# Summary of the basic things covered in the lecture (Phy 410) for Exam 2, April 2, 2010

#### **Chapter 4**

1. Planck Distribution function (photon gas)

Thermal average number of photons and average energy in a single mode of frequency  $\omega$  and energy  $\hbar \omega$  are given by

$$\langle s \rangle = \frac{1}{e^{\hbar \omega/\tau} - 1}; \langle \varepsilon \rangle = \frac{\hbar \omega}{e^{\hbar \omega/\tau} - 1}$$

- 2. Counting and summing of modes in a cavity (cube of volume V=L<sup>3</sup>);  $\vec{n} = (n_x, n_y, n_z)$ .
- 3. The frequency of EM waves in a cavity  $\omega_{\bar{n}} = \frac{\pi}{L} cn; n = \sqrt{n_x^2 + n_y^2 + n_z^2}$  (3d)
- 4. Review how you will do calculations for 2 dimensional photon and phonon gas (see 6 below).

In 2-dimensions one has to sum over  $\vec{n} = (n_x, n_y)$  and replace volume V by area

A= L<sup>2</sup>, and 
$$\omega_{\bar{n}} = \frac{\pi}{L} cn; n = \sqrt{n_x^2 + n_y^2}$$

- 5. For photon gas in 3d review energy density u=U/V, spectral density  $u_{\omega}$ , entropy, and energy flux density  $J_U = \sigma_B T^4$  (Note it is  $T \operatorname{not} \tau$ ).
- 6. Phonons in a solid. Debye law for heat capacity at low T, concept of Debye temperature. In 3-dimensions

$$U(\tau) \approx \frac{3\pi^4 N \tau^4}{5(k_B \theta)^3} \text{ for } \tau \ll k_B \theta \text{ or } T \ll \theta$$
$$C_V = \frac{12\pi^4 N}{5} \left(\frac{\tau}{k_B \theta}\right)^3; \tau^3 \text{ or } T^3 \text{ law}$$

In 2-dimensions it is  $T^2$  law and in 1-dimension it is T law.

## Chapter 5

7. Chemical potential

$$\mu(\tau, V, N) = \left(\frac{\partial F}{\partial N}\right)_{\tau, V}$$

For classical ideal gas (no internal degrees of freedom), calculate F and obtain

$$\mu = \mu_{ext} = \tau \ln(n / n_Q);$$
  
In 3-dimension  $n = N / V; n_Q = \frac{1}{\lambda_{th}^3};$   
In 2-dimension  $n = N / A; n_Q = \frac{1}{\lambda_{th}^2}$   
Where thermal wave length  $\lambda_{th} = \left(\frac{2\pi \lambda^2}{M\tau}\right)^{1/2}$ 

When there is internal degrees of freedom

$$\mu_{tot} = \mu_{ext} + \mu_{int}$$

#### For diffusive equilibrium the total chemical potential must be constant.

8. For an open system with a given temperature  $\tau$  and given chemical potential  $\mu$ . Gibbs factor, Gibbs sum, Grand Sum, Grand Partition function, average number of particles in systems where the particle number can fluctuate.

$$P(N, \mathcal{E}_{s(N)}) = \frac{e^{[N\mu - \mathcal{E}_{s(N)}]/\tau}}{\widetilde{Z}} = \frac{\lambda^{N} e^{-\mathcal{E}_{s(N)}/\tau}}{\widetilde{Z}}$$
$$\widetilde{Z} = \sum_{N} \sum_{s(N)} e^{[N\mu - \mathcal{E}_{s(N)}]/\tau} = \sum_{N} \lambda^{N} \sum_{s(N)} e^{-\mathcal{E}_{s(N)}/\tau} = \sum_{N} \lambda^{N} Z(N, \tau, V)$$
$$P(N) = \frac{\lambda^{N} Z(N, \tau, V)}{\widetilde{Z}}$$

Average number of particles <N>

$$< N >= \sum_{N} NP(N) = \frac{\sum_{N} N\lambda^{N} Z(N, \tau, V)}{\widetilde{Z}} = \frac{\sum_{N} N \sum_{s(N)} e^{[N\mu - \varepsilon_{s(N)}]}}{\widetilde{Z}} = \tau \frac{\partial \ln \widetilde{Z}}{\partial \mu} = \lambda \frac{\partial}{\partial \lambda} \ln \widetilde{Z}$$

where absolute activity  $\lambda$  is related to the hemical potential  $\mu$  and  $\tau$ i.e.  $\lambda = e^{\mu/\tau}$ 

Note: The notation for Gibbs Sum  $\widetilde{Z}$  is different from that used in the lecture and the book, Sorry!) Review the pressure dependence of binding of atoms to surfaces, myoglobin, hemoglobin etc.

## **Chapter 6**

9. Ideal Gas, Fermions, Bosons, Classical regime Fermi Dirac distribution:  $f_{FD}(\varepsilon) = \frac{1}{e^{(\varepsilon-\mu)/\tau} + 1}$ Bose Einstein distribution:  $f_{BE}(\varepsilon) = \frac{1}{e^{(\varepsilon-\mu)/\tau} - 1}$ For very low densities  $n \ll n_O \ OR \ v \gg \lambda_{th}^3$ 

$$f_{FD}(\varepsilon) = f_{BE}(\varepsilon) = e^{-(\varepsilon - \mu)/\tau} = \lambda e^{-\varepsilon/\tau}$$

- 10. Obtain classical ideal gas thermal properties starting from the chemical potential as a function of  $N = \langle N \rangle, V, \tau$  and using thermodynamic relations (see below).
- 11. Thermodynamic identity and how to use it to calculate changes in enthalpy H and Helmholtz free energy F, and F itself.

$$\tau \, d\sigma = dU + p \, dV - \mu \, dN$$

$$H = U + pV; dH = \tau \, d\sigma + V \, dp + \mu \, dN; H(\sigma, p, N)$$

$$F = U - \tau\sigma; dF = -\sigma \, d\tau - p \, dV + \mu \, dN; F(\tau, V, N)$$
The Chemical potential  $\mu(\tau, N, V) = \left(\frac{\partial F}{\partial N}\right)_{\tau, V}$ 

$$F(N, \tau, V) = \int_{0}^{N} \mu(\tau, N', V) dN'$$

12. For system in 2 and 1 dimensions, one has to extend the thermodynamic identity and the definitions of some of the free energies.

$$\tau d\sigma = dU + \pi dA - \mu dN (2d)$$
  

$$\tau d\sigma = dU + TdL - \mu dN(1d)$$
  

$$H = U + \pi A (2d); H = U + TL (1d)$$
  

$$F = U - \tau \sigma$$

In the above equation  $\pi$  is called the spreading pressure and T is called the tension.