

## HW#11

### 11.1

(i)

The Hydrogen atom consists of one electron and one proton. They both have  $1/2$  spin. Thus, the spin of Hydrogen atom is  $1/2+1/2=1$ . We can treat it as a boson.

(ii)

Let the average distance to be  $d$ , every atom take a average volume of a sphere with radius  $r = d / 2$ , thus

$$N \times \frac{4}{3} \pi r^3 = V \Rightarrow \frac{N}{V} \times \frac{4}{3} \pi r^3 = 1$$

For  $\frac{N}{V} = 1.8 \times 10^{14} \text{ atoms} / \text{cm}^3$ , we can obtain

$$r = 1.1 \times 10^{-5} \text{ cm} \Rightarrow d = 2r = 2.2 \times 10^{-5} \text{ cm}$$

(iii)

We have

$$\frac{N}{V} = 1.8 \times 10^{14} / \text{cm}^3 = 1.8 \times 10^{20} / \text{m}^3$$

$$M = 1.673 \times 10^{-27} \text{ Kg}$$

From (7.72), we can obtain

$$\begin{aligned} \tau_E &= \frac{2\pi\hbar^2}{M} \left( \frac{N}{2.612V} \right)^{2/3} \\ \Rightarrow T_E &= \frac{2\pi\hbar^2}{k_B M} \left( \frac{N}{2.612V} \right)^{2/3} = 5.09 \times 10^{-5} \text{ K} = 50.9 \mu\text{K} \end{aligned}$$

The result is very close to the measured value of  $50 \mu\text{K}$ .

### 11.2

We have

$$M = 6.697 \times 10^{-27} \text{ Kg}$$

$$\frac{N}{V} = \frac{M_{\text{all}} / M}{M_{\text{all}} / \rho} = \frac{\rho}{M} = 2.165 \times 10^{28} / \text{m}^3$$

From (7.72), we can obtain

$$\begin{aligned} \tau_E &= \frac{2\pi\hbar^2}{M} \left( \frac{N}{2.612V} \right)^{2/3} \\ \Rightarrow T_E &= \frac{2\pi\hbar^2}{k_B M} \left( \frac{N}{2.612V} \right)^{2/3} = 3.1 \text{ K} \end{aligned}$$

The result is greater than experimental value of  $2.17 \text{ K}$ .

11.3

From (7.65), we have

$$D(\varepsilon) = \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \varepsilon^{1/2}$$

Thus, we can get the energy

$$\begin{aligned} U &= \int_0^\infty \varepsilon D(\varepsilon) f(\varepsilon) d\varepsilon = \int_0^\infty \varepsilon \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \varepsilon^{1/2} \frac{1}{e^{(\varepsilon-\mu)/\tau} - 1} d\varepsilon \\ &= \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \int_0^\infty \varepsilon^{3/2} \frac{1}{\lambda^{-1} e^{\varepsilon/\tau} - 1} d\varepsilon \end{aligned}$$

With  $x = \varepsilon / \tau, \lambda = 1$ , we have

$$U = \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \tau^{5/2} \int_0^\infty \frac{x^{3/2}}{e^x - 1} dx$$

For the integral, we have

$$\int_0^\infty \frac{x^{3/2}}{e^x - 1} dx = \Gamma(5/2) \zeta(5/2) \approx \frac{3}{4} \sqrt{\pi} \times 1.341 \approx 1.783$$

Thus

$$U = 1.783 \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \tau^{5/2}$$

For  $d\sigma = \left(\frac{dU}{\tau}\right)_V$ , we can obtain

$$\begin{aligned} \sigma &= \int \frac{dU}{\tau} + \text{constant} \\ &= 1.783 \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \int_0^\tau \frac{5}{2} \tau'^{3/2} \frac{d\tau'}{\tau'} + \text{constant} \\ &= 1.783 \frac{5}{2} \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \int_0^\tau \tau'^{1/2} d\tau' + \text{constant} \\ &= 1.783 \frac{5}{2} \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \frac{2}{3} \tau'^{3/2} \Big|_0^\tau + \text{constant} \\ &= 1.783 \frac{5}{2} \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \frac{2}{3} \tau^{3/2} \Big|_0^\tau + \text{constant} \\ &= 1.783 \frac{5}{3} \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \tau^{3/2} + \text{constant} \end{aligned}$$

