

HW#8

8.1

(a)

The system has two energy states $(0, \varepsilon)$, the accessible states are below.

| State number | Description | N | Energy |
|--------------|-------------|---|---------------|
| 1 | $(0, 0)$ | 0 | 0 |
| 2 | $(1, 0)$ | 1 | 0 |
| 3 | $(0, 1)$ | 1 | ε |

Thus, the Gibbs sum is

$$Z = \exp(0/\tau) + \exp(\mu/\tau) + \exp[(\mu - \varepsilon)/\tau] \\ = 1 + \lambda + \lambda \exp(-\varepsilon/\tau)$$

(b)

$$\langle N \rangle = \frac{0 \times 1 + 1 \times \lambda + 1 \times \lambda \exp(-\varepsilon/\tau)}{Z} = \frac{\lambda + \lambda \exp(-\varepsilon/\tau)}{Z}$$

(c)

At energy ε , we just need the last term in the expression for $\langle N \rangle$

$$\langle N(\varepsilon) \rangle = \frac{1 \times \lambda \exp(-\varepsilon/\tau)}{Z} = \frac{\lambda \exp(-\varepsilon/\tau)}{Z}$$

(d)

The thermal average energy is

$$\langle U \rangle = \frac{0 \times 1 + 0 \times \lambda + \varepsilon \times \lambda \exp(-\varepsilon/\tau)}{Z} = \frac{\varepsilon \lambda \exp(-\varepsilon/\tau)}{Z}$$

(e)

In this case, the accessible states are below.

| State number | Description | N | Energy |
|--------------|-------------|---|---------------|
| 1 | $(0, 0)$ | 0 | 0 |
| 2 | $(1, 0)$ | 1 | 0 |
| 3 | $(0, 1)$ | 1 | ε |
| 4 | $(1, 1)$ | 2 | ε |

Thus, the Gibbs sum is

$$Z = \exp(0/\tau) + \exp(\mu/\tau) + \exp[(\mu - \varepsilon)/\tau] + \exp[(2\mu - \varepsilon)/\tau] \\ = 1 + \lambda + \lambda \exp(-\varepsilon/\tau) + \lambda^2 \exp(-\varepsilon/\tau) \\ = (1 + \lambda)[1 + \lambda \exp(-\varepsilon/\tau)]$$

8.2

(a)

We have the Gibbs sum is

$$Z = 1 + \lambda(O_2) \exp(-\varepsilon_A/\tau) + \lambda(CO) \exp(-\varepsilon_B/\tau)$$

When consider the system in the absence of CO, we have

$$Z = 1 + \lambda(O_2) \exp(-\varepsilon_A/\tau)$$

$$\begin{aligned}
& \frac{\lambda(O_2)\exp(-\varepsilon_A/\tau)}{1+\lambda(O_2)\exp(-\varepsilon_A/\tau)} = 0.9 \\
& \Rightarrow \lambda(O_2)\exp(-\varepsilon_A/\tau) = 9 \\
& \Rightarrow \exp(-\varepsilon_A/\tau) = 9 \times 10^5 \\
& \Rightarrow -\varepsilon_A/\tau = 13.71 \\
& \Rightarrow \varepsilon_A = -13.71 \times k_B T = -13.71 \times 8.617 \times 10^{-5} eV / K \times 310.15 K = -0.366 eV
\end{aligned}$$

(b)

In this case, we have

$$\begin{aligned}
& \frac{\lambda(O_2)\exp(-\varepsilon_A/\tau)}{1+\lambda(O_2)\exp(-\varepsilon_A/\tau)+\lambda(CO)\exp(-\varepsilon_B/\tau)} = 0.1 \\
& \Rightarrow \frac{9}{10+\lambda(CO)\exp(-\varepsilon_B/\tau)} = 0.1 \\
& \Rightarrow \lambda(CO)\exp(-\varepsilon_B/\tau) = 80 \\
& \Rightarrow \exp(-\varepsilon_B/\tau) = 8 \times 10^8 \\
& \Rightarrow -\varepsilon_B/\tau = 20.5 \\
& \Rightarrow \varepsilon_B = -20.5 \times k_B T = -20.5 \times 8.617 \times 10^{-5} eV / K \times 310.15 K = -0.548 eV
\end{aligned}$$

