## PHYS852 Quantum Mechanics II, Spring 2010 HOMEWORK ASSIGNMENT 13

Topics covered: Hilbert-space Frame Transformations, Time-Dependent Perturbation Theory

1. The Hamiltonian for a driven two-level system is

$$H = \hbar\omega_0 |2\rangle \langle 2| + \hbar\Omega \cos(\omega t) \left( |1\rangle \langle 2| + |2\rangle \langle 1| \right), \tag{1}$$

where  $\omega_0$  is the separation between the bare levels, and  $\omega$  is the driving frequency.

- (a) Make a frame transformation generated by the operator  $G = \hbar \omega |2\rangle \langle 2|$ , and determine the equation of motion for the state-vector in the new frame, defined by  $|\psi_G(t)\rangle = U_G(t)|\psi(t)\rangle$ .
- (b) Make the rotating wave approximation (RWA) by assuming that  $\omega \approx \omega_0$ , and dropping any terms that oscillate at or near  $2\omega_0$ . Write, in terms of the detuning  $\Delta = \omega_0 - \omega$ , the effective time-independent Hamiltonian,  $H_G$ , that then governs the time evolution of  $|\psi_G(t)\rangle$ .
- (c) Assume that the system begins at time t = 0 in the ground-state of  $H_G$ , and calculate  $|\psi_G(t)\rangle$ . Is this a stationary state in the rotating frame? Now use  $|\psi_S(t)\rangle = U_G^{\dagger}(t)|\psi_G(t)\rangle$  to see what this state looks like in the Scrödinger picture. Is it a stationary state in the Schrödinger picture?
- (d) Assuming the system begins in the ground state of  $H_G$ , use second-order time-dependent perturbation theory to treat the fast-oscillating terms that were discarded in the RWA, and compute the probability to find the system in the excited state of  $H_G$ , at time t > 0.
- (e) Assume that at time t = 0, we have  $\Delta > 0$ ,  $\Omega = 0$ , and  $|\psi_G(0)\rangle = |1\rangle$ . If  $\Omega$  is smoothly increased from zero to  $\Omega_0$  on time-scale  $T >> 1/\Delta$ , what is the state of the system at time t = T?
- 2. Consider a system described by the Hamiltonian:

$$H = -\frac{\hbar\omega}