

# PHY294H

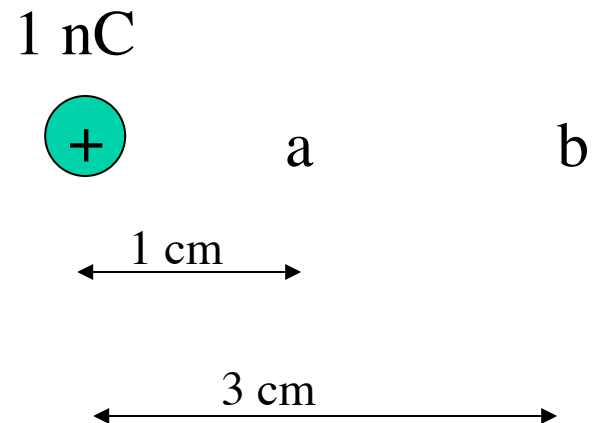
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- Professor: Joey Huston
- email: [huston@msu.edu](mailto:huston@msu.edu)
- office: BPS3230
- Homework will be with Mastering Physics (and an average of 1 hand-written problem per week)
  - ◆ **Added problem 28.68 for 3<sup>rd</sup> MP assignment due Wed Feb. 3 as a hand-in problem**
  - ◆ **Help-room hours: 12:40-2:40 Tues; 3:00-4:00 PM Friday**
- Quizzes by iclicker (sometimes hand-written)
- Course website: [www.pa.msu.edu/~huston/phy294h/index.html](http://www.pa.msu.edu/~huston/phy294h/index.html)
  - ◆ lectures will be posted frequently, mostly every day if I can remember to do so

# Another example before we leave this chapter

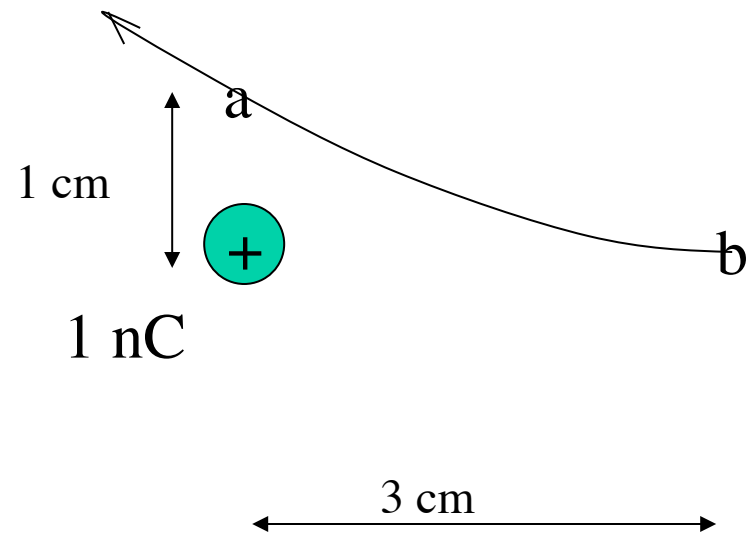
- What is:

- ◆ the potential at points a and b
- ◆ the potential difference between a and b
- ◆ the potential energy of a proton at a and b
- ◆ the speed at point b of a proton that was moving to the right at point a with a speed of  $4 \times 10^5$  m/s
- ◆ the speed at point a of a proton that was moving to the left at point b with a speed of  $5.3 \times 10^5$  m/s



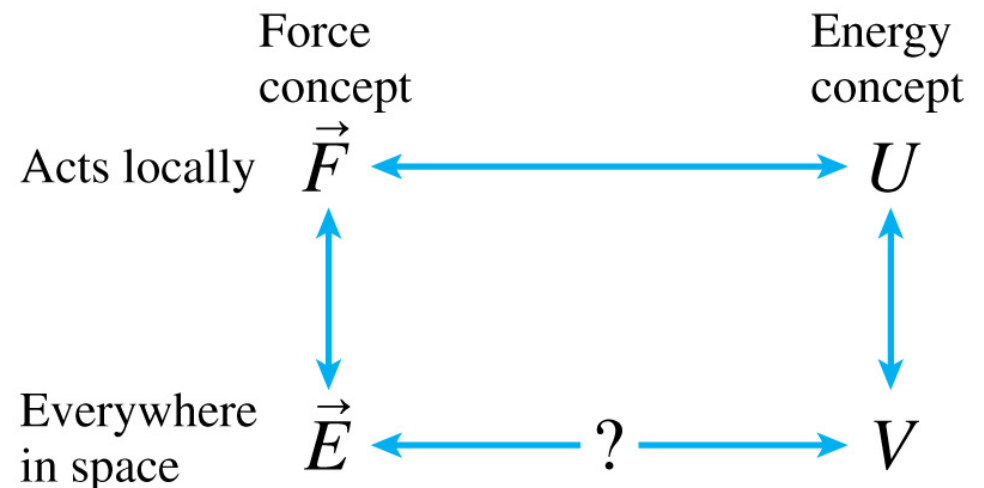
## Another example before we leave this chapter

- Suppose the proton isn't heading towards the positive charge head-on, but at a small angle
- Then the path travelled will look like that on the right
- Suppose the proton's speed at point b is  $5.3 \times 10^5$  /ms
- What is the speed at point a?



