

PHY481 - Lecture 28: Dielectric materials - problem solving Griffiths: Chapter 4

Boundary value problems

The boundary conditions on the displacement field and electric field across a surface that has surface free charge density σ_f and surface bound charge density σ_b are (using arguments like those used in Lecture 9),

$$D_{above}^\perp - D_{below}^\perp = \sigma_f \quad \text{so that} \quad \epsilon_{above} E_{above}^\perp - \epsilon_{below} E_{below}^\perp = \sigma_f \quad (1)$$

$$\vec{E}_{above}^\parallel - \vec{E}_{below}^\parallel = 0; \quad E_{above}^\perp - E_{below}^\perp = (\sigma_f + \sigma_b)/\epsilon_0 \quad (2)$$

The electrostatic potential is continuous across the surface.

A uniform electric field applied to a uniform dielectric sphere - Use a similar method for Problem 4.22

We consider an uncharged uniform dielectric sphere of radius a and dielectric constant ϵ , in a constant applied field

$$\vec{E}_0 = E_0 \hat{z}, \quad \text{so that} \quad V = -E_0 z = -E_0 r \cos\theta. \quad (3)$$

Since the dielectric sphere is uniform and there is no free charge, we have,

$$\vec{\nabla} \cdot \vec{D} = 0 = \vec{\nabla} \cdot (\epsilon \vec{E}) = \epsilon \vec{\nabla} \cdot \vec{E} = 0. \quad \text{uniform } \epsilon \quad (4)$$

Using $\vec{E} = -\vec{\nabla}V$ we then find that V still obeys Laplace's equation, so we try the solutions,

$$V_{int} = -C_1 r \cos\theta, \quad V_{ext} = -E_0 r \cos\theta + \frac{C_2 a^3 \cos\theta}{r^2} \quad (5)$$

We impose continuity of V , and the condition $\epsilon_0 E_n^{ext}(a, \theta) = \epsilon E_n^{int}(a, \theta)$, to find,

$$-E_0 + C_2 = -C_1; \quad -E_0 - 2C_2 = -\frac{\epsilon}{\epsilon_0} C_1 \quad (6)$$

which lead to,

$$C_1 = E_0 \frac{3}{\epsilon_r + 2}; \quad C_2 = E_0 \frac{\epsilon_r - 1}{\epsilon_r + 2} \quad (7)$$

where $\epsilon_r = \epsilon/\epsilon_0$. Using the fact that $V_{int} = -C_1 z$, we find that the electric field inside the sphere is uniform

$$\vec{E}_{int} = -\frac{\partial V_{int}}{\partial z} \hat{z} = E_0 \frac{3}{\epsilon_r + 2} \hat{z} \quad (8)$$

and that the dipole of the sphere is

$$\vec{p}_{sphere} = \frac{C_2 a^3}{k} \hat{z} = 4\pi\epsilon_0 E_0 a^3 \frac{\epsilon_r - 1}{\epsilon_r + 2} \hat{z} \quad (9)$$

The polarization density is then,

$$\vec{P} = \frac{\vec{p}_{sphere}}{volume} = \frac{\epsilon_0 E_0}{3} \frac{\epsilon_r - 1}{\epsilon_r + 2} \hat{z} \quad (10)$$

Again it is useful to compare this to the limit of a metal where $\epsilon_r \rightarrow \infty$.

Other homework hints

Problem 4.26 Use the formula for the energy density, $u = \epsilon E^2/2$, and integrate.

Problem 4.28: The energy gain when the dielectric oil enters the cylindrical cavity is $[C_{after} - C_{before}]V^2/2$. Since V is fixed, we need the expression for the capacitance of a cylindrical capacitor when part of the cylindrical cavity has dielectric susceptibility χ_e . The capacitance of a cylindrical capacitor uniformly filled with material of dielectric constant ϵ is $C = 2\pi L\epsilon/\ln(b/a)$, where L is the length of the cylinder. When the oil fills the cavity, the

capacitance is like two capacitors in parallel. The energy gain when the oil fills the cavity is then $(C(h) - C(0))V^2/2$. The total energy change when the oil enters the cylindrical cavity to height h is then $U = \delta U_{grav} - \delta U_{cap}$. The equilibrium condition is $-\vec{\nabla}U = 0$ which sets the force to zero. In this problem the gradient reduces to a derivative with respect to h .

Problem 4.32 Straightforward application of formulae.

Problem 4.34 The potential has a dipole form outside the sphere and has the form of a dipole plus a constant electric field (along the dipole axis) inside the sphere. (Give arguments as to why this is true) We then take,

$$V_{in} = (a_1 r + \frac{p}{4\pi\epsilon r^2})\cos\theta; \quad V_{out} = \frac{b_1}{r^2}\cos\theta \quad (11)$$

and use the potential and displacement field boundary conditions to find a_1 and $b - 1$.

Problem 4.40 Much of this is done in Lecture 26 of the online notes.