8.3

(a)

From (6.4), we have

$$\begin{aligned}
f(\varepsilon) &= \frac{1}{\exp[(\varepsilon-\mu)/\tau]+1} \\
\Rightarrow \frac{\partial f}{\partial \varepsilon} &= -\frac{1}{\{\exp[(\varepsilon-\mu)/\tau]+1\}^2} \times \exp[(\varepsilon-\mu)/\tau] \times 1/\tau \\
&= -\frac{\exp[(\varepsilon-\mu)/\tau]}{\tau \{\exp[(\varepsilon-\mu)/\tau]+1\}^2}
\end{aligned}$$

When $\varepsilon = \mu$, we have

$$\begin{aligned}
& \exp[(\varepsilon-\mu)/\tau] = 1 \\
\Rightarrow -\frac{\partial f}{\partial \varepsilon} &= \frac{1}{4\tau}
\end{aligned}$$

(b)

$$\begin{aligned}
f(\mu + \delta) &= \frac{1}{\exp[(\mu + \delta - \mu)/\tau] + 1} = \frac{1}{\exp(\delta/\tau) + 1} \\
f(\mu - \delta) &= \frac{1}{\exp[(\mu - \delta - \mu)/\tau] + 1} = \frac{1}{\exp(-\delta/\tau) + 1} = \frac{\exp(\delta/\tau)}{\exp(\delta/\tau) + 1} \\
\Rightarrow f(\mu + \delta) + f(\mu - \delta) &= \frac{1}{\exp(\delta/\tau) + 1} + \frac{\exp(\delta/\tau)}{\exp(\delta/\tau) + 1} = 1 \\
\Rightarrow f(\mu + \delta) &= 1 - f(\mu - \delta)
\end{aligned}$$

8.4

(a)

The accessible states are below.

| State number | N | Energy |
|--------------|---|----------------|
| 1 | 0 | 0 |
| 2 | 1 | ε |
| 3 | 2 | 2ε |

Thus, the Gibbs sum is

$$\begin{aligned}
Z &= \exp(0/\tau) + \exp[(\mu - \varepsilon)/\tau] + \exp[(2\mu - 2\varepsilon)/\tau] \\
&= 1 + \lambda \exp(-\varepsilon/\tau) + \lambda^2 \exp(-2\varepsilon/\tau)
\end{aligned}$$

The ensemble average occupancy is

$$\langle N \rangle = \frac{0 \times 1 + 1 \times \lambda \exp(-\varepsilon/\tau) + 2 \times \lambda^2 \exp(-2\varepsilon/\tau)}{Z} = \frac{\lambda \exp(-\varepsilon/\tau) + 2\lambda^2 \exp(-2\varepsilon/\tau)}{Z}$$

(b)

The accessible states are below.

| State number | Description | N | Energy |
|--------------|-------------|---|----------------|
| 1 | (0, 0) | 0 | 0 |
| 2 | (1, 0) | 1 | ε |
| 3 | (0, 1) | 1 | ε |
| 4 | (1, 1) | 2 | 2ε |

Thus, the Gibbs sum is

$$\begin{aligned}
Z &= \exp(0/\tau) + \exp[(\mu - \varepsilon)/\tau] + \exp[(\mu - \varepsilon)/\tau] + \exp[(2\mu - 2\varepsilon)/\tau] \\
&= 1 + 2\lambda \exp(-\varepsilon/\tau) + \lambda^2 \exp(-2\varepsilon/\tau)
\end{aligned}$$

The ensemble average occupancy is

$$\begin{aligned}
\langle N \rangle &= \frac{0 \times 1 + 1 \times \lambda \exp(-\varepsilon/\tau) + 1 \times \lambda \exp(-\varepsilon/\tau) + 2 \times \lambda^2 \exp(-2\varepsilon/\tau)}{Z} \\
&= \frac{2\lambda \exp(-\varepsilon/\tau) + 2\lambda^2 \exp(-2\varepsilon/\tau)}{Z}
\end{aligned}$$