# Electric potential and electric field

- We've talked so far about the connections between force and potential energy, between force and the electric field, and between the potential and the potential energy
- Now in this chapter, we'll study the relationship between the electric field and the electric potential



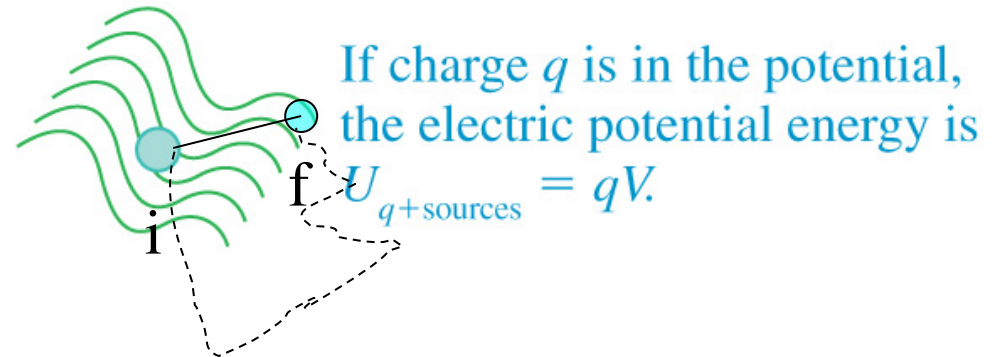
# Potential

- We can write the change in potential energy in terms of the work done by the field

$$\Delta U = -W(i \rightarrow f) = -\int_{s_i}^{s_f} \vec{F} \cdot d\vec{s}$$

- But we can also write the potential energy as  $U=qV$  and the force as  $F=qE$
- Thus, we can write the potential difference in terms of an integral over the electric field

$$\Delta V = V(s_f) - V(s_i) = -\int_{s_i}^{s_f} \vec{E} \cdot d\vec{s} = -\int_{s_i}^{s_f} E_s ds$$

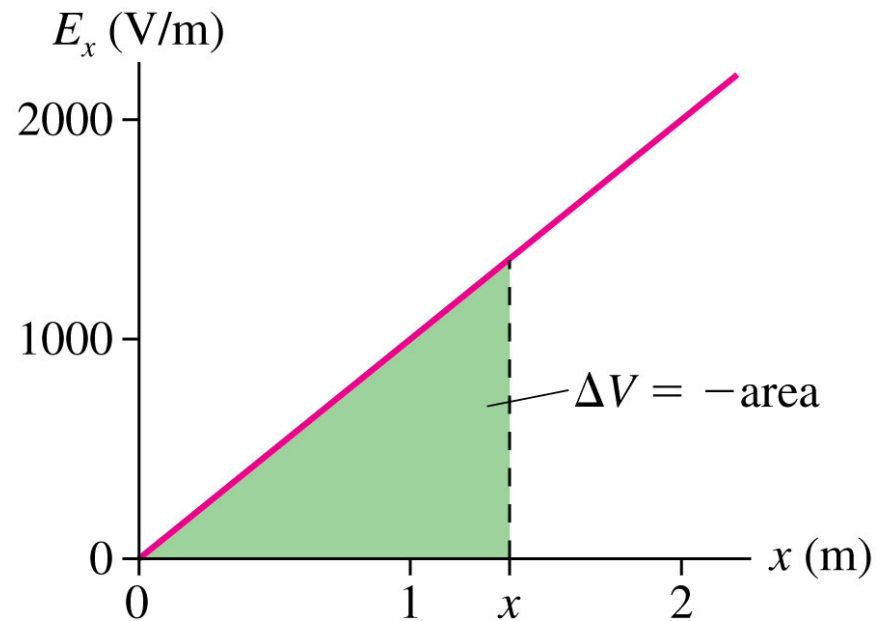


Because electrostatic forces are conservative, the potential difference is independent of the path taken

For a uniform electric field, we can write simply  $\Delta V = -E_s \Delta s$

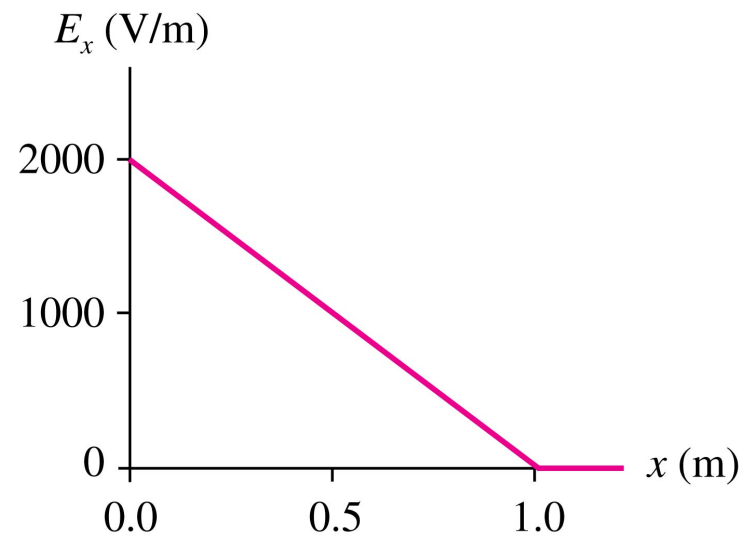
- The electric field is the negative of the area given in green (the integral of the electric field over  $x$ )

$$\Delta V = V(s_f) - V(s_i) = -\int_{s_i}^{s_f} E ds = -\int_{s_i}^{s_f} E_s ds$$

