

GENERAL PURPOSE DETECTOR @ VLHC

2 DETECTORS ... BOTH WILL BE G.P.
(+ SOME SPECIALIZED ?)

OBJECTS : JETS, γ , e, μ , τ , (b-jets), \cancel{E}_T

MAGNET (μ , G.P.) INNER SOLENOID, FWD + OUTER TOROIDS

GENERIC TRACKING (e, μ , τ , b, h...) Si/O ... STRIPS
PIXELS
3D PIXELS

b-JET TAGGING (H, t, SUSY...) VERTICES (mm), e/ μ IN JETS

e.m. CALORIMETRY e-IDENT+MEAS γ -IDENT+MEAS EM PART OF JETS
+ SHOWER-MAX. LAYER

had CALORIMETRY JET SPECTROSCOPY, QCD, COMPOSITENESS, BH.
 $b\bar{b} \rightarrow jj, j\bar{Z}$ DIRECTION (CORE) & E
ISOLATED μ IDENT/MEAS.

μ ON MEASUREMENT

DO NOT GIVE UP ON THIS FUNDAMENTAL FERMION!

\bar{F}_μ FOR $X \rightarrow n\mu + \gamma$, \bar{X}_μ

ALSO FOR \cancel{E}_T !! GIVE UP ON μ & GIVE UP ON \cancel{E}_T !

$\frac{\sigma}{\sigma_{SM}} \approx 20\%$ @ 5 [10?] TeV

MEASURE INSIDE COIL, VERIFY BEHIND CALO.

(SAGITTA σ (5 TeV, 4T, 2m.) $\approx 120 \mu\text{m}$)

INTERIM REPORT GP. 4C DETECTORS

DENISOV / ALBROW

D. DENISOV	PREVIOUS VLHC DETECTOR EFFORTS
M. ALBROW	VLHC CALORIMETRY
A. SKUJA	CMS CALORIMETRY
O. LOBBAN	CAL RESPONSE OPTIMIZN.; Q-SC HYBRID
S. PARKER	3D PIXEL DETECTOR
D. WOOD	μ DETECTION, $D\bar{X} \rightarrow$ VLHC
M. FORTNER	μ TRIGGERING, RECONSTRUCTION
D. CHRISTIAN	PIXEL DETECTORS
N. MOKHOV	ENERGY DEPOSITION @ VLHC
D. DENISOV	RADIATION DOSES

What we would like to measure?

- Because there is no specific physics process to be studied, general purpose detector(s) is a must for an energy frontier machine
- In order to set benchmarks different process are used: top studies, Higgs/SUSY discovery, etc. Let's consider 'recent' top quark discovery
- In order to discover top quark both CDF and D0 had to detect: leptons (e/μ), jets, tag b-quarks (leptons or displaced vertex), and neutrinos (missing E_T):

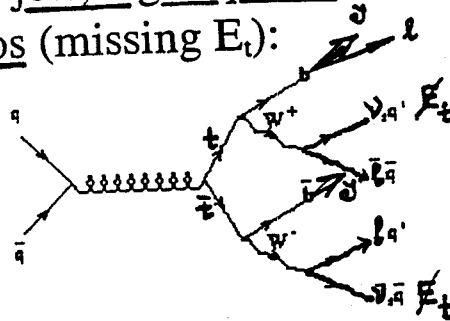


Figure 1: Tree level top quark production by $q\bar{q}$ annihilation followed by the Standard Model top quark decay chain.

Decay mode	Branching ratio
$t\bar{t} \rightarrow (q\bar{q}b)(q\bar{q}b)$	36/81
$t\bar{t} \rightarrow (q\bar{q}b)(c\bar{c}b)$	12/81
$t\bar{t} \rightarrow (q\bar{q}b)(\mu\nu b)$	12/81
$t\bar{t} \rightarrow (q\bar{q}b)(\tau\nu b)$	12/81
$t\bar{t} \rightarrow (c\bar{c}b)(\mu\nu b)$	2/81
$t\bar{t} \rightarrow (c\bar{c}b)(\tau\nu b)$	2/81
$t\bar{t} \rightarrow (\mu\nu b)(\tau\nu b)$	2/81
$t\bar{t} \rightarrow (c\bar{c}b)(c\bar{c}b)$	1/81
$t\bar{t} \rightarrow (\mu\nu b)(\mu\nu b)$	1/81
$t\bar{t} \rightarrow (\tau\nu b)(\tau\nu b)$	1/81

So, we needed tracker (e/μ , vertex), calorimeter (jets, electrons), and muon system

BACKGROUNDS & RADIATION LOADS

-NIKOLAI MOKHOV

DPMJET / MARS CALCULATIONS.