For  $\sigma \rightarrow 0$  as  $\tau \rightarrow 0$ , thus

$$\sigma = 1.783 \frac{5}{3} \frac{V}{4\pi^2} \left(\frac{2M}{\hbar^2}\right)^{3/2} \tau^{3/2}$$

#### 11.4

For Bose-Einstein distribution function (7.53), we have

$$f(\varepsilon, \tau) = \frac{1}{\exp[(\varepsilon - \mu) / \tau] - 1}$$

Thus,

$$f(0) = \frac{1}{\exp(-\mu / \tau) - 1} \Rightarrow \exp(-\mu / \tau) = 1 + \frac{1}{f(0)}$$

$$f(\varepsilon) = \frac{1}{\exp[(\varepsilon - \mu) / \tau] - 1} \Rightarrow \exp[(\varepsilon - \mu) / \tau] = 1 + \frac{1}{f(\varepsilon)}$$

Also, we have

$$\frac{f(0)}{f(\varepsilon)} = 2 \Rightarrow f(0) = \frac{2}{3}N, f(\varepsilon) = \frac{1}{3}N$$

For  $N \gg 1$ , we can make some approximation

$$\exp(-\mu / \tau) = 1 + 3 / 2N, \exp[(\varepsilon - \mu) / \tau] = 1 + 3 / N$$

$$\Rightarrow \exp(\varepsilon / \tau) = \frac{\exp[(\varepsilon - \mu) / \tau]}{\exp(-\mu / \tau)} = \frac{1 + 3 / N}{1 + 3 / 2N} = 1 + \frac{3 / 2N}{1 + 3 / 2N} \approx 1 + 3 / 2N$$

$$\Rightarrow \varepsilon / \tau = \ln(1 + 3 / 2N) \approx 3 / 2N$$

$$\Rightarrow \tau = 2N\varepsilon / 3$$

#### 11.5

(a)

For reversible heat pump, from (8.4), we have

$$\sigma_l = \sigma_h \Rightarrow Q_l / \tau_l = Q_h / \tau_h$$

$$\Rightarrow Q_l = \frac{\tau_l}{\tau_h} Q_h$$

Thus

$$W = Q_h - Q_l = Q_h - \frac{\tau_l}{\tau_h} Q_h = \frac{\tau_h - \tau_l}{\tau_h} Q_h = \eta_c Q_h$$

$$\Rightarrow \frac{W}{Q_h} = \eta_c = \frac{\tau_h - \tau_l}{\tau_h}$$

For irreversible heat pump, we have

$$\sigma_l < \sigma_h \Rightarrow Q_l / \tau_l < Q_h / \tau_h$$

$$\Rightarrow Q_l < \frac{\tau_l}{\tau_h} Q_h$$

Thus

$$W = Q_h - Q_l > Q_h - \frac{\tau_l}{\tau_h} Q_h = \frac{\tau_h - \tau_l}{\tau_h} Q_h = \eta_c Q_h$$

$$\Rightarrow \frac{W}{Q_h} > \eta_c = \frac{\tau_h - \tau_l}{\tau_h}$$

More work is required for irreversible heat pump.

(b)

We have

$$\frac{W}{Q_h} = \frac{\tau_h - \tau_l}{\tau_h}, \frac{W}{Q_{hh}} = \frac{\tau_{hh} - \tau_l}{\tau_{hh}}$$

Thus

$$\frac{Q_{hh}}{Q_h} = \frac{W / Q_h}{W / Q_{hh}} = \frac{\tau_h - \tau_l}{\tau_h} \times \frac{\tau_{hh}}{\tau_{hh} - \tau_l} = \frac{300 - 270}{300} \times \frac{600}{600 - 270} = \frac{2}{11}$$

(c)