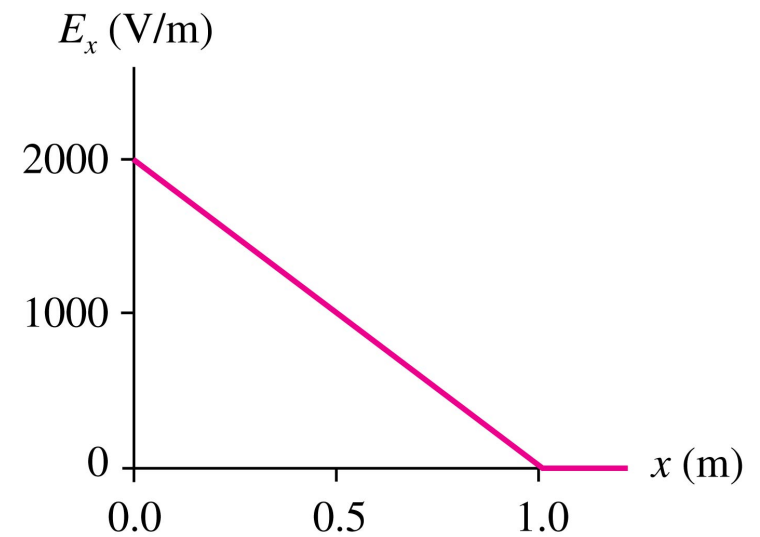


This is a graph of the  $x$ -component of the electric field along the  $x$ -axis. The potential is zero at the origin. What is the potential at  $x = 1\text{m}$ ?

- A. 2000 V.
- B. 1000 V.
- C. 0 V.
- D. -1000 V.
- E. -2000 V.



This is a graph of the  $x$ -component of the electric field along the  $x$ -axis. The potential is zero at the origin. What is the potential at  $x = 1\text{m}$ ?



- A. 2000 V.
- B. 1000 V.
- C. 0 V.
- ✓ D. **-1000 V.**     $\Delta V = -\text{area under curve}$
- E. -2000 V.



# Potential of a point charge

- Let's first define the reference potential to be 0 at  $r=\text{infinity}$

- And then write

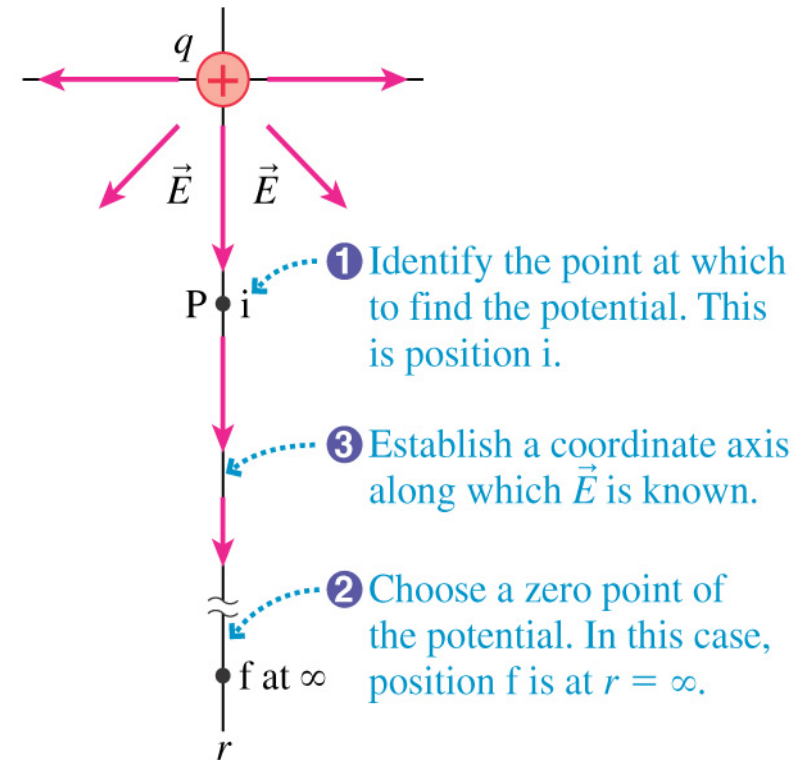
$$\Delta V = V(\text{infinity}) - V(r) = - \int_r^{\text{infinity}} E_r dr$$

- We know the electric field from a point charge

$$E_r = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

- So we can then write

$$V(r) = V(\text{infinity}) + \frac{q}{4\pi\epsilon_0} \int_r^{\text{infinity}} \frac{1}{r^2} dr = V(\text{infinity}) + \frac{q}{4\pi\epsilon_0} \left. \frac{-1}{r} \right|_r^{\text{infinity}} = \frac{1}{4\pi\epsilon_0} \frac{q}{r}$$