FROM COLLISIONS AT I.P.

	<u>STAGE 1</u>	<u>STAGE 2</u>
	40 TeV	175 TeV
	$1 \cdot 10^{34}$	$2 \cdot 10^{34}$
CENTRAL TRACKER		
Φ_{+} $\text{cm}^{-2} \text{s}^{-1} @ 10 \text{cm}$	$3 \cdot 10^7$	$1 \cdot 10^8$
Φ_n	$3 \cdot 10^6$	$1 \cdot 10^7$
MRAD / YEAR \approx	10	30

SMALL ANGLES (END CAP CLOSE TO PIPE)

$\times 100 \rightarrow \times 1000 !!$

FORWARD μ SYSTEM - MUST HAVE

WELL DESIGNED SHIELDING - THEN NON-ISSUE.

FROM BEAM LOSS: $500 \text{ m}^{-1} \text{s}^{-1}$ IN WARM STRAIGHT

APPROPRIATE SHIELDING PLUG IN TUNNEL-HALL,
AROUND FRONT COLLIMATOR, ENDCAP CALORIMETER etc.

THEN MACHINE B/G
COLLISIONS \sim few % , OK.

Conclusions

1. Cross sections and multiplicities changing with energy increase very slowly.
2. Multiplicity distributions have long non-Gaussian tails.
3. Radiation dose in the central region is function of luminosity, "not" \sqrt{s} .
4. Most of charged tracks are low ^{transverse} momentum
 $\langle p_T \rangle \approx 0.6 \text{ GeV}$
5. All parameters discussed are based on "extrapolation":
 we could see surprises at very high energy!

TRACKING

TRACKING USING DRIFT CELLS \approx mm
PROB. TOO SLOW FOR CENTRAL TRACKER.
OK FOR MUON SYSTEM.

ALL Silicon (\diamond !?) CENTRAL TRACKING
STRIPS ... MINISTRIPS (eg $50\mu \times 5\text{mm}$)

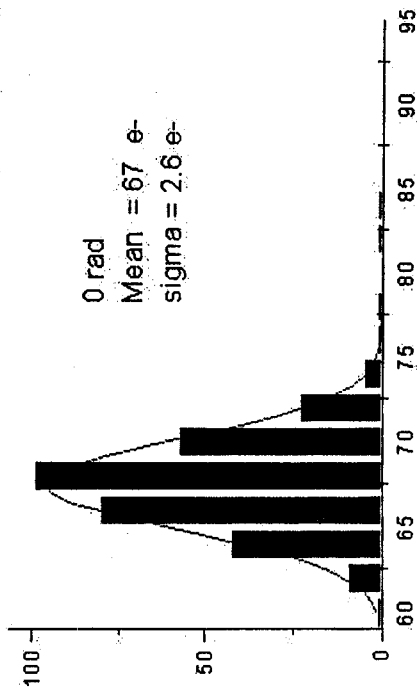
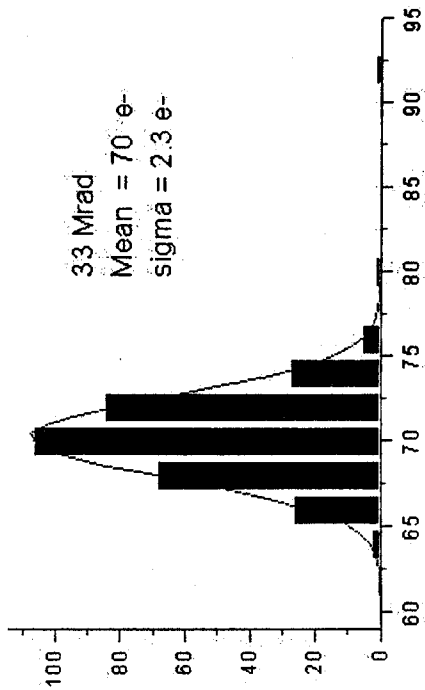
PIXELS — 'PAD' — D. CHRISTIAN
— 'COLUMN' — S. PARKER.

