

HW#6

6.1

We have

$$N = \sum_n \langle s_n \rangle = \sum_n \frac{1}{\exp(\hbar\omega_n / \tau) - 1} = \sum_n \frac{1}{\exp(x) - 1}$$

$$\omega_n = n\pi c / L, x = \hbar\omega_n / \tau, n = \frac{\tau L}{\hbar\pi c} x$$

Similar to (4.17), a factor 2 will be applied for 2 independent polarizations.

$$N = \sum_n \frac{1}{\exp(x) - 1} \approx \pi \int_0^\infty \frac{1}{\exp(x) - 1} n^2 dn = \pi \left(\frac{\tau L}{\hbar\pi c}\right)^3 \int_0^\infty \frac{x^2}{\exp(x) - 1} dx = \pi^{-2} V \left(\frac{\tau}{\hbar c}\right)^3 \int_0^\infty \frac{x^2}{\exp(x) - 1} dx$$

We can get the integral value

$$\begin{aligned} \int_0^\infty \frac{x^2}{\exp(x) - 1} dx &= \int_0^\infty \lim_{n \rightarrow \infty} \left[\frac{1 - \exp(-nx)}{1 - \exp(-x)} - 1 \right] x^2 dx = \sum_{n=1}^\infty \int_0^\infty \exp(-nx) x^2 dx \\ &= \sum_{n=1}^\infty \left. -\frac{\exp(-nx)(2 + 2nx + n^2 x^2)}{n^3} \right|_0^\infty = \sum_{n=1}^\infty \frac{2}{n^3} = 2\zeta(3) \approx 2.40411 \end{aligned}$$

Thus,

$$N \approx 2.404\pi^{-2} V \left(\frac{\tau}{\hbar c}\right)^3$$

6.2

We have

$$C_L = \left(\frac{\partial U}{\partial \tau} \right)_L$$

$$U = \sum_n \langle s_n \rangle \hbar\omega_n = \sum_n \frac{\hbar\omega_n}{\exp(\hbar\omega_n / \tau) - 1} = \tau \sum_n \frac{x}{\exp(x) - 1}$$

$$\omega_n = n\pi v / L, x = \hbar\omega_n / \tau, n = \frac{\tau L}{\hbar\pi v} x$$

Thus,

$$U = \tau \sum_n \frac{x}{\exp(x) - 1} \approx \tau \int_0^\infty \frac{x}{\exp(x) - 1} dn = \frac{\tau^2 L}{\hbar\pi v} \int_0^\infty \frac{x}{\exp(x) - 1} dx$$

Similar to problem 1, we can get

$$\begin{aligned} \int_0^\infty \frac{x}{\exp(x) - 1} dx &= \int_0^\infty \lim_{n \rightarrow \infty} \left[\frac{1 - \exp(-nx)}{1 - \exp(-x)} - 1 \right] x dx = \sum_{n=1}^\infty \int_0^\infty \exp(-nx) x dx \\ &= \sum_{n=1}^\infty \left. -\frac{\exp(-nx)(1 + nx)}{n^2} \right|_0^\infty = \sum_{n=1}^\infty \frac{1}{n^2} = \zeta(2) = \frac{\pi^2}{6} \end{aligned}$$

Thus,

$$U = \frac{\pi\tau^2 L}{6\hbar\nu}$$

$$\Rightarrow C_L = \left(\frac{\partial U}{\partial \tau} \right)_L = \frac{\pi\tau L}{3\hbar\nu}$$

6.3

For capacity of matter, we have

$$C_v = \frac{3}{2} Nk_B$$

For radiation, from (4.20)

$$\frac{U}{V} = \frac{\pi^2}{15\hbar^3 c^3} \tau^4$$

$$\Rightarrow U = \frac{\pi^2 V}{15\hbar^3 c^3} k_B^4 T^4$$

$$\Rightarrow C_{rad} = \left(\frac{\partial U}{\partial T} \right)_V = \frac{4\pi^2 V}{15\hbar^3 c^3} k_B^4 T^3$$

Thus, the ratio of heat capacity of matter to that of radiation is

$$\frac{C_v}{C_{rad}} = \frac{3}{2} Nk_B \times \frac{15\hbar^3 c^3}{4\pi^2 V k_B^4 T^3} = \frac{45n\hbar^3 c^3}{8\pi^2 k_B^3 T^3} \approx 2.8 \times 10^{-10}$$

This ratio is $\sim 10^{-9}$.

6.4

The partition function is

$$Z = \sum_n s_n \exp(-\hbar\omega_n / \tau)$$

s_n is the number of photons in the mode n.

For

$$F = U - \tau\sigma, F = -\tau \ln Z$$

We can get

$$p = \tau \left(\frac{\partial \sigma}{\partial V} \right)_U = - \left(\frac{\partial (U - \tau\sigma)}{\partial V} \right)_{U, \tau} = - \left(\frac{\partial F}{\partial V} \right)_\tau = \tau \left(\frac{\partial \ln Z}{\partial V} \right)_\tau = \frac{\tau}{Z} \left(\frac{\partial Z}{\partial V} \right)_\tau$$

$$= \frac{\tau}{Z} \sum_n s_n \exp(-\hbar\omega_n / \tau) \times -\hbar / \tau \times \frac{d\omega_n}{dV}$$

$$= -\frac{1}{Z} \sum_n s_n \exp(-\hbar\omega_n / \tau) \times \hbar \times \frac{d\omega_n}{dV}$$

From (3.15)

$$\omega_n = n\pi c / L$$

$$\Rightarrow \ln \omega_n = -\ln L + \text{constant} = -\frac{1}{3} \ln V + \text{constant}$$

$$\Rightarrow \frac{d \ln \omega_n}{dV} = \frac{1}{\omega_n} \frac{d\omega_n}{dV} = -\frac{1}{3V}$$

$$\Rightarrow \frac{d\omega_n}{dV} = -\frac{\omega_n}{3V}$$

Thus,

$$p = -\frac{1}{Z} \sum_n s_n \exp(-\hbar\omega_n / \tau) \times \hbar \times \frac{d\omega_n}{dV}$$

$$= -\frac{1}{Z} \sum_n s_n \exp(-\hbar\omega_n / \tau) \times \hbar \times -\frac{\omega_n}{3V}$$

$$= \frac{1}{3V} \frac{1}{Z} \sum_n s_n \exp(-\hbar\omega_n / \tau) \times \hbar\omega_n$$

$$= \frac{U}{3V}$$