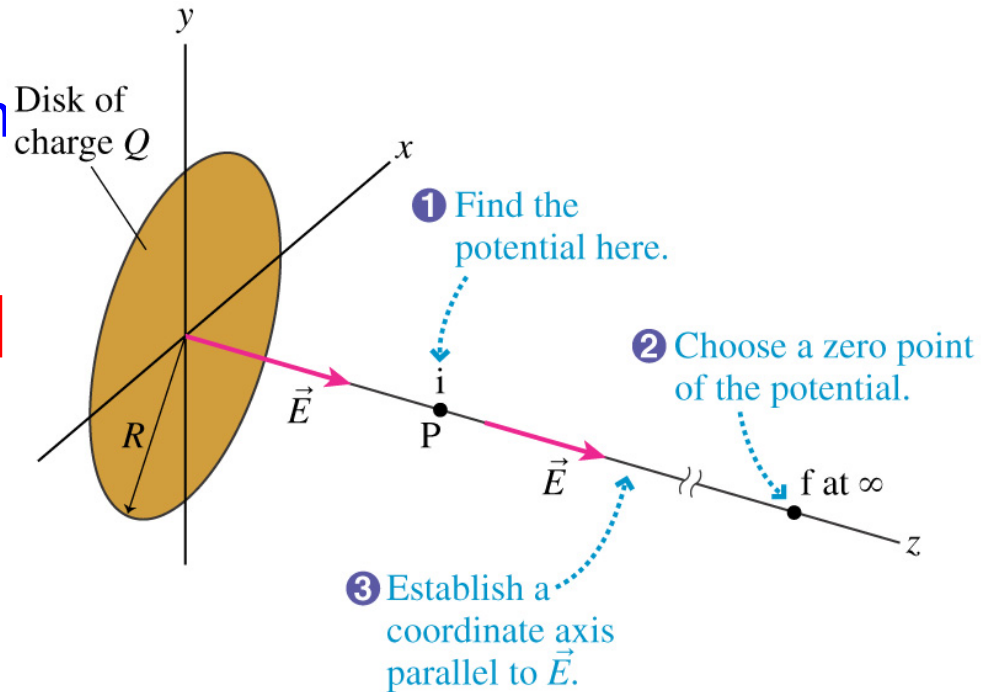


# Let's go back to the charged disk

- We found that the electric field had the form of

$$E = \frac{Q}{2\pi R^2 \epsilon_0} \left[ 1 - \frac{z}{(z^2 + R^2)^{1/2}} \right]$$

- Can we calculate the potential from the field?



$$V(z) = V(\infty) + \int_z^\infty E_z(z) dz$$

$$= 0 + \frac{Q}{2\pi R^2 \epsilon_0} \int_z^\infty \left[ 1 - \frac{z}{(z^2 + R^2)^{1/2}} \right]$$

$$\text{let } u = z^2 + R^2$$

$$\begin{aligned} V(z) &= \frac{Q}{2\pi R^2 \epsilon_0} \left[ \int dz - \int \frac{\frac{1}{2} du}{u^{1/2}} \right]_z^\infty = \frac{Q}{2\pi R^2 \epsilon_0} \left[ z - \sqrt{z^2 + R^2} \right]_z^\infty \\ &= \frac{Q}{2\pi R^2 \epsilon_0} \left[ \sqrt{z^2 + R^2} - z \right] \end{aligned}$$

# Potential from a uniform line of charge

- What is the potential from an infinite line of charge with charge density  $\lambda$  as a function of  $R$ , the distance perpendicular to the line of charge?
- We know the electric field from an infinite line of charge; it has only a radial component

$$E = \frac{\lambda}{2\pi\epsilon_0 R}$$

$$\Delta V = -\int E_r dr = -\frac{\lambda}{2\pi\epsilon_0} \int \frac{dr}{r}$$

$$\Delta V = -\frac{\lambda}{2\pi\epsilon_0} \int_a^R \frac{dr}{r} = -\frac{\lambda}{2\pi\epsilon_0} \ln r \Big|_a^R$$

$$\Delta V = -\frac{\lambda}{2\pi\epsilon_0} \ln \frac{R}{a}$$

y

x x  
a R

Let me choose the potential to be zero at  $r=a$

Note that I can't set  $V$  to be zero at infinity since the log of infinity is infinite

Since the line is infinite, we can't get *far away* from it

# Electric field from the potential

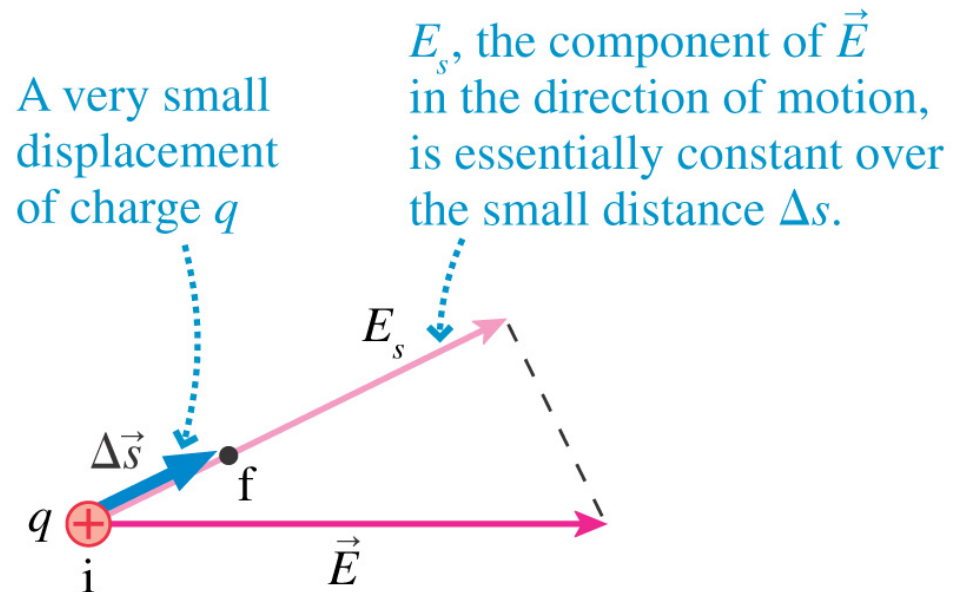
- We saw earlier that

- ◆  $\Delta V = -E_s \Delta s$
- ◆ Thus,  $E_s = -\Delta V/\Delta s$
- ◆ ...in the  $s$  direction