$10^7 - 10^9$ ELEMENTS

\uparrow
~40 layers ~4 m long ~2 m radius

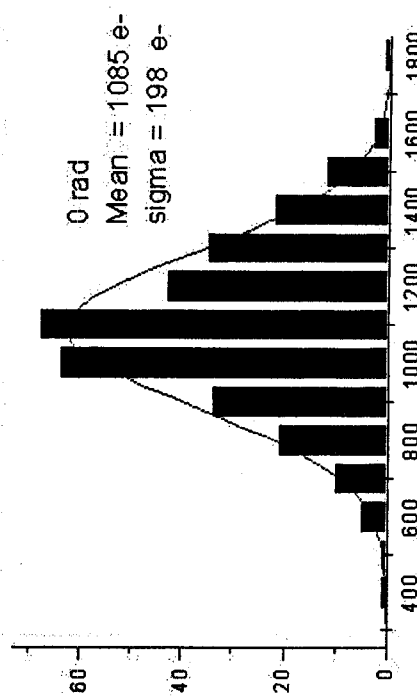
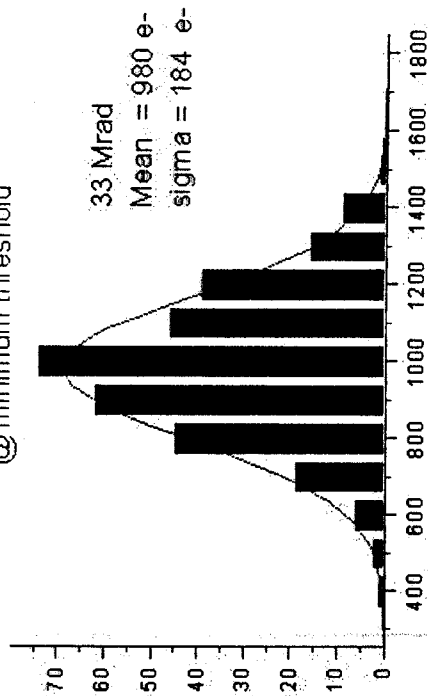
(pre)FPiX2 Noise and discriminator threshold distributions

NOISE DISTRIBUTION vs DOSE



THRESHOLD DISTRIBUTION vs DOSE

@ minimum threshold



Speed -- 3D vs. planar

1. 3D lateral cell size can be smaller than wafer thickness, so:

shorter collection distances

2. in 3D, field lines end on cylinders rather than on circles, so:

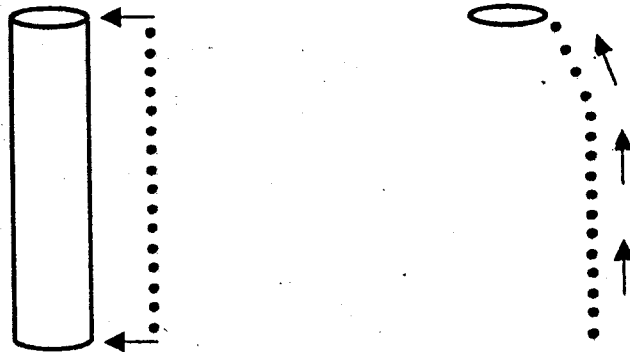
higher average fields for any given max field

(price: larger $C_{\text{electrode}}$: we now have 0.2 pF / 121 μm long electrode)

3. most of the signal is induced when the charge is close to the electrode, where the electrode solid angle is large, so:

planar signals are spread out in time as the charge arrives

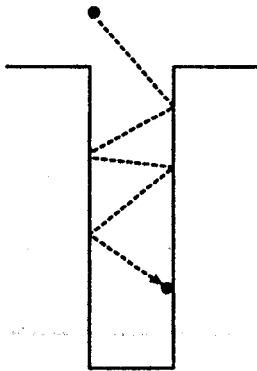
3D signals are concentrated in time as the track arrives



