega^4 a^8 / (8 c^3) \circ \{ (J_0 - J_2)^2 + \cos^2 \theta (J_0 + J_2)^2 \}\end{aligned}$$