- In the limit that  $\Delta s$  goes to zero, we can write

- ◆  $E_s = -dV/ds$
- ◆ so for  $s = r$ , we can write for a point charge

$$E_r = -\frac{dV}{dr} = -\frac{d}{dr}\left(\frac{1}{4\pi\epsilon_0} \frac{q}{r}\right) = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}$$



# Electric field from the potential

- In general...

$$d\vec{s} = dx\vec{i} + dy\vec{j} + dz\vec{k}$$

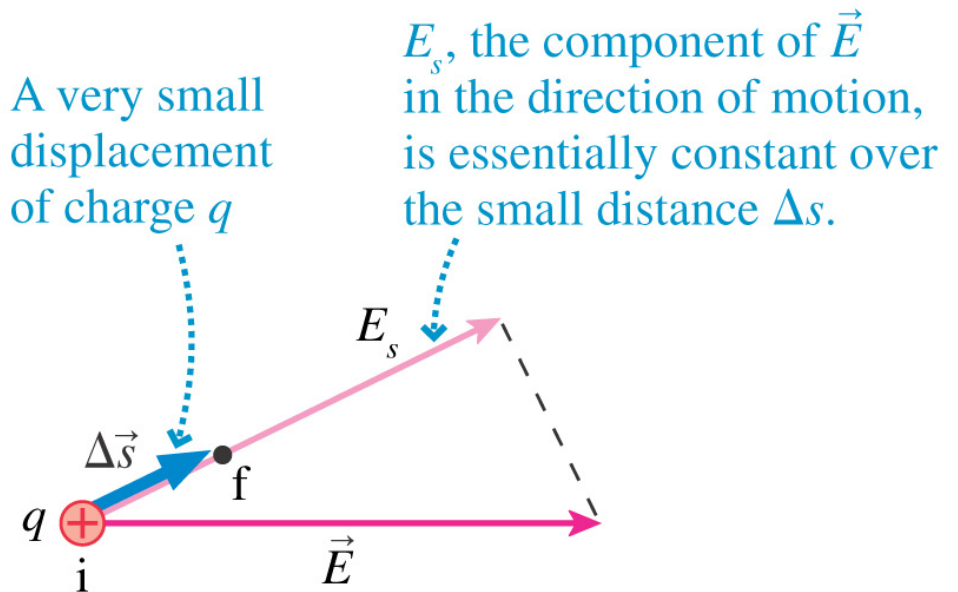
$$dV = -Eds = -E_x dx - E_y dy - E_z dz$$

- So we can write

$$dV = \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz$$

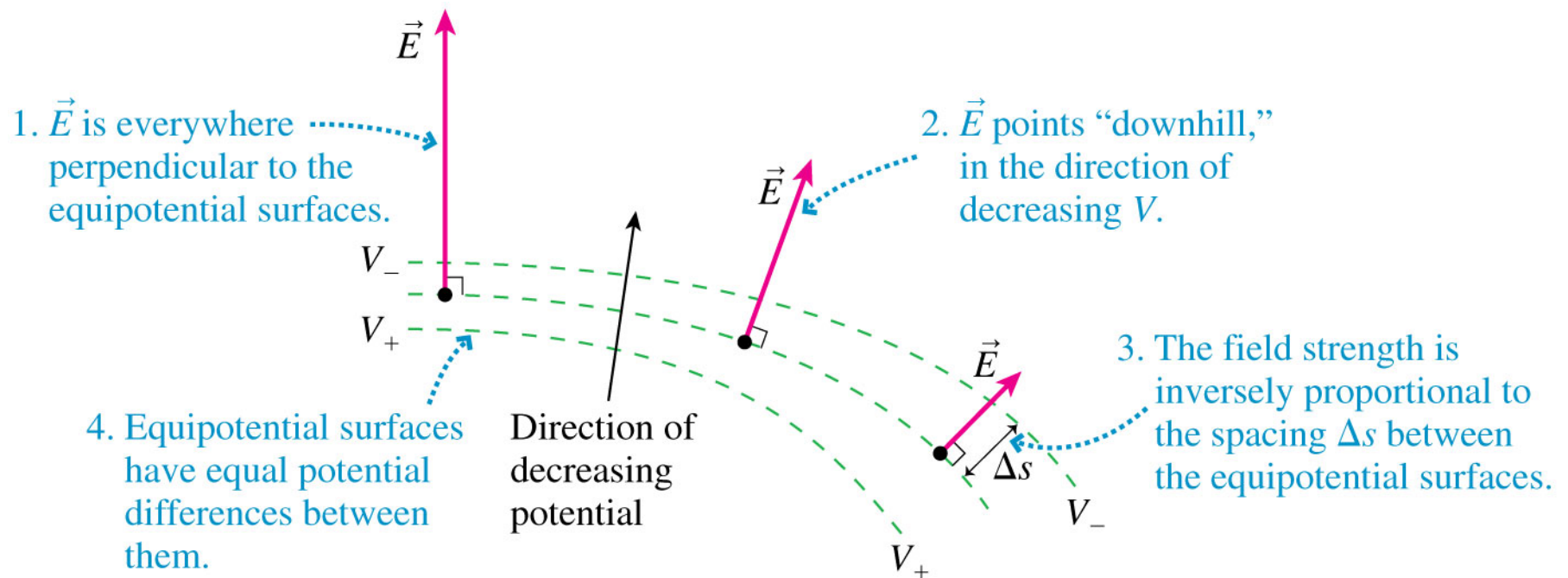
- and

$$E_x = -\frac{\partial V}{\partial x}, E_y = -\frac{\partial V}{\partial y}, E_z = -\frac{\partial V}{\partial z}$$



In three dimensions, we can find the electric field from the electric potential as:

$$\vec{E} = E_x \hat{i} + E_y \hat{j} + E_z \hat{k} = -\left( \frac{\partial V}{\partial x} \hat{i} + \frac{\partial V}{\partial y} \hat{j} + \frac{\partial V}{\partial z} \hat{k} \right)$$



# Example

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- Suppose  $V = Axy^2 - Byz$
- What is  $E$ ?