4. if readout chip has inputs from both n^+ and p^+ electrodes,

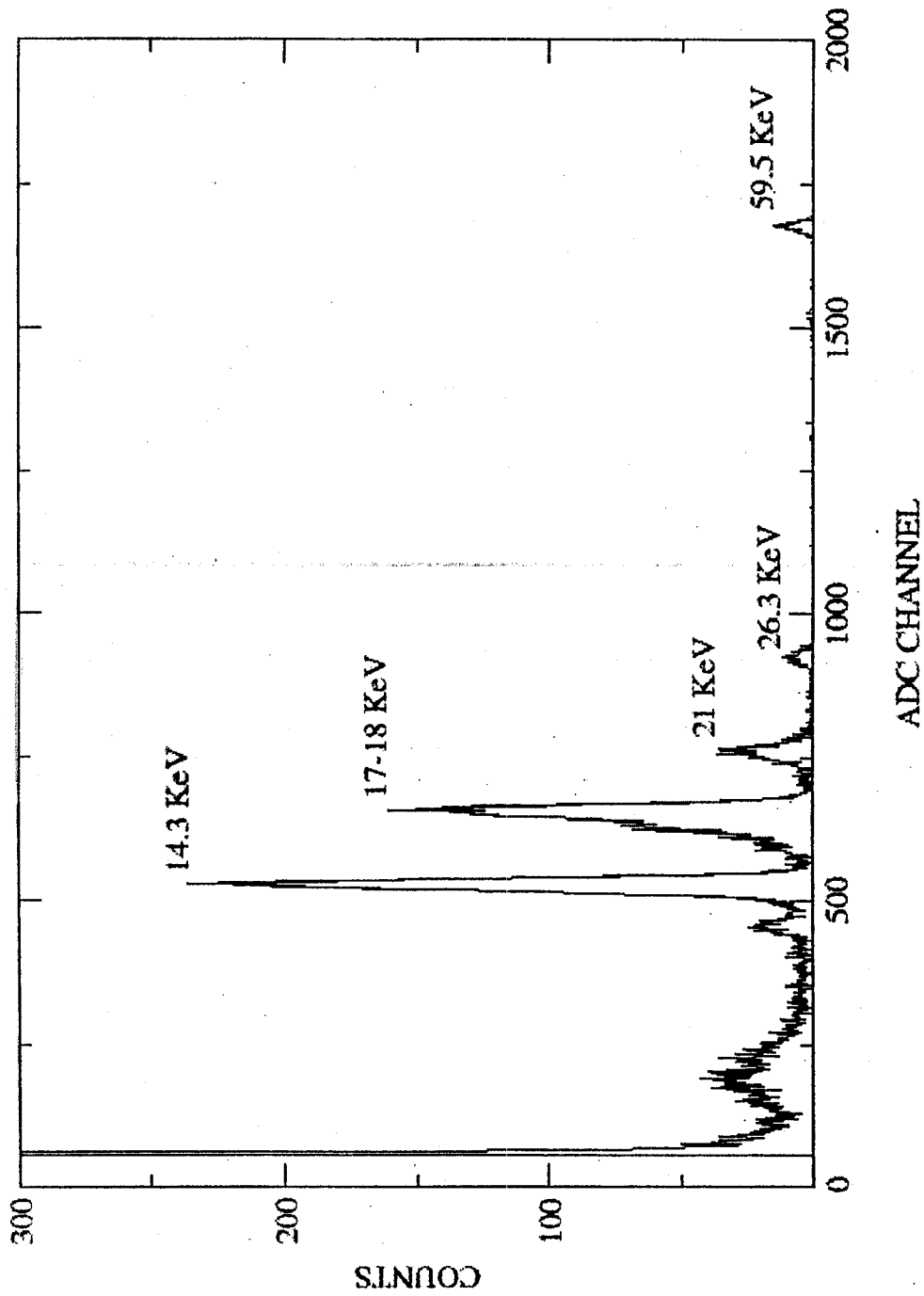
drift time correction can be made

Keys to the technology:

1. Can now etch deep, near-vertical holes.
2. With low pressure and moderate temperatures, gas molecules used to form polysilicon such as SiH_4 , SiH_2Cl_2 , SiHCl_3 , and SiCl_4 will bounce off the hole walls thousands of times before they stick. Mostly they enter, bounce a number of times, and leave. When they stick, it can be anywhere, and so a conformal coat of polysilicon is formed as the H (or Cl) leaves and the silicon migrates to a lattice site.
3. Gasses such as diborane (B_2H_6) and phosphine (PH_3) can be added to the silane. They also come out in a conformal layer, and make p^+ and n^+ doped polysilicon.
4. Heating will drive the dopants into the surrounding single crystal silicon, forming the p - n junctions and ohmic contacts in high-quality silicon, keeping electric fields away from the short-lifetime poly.

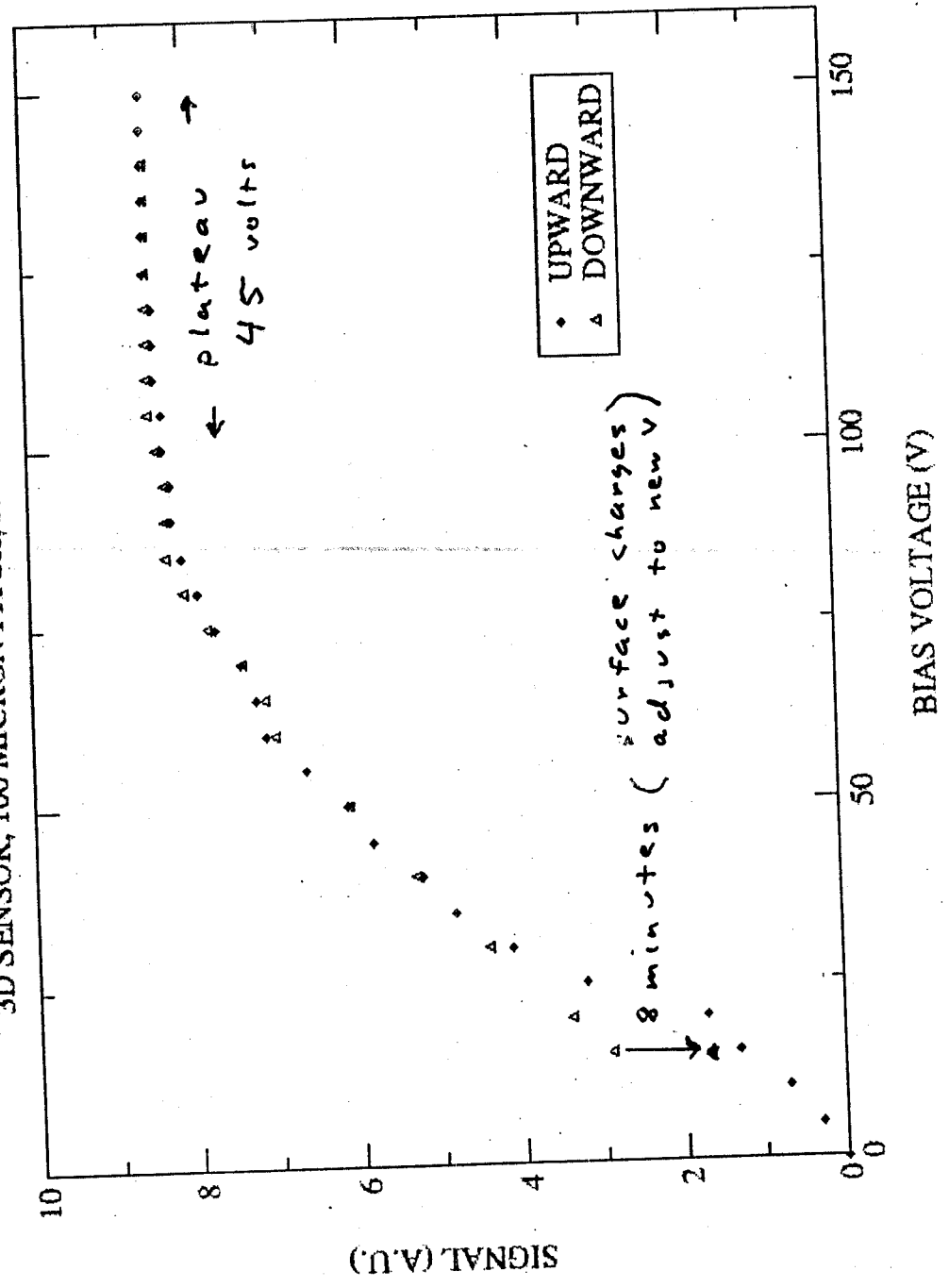
AMERICIUM-241

3D - 200 MICRON PITCH



FAKKEK

1E15 55 MeV PROTONS/CM2
3D SENSOR, 100 MICRON PITCH, ROOM TEMPERATURE



• Snowmass, July 2001

CONCLUSIONS

1. THE FIRST 3D DETECTORS HAVE BEEN SUCCESSFULLY FABRICATED
2. THEY HAVE REASONABLE LEAKAGE CURRENTS: $\frac{1}{4}$ - 1 nA PER MM^3 AT ROOM TEMPERATURE
3. THEY DEplete AT LOW VOLTAGES AS EXPECTED (TYPICALLY 5 - 10 V)
4. THEY HAVE WIDE PLATEAUS FOR INFRARED MICROBEAM SIGNALS
5. BETA, X-RAY, AND GAMMA SIGNALS HAVE BEEN SEEN
6. THE 14.3 KEV X-RAY LINE (^{241}Am SOURCE) FITS A SYMMETRIC GAUSSIAN WITH A SIGMA OF 282 eV
7. A SENSOR WITH 100-MICRON CELL SIZE HAD ITS DEPLETION VOLTAGE INCREASE FROM 5 V TO 105V (WITH A 45 V WIDE PLATEAU) AFTER IRRADIATION WITH 10^{15} - 55 MEV PROTONS / CM^2 (PRIOR TO BENEFICIAL ANNEALING)
8. ACTIVE EDGE PREVIEW: AN INFRARED MICROBEAM SCAN SHOWS A SENSOR WITH SETS OF NORMAL 3D CYLINDRICAL ELECTRODES BETWEEN TWO WALL ELECTRODES IS SENSITIVE TO WITHIN SEVERAL MICRONS OF THE WALL ELECTRODE.
9. CALCULATIONS INDICATE PULSE DURATIONS OF ABOUT 1 NS
10. THE 2nd FABRICATION RUN HAS STARTED

CALORIMETRIC TECHNIQUES

SCINTILLATOR + WLS FIBERS + PMT/APD/HPD
(SOLID / LIQUID)

WELL ESTABLISHED. RAD HARDNESS. CENTRAL.

$PbWO_4$ CRYSTALS (em) OR OTHER CRYSTALS

CMS - GOOD EM CAL, COMPROMISED JET σ

QUARTZ FIBER IN Cu (e.g.)

LOW SIGNAL BUT OK FORWARD. V. RAD HARD

SILICON PADS

RAD HARDNESS? COMPACT, SAFE

DIAMOND PADS

COST! BUT BY 2020? R&D!

HIGH PRESSURE GAS TUBES

CHEAP, V. RAD HARD, PROMISING esp. FORWARD.

LIQUID ARGON

ATLAS - SLOW ϕ 18ms

QUARTZ FIBRES (em)
+ SCINTILLATING FIBERS
(had)



MAY BE OPTIMUM FOR em & JETS, WITHOUT DEPTH SEGMENTATION.

What is important for the experiment?

- With traditional techniques:
 - If it is important to optimize electromagnetic resolution, then a **large sampling fraction** is required.
 - If it is important to optimize hadronic resolution, then a **small sampling fraction** is required.
- **These two goals exactly conflict with each other!**
- In the past, the electromagnetic energy resolution has generally been given priority.
- If an electromagnetic calorimeter has a high sampling fraction, and thus a large e/h value, then it makes sense to build a hadronic calorimeter with a similarly large e/h value.

If hadronic response and resolution is a priority...

- One should choose a **compensating** and **longitudinally unsegmented** calorimeter.
 - It should be compensating in order to eliminate fluctuations in the electromagnetic component of hadron showers (also important for jets)
 - It should be unsegmented because calibrating longitudinal segments is a **tricky business**.

