

Resonance terms in $R(\sqrt{s})$

In examining the graph of the R value for $e^+e^- \rightarrow \text{hadrons}$, which shows a region from about 1 GeV to 3 GeV that is consistent with light quark production, there is region at lower energy with several peaks. These peaks occur because of vector boson production. The ρ , ω and ϕ vector bosons

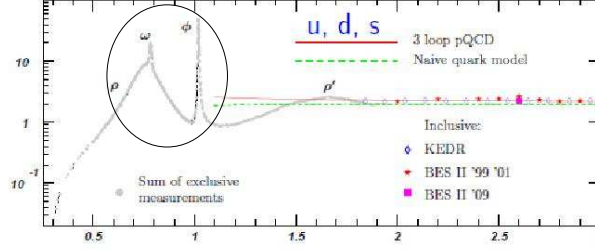


Figure 1: The R value for $\sqrt{s} \leq 3.0$ GeV is shown.

have the same parity and charge conjugation quantum numbers as the (virtual) photon produced in e^+e^- annihilation and therefore couple to it. For arbitrary center of mass energy \sqrt{s} , this coupling can produce a virtual vector boson with very small probability. However, when \sqrt{s} is near the vector boson's physical mass there is a high probability and a huge spike in the cross section. All these vector mesons decay by virtue of the strong interaction. What is seen in the detector at these values of \sqrt{s} is the decay products of the vector meson that have the invariant mass of the vector meson. These decay products predominantly consist of two or three pions or a pair of kaons.

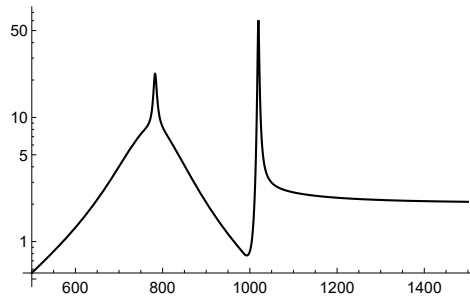


Figure 2: The effect of the ρ , ω and ϕ resonances on the $R(\sqrt{s})$ ratio is shown.

In addition to their dominant hadronic decay modes, all three vector bosons have a small decay width $\Gamma_{e^+e^-}$ into e^+e^- . The strength of the photon-vector boson coupling, f_V , is determined from measuring $\Gamma_{e^+e^-}$ and using the relation

$$\Gamma_{e^+e^-} = \frac{4\pi\alpha^2 m_V}{3f_V^2}. \quad (1)$$

In a similar way, the hadronic decay strength g_V can be obtained from its total width Γ_V . For the ρ meson, which decays almost exclusively into $\pi^+\pi^-$, Γ_ρ is

$$\Gamma_\rho = \frac{g_\rho^2(m_\rho^2 - 4m_\pi^2)^{3/2}}{48\pi m_\rho^2} = 149.1 \text{ MeV}. \quad (2)$$

Figure [2] shows the theoretical prediction for this resonance region. The primary decay modes are $\rho \rightarrow \pi^+\pi^-$, $\omega \rightarrow \pi^+\pi^-\pi^0$ and $\phi \rightarrow K\bar{K}$. In addition, both the ω and the ϕ have measurable $\pi^+\pi^-$ decay modes. This means that the ρ , ω and ϕ have $\pi^+\pi^-$ decay amplitudes that interfere with one another. The square of the scattering amplitude has the form

$$|\mathcal{M}|^2 = |M_\rho(\pi^+\pi^-) + M_\omega(\pi^+\pi^-) + M_\phi(\pi^+\pi^-)|^2 + |M_\omega(\pi^+\pi^-\pi^0)|^2 + |M_\phi(K\bar{K})|^2, \quad (3)$$

and each contribution is multiplied by the appropriate phase space factor. A typical vector boson amplitude looks like

$$M_\rho(\pi^+\pi^-) = \frac{e^2 m_\rho^2}{f_\rho s} \frac{\bar{v}(p_2)\gamma_\mu u(p_1)H_\mu(k_1, k_2, \dots)}{s - m_\rho^2 + im_\rho\Gamma_\rho}, \quad (4)$$

where $H_\mu(k_1, k_2, \dots)$ contains the kinematics of final state. The appearance of the factors like

$$\frac{1}{(s - m_\rho^2)^2 + m_\rho^2\Gamma_\rho^2}, \quad (5)$$

in the squared amplitude produces the peak when s is in the vicinity of m_ρ^2 .