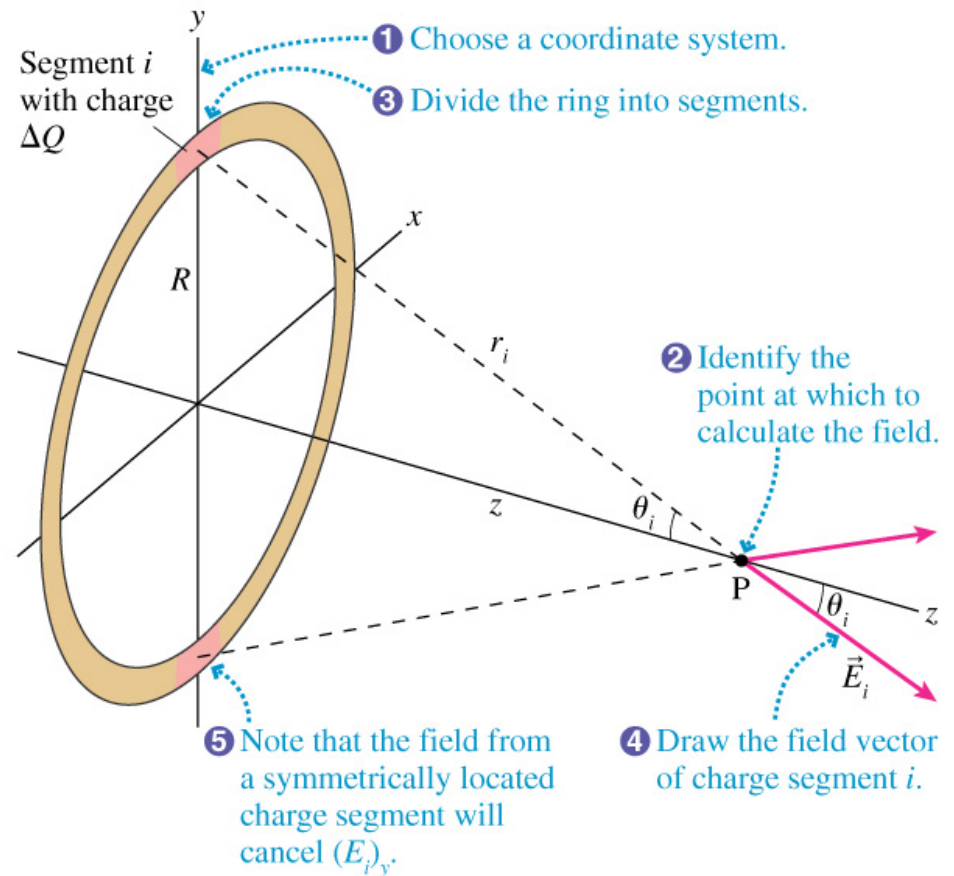
# Electric field from a charged ring

- We know the potential is

$$V_{ring} = \frac{1}{4\pi\epsilon_0} \frac{Q}{\sqrt{z^2 + R^2}}$$

$$E = E_z = -\frac{dV}{dz} = -\frac{d}{dz} \left( \frac{1}{4\pi\epsilon_0} \frac{Q}{\sqrt{z^2 + R^2}} \right)$$

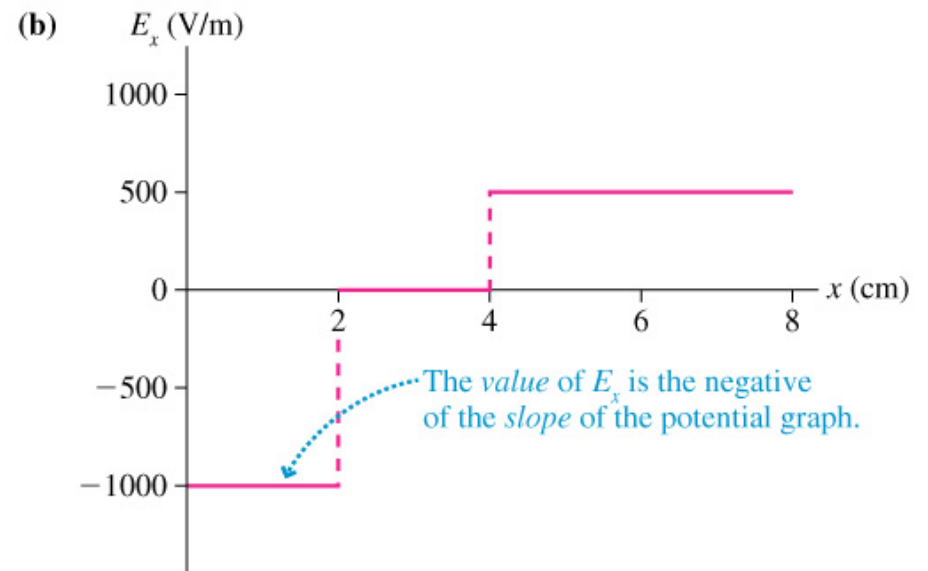
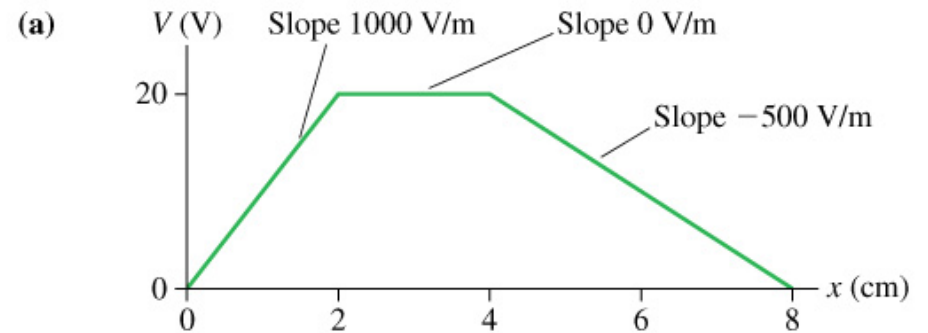
$$E_{ring} = \frac{1}{4\pi\epsilon_0} \frac{zQ}{(z^2 + R^2)^{3/2}}$$