Have Your Cake and Eat it, Too!!!

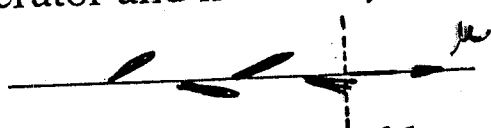
With dual readout calorimetry,

- Achieve essential advantages of compensation.
 - Eliminate dominating source of fluctuations in non-compensating calorimeter (fluctuations in f_{em}).
 - Eliminate signal nonlinearity.
- With this technique, the sampling fraction could be as large as needed for optimizing the calorimeter's resolution for electromagnetic showers, and the hadronic resolution **WOULD NOT** suffer!

Resolution of Muons

are important for search of heavy
muons for t, b quarks tagging

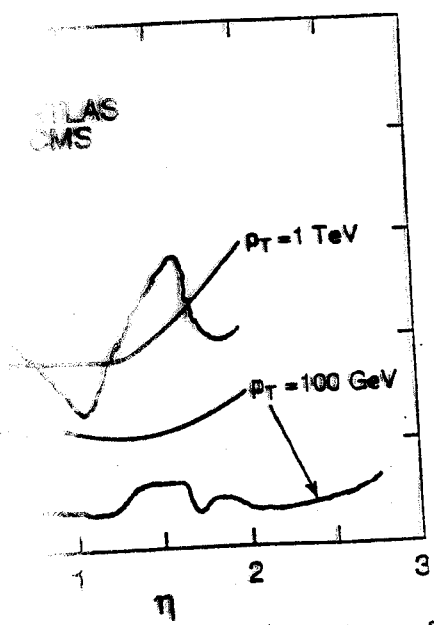
loss of a problem for muon detectors,
(from accelerator and neutrons) have



to irradiate γ (showers) and loose
ionization losses at a few hundred GeV, ^{then due to ionization}
backgrounds and requires corrections

Resolution is limited by factors similar to

$$\frac{\sigma_{det}}{L} \oplus m.s.$$



Stand alone
muon system resolution
 $\sim 10\% @ p_t \sim 1 \text{ TeV}$

Resolution of ATLAS and CMS detectors for stand-
alone function of pseudorapidity, for two values of the transverse

MUON DETECTORS ~~DO~~ → VLHC

DARIEN WOOD

BEST TO MEASURE p_μ BEFORE CALORIMETER.

MULTIPLE SCATTERING & ΔE FLUCTUATIONS
AFFECT \vec{p} (BEHIND CAL.)

& BACKGROUNDS FROM PUNCHTHRO' & DELAY OF p_μ
ALSO BEAM HALO, BACKSCATTER.

↗ ↘
DESIGN OF SHIELDING ETC.

~~DO~~ : 2 Tm, 10 POINTS, $\sigma \sim 1\text{mm}$, SIGN TO 300 GeV
WANT SIGN TO $\sim 10\text{TeV}$ (?)

COMBINATION OF : Tm ($\sim 4 \times 2 = 8\text{Tm}$?)

POINTS 40-50 ?

PRECISION $\sigma \sim 50\mu\text{m}$?

& MEASURE, FOR ISOLATED μ , SHOWERING

CRITICAL ENERGY, SHOWERING LOSSES = $\partial E_{\partial x}$
 $\sim 350\text{GeV}$ IN Fe

BIG FLUCTUATIONS IN $\frac{P_{\text{OUT}}}{P_{\text{IN}}}$ FOR $1\text{TeV}/c$ μ .
(20-30 GeV)

NO CONCLUSIONS YET

- ATLAS & CMS, IF WORK AT LHC, WOULD WORK AT $20+20 \text{ TeV } 10^{34}$ VLHC I.
- IF LHC DETECTORS SCALED UP 1.4 LINEARLY (2.7 IN VOLUME) WOULD BE FAIR STARTING POINT FOR VLHC II $100+100 \text{ TeV } 2 \cdot 10^{34}$
- \uparrow BUT : MORE RAD HARD (R&D) TECHNOLOGY
FASTER
MORE PRECISION / GRANULARITY
10 YEARS TO BUILD (2012-2022?) R&D \rightarrow 2012
\$ 0.5-1.0 B ? PER DETECTOR