# Example

- $E_s = -dV/ds$
- ...or in this case
- $E_x = -dV/dx$

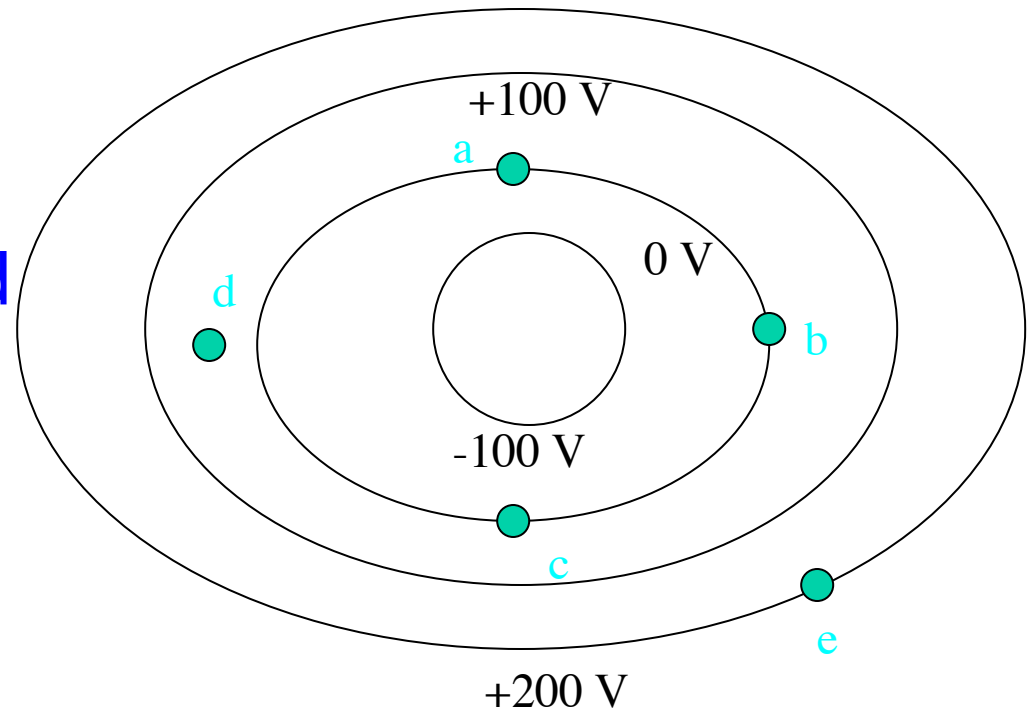


# Useful table

ELECTRIC FIELDS AND POTENTIALS FOR VARIOUS CHARGE CONFIGURATIONS			
Charge Configuration	Magnitude of Electric Field	Electric Potential	Location of Zero Potential
Point charge	$\frac{q}{4\pi\epsilon_0 r^2}$	$\frac{q}{4\pi\epsilon_0 r}$	$\infty$
Infinite line of uniform charge density $\lambda$	$\frac{\lambda}{2\pi\epsilon_0 r}$	$-\frac{\lambda}{2\pi\epsilon_0} \ln \frac{r}{a}$	$r = a$
Parallel, oppositely charged plates of uniform charge density $\sigma$ , separation $d$	$\frac{\sigma}{\epsilon_0}$	$\Delta V = -Ed = -\frac{\sigma d}{\epsilon_0}$	Anywhere
Charged disk of radius $R$ , along axis at distance $x$	$\frac{Q}{2\pi\epsilon_0} \left( \frac{\sqrt{R^2 + x^2} - x}{\sqrt{R^2 + x^2}} \right)$	$\frac{Q}{2\pi\epsilon_0 R^2} (\sqrt{R^2 + x^2} - x)$	$\infty$
Charged spherical shell of radius $R$	$r \geq R: \frac{Q}{4\pi\epsilon_0 r^2}$	$r \geq R: \frac{Q}{4\pi\epsilon_0 r}$	$\infty$
	$r < R: 0$	$r \leq R: \frac{Q}{4\pi\epsilon_0 R}$	$\infty$
Electric dipole	Along bisecting axis only, far away:	Everywhere, far away:	$\infty$
	$\frac{p}{4\pi\epsilon_0 r^3}$	$\frac{p \cos \theta}{4\pi\epsilon_0 r^2}$	
Charged ring of radius $R$ , along axis	$\frac{Qx}{4\pi\epsilon_0 (R^2 + x^2)^{3/2}}$	$\frac{Q}{4\pi\epsilon_0 \sqrt{R^2 + x^2}}$	$\infty$
Uniformly charged nonconducting solid sphere of radius $R$	$r \geq R: \frac{Q}{4\pi\epsilon_0 r^2}$	$r \geq R: \frac{Q}{4\pi\epsilon_0 r}$	$\infty$
	$r < R: \frac{Qr}{4\pi\epsilon_0 R^3}$	$r < R: \frac{Q}{8\pi\epsilon_0} \left( 3 - \frac{r^2}{R^2} \right)$	$\infty$

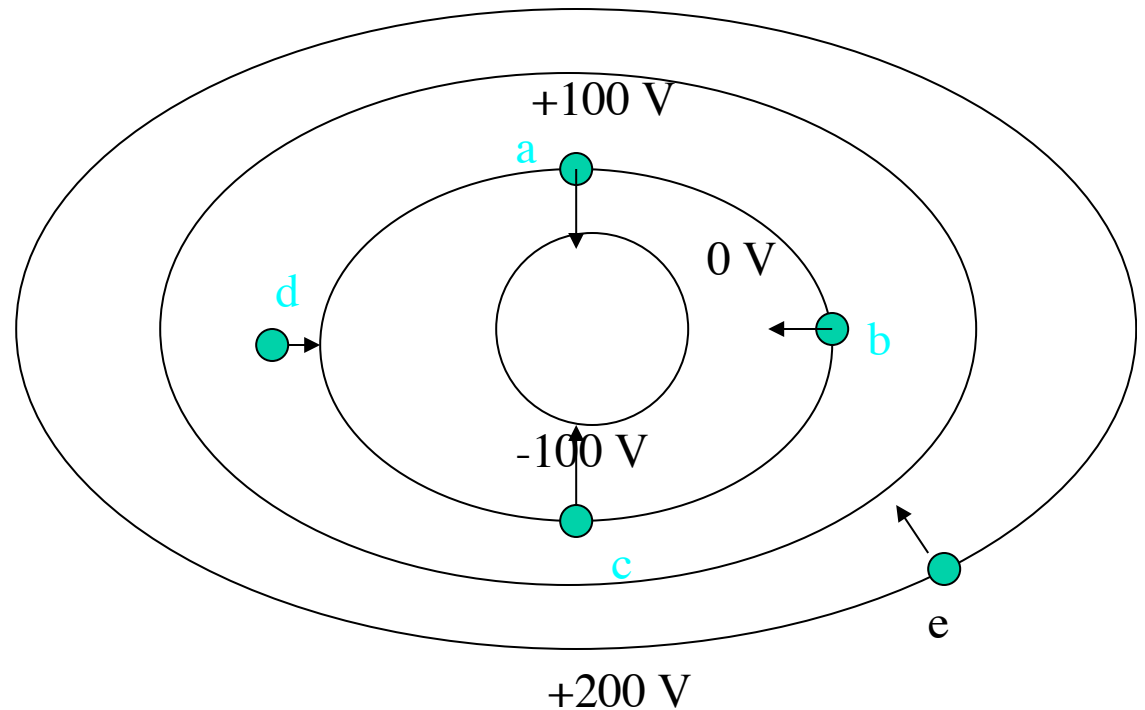
# Equipotential surfaces

- Suppose I have the equipotential surfaces shown to the right
- Let me label some points
- What are the E field directions?
- Is  $E_a = E_b$ ?
- What about  $E_c$  and  $E_d$ ?



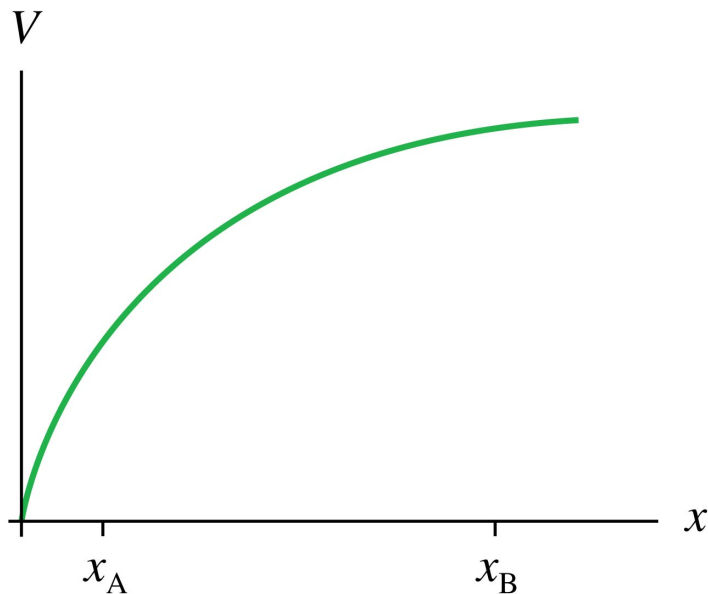
# Equipotential surface

- $E_a$  and  $E_c$  look to be about the same magnitude
- $E_b < E_{a,c}$
- $E_b > E_d$



At which point is the electric field stronger?

- A. At  $x_A$ .
- B. At  $x_B$ .
- C. The field is the same strength at both.
- D. There's not enough information to tell.



At which point is the electric field stronger?

- ✓ A. At  $x_A$ .  $|E| = \text{slope of potential graph}$
- B. At  $x_B$ .
- C. The field is the same strength at both.
- D. There's not enough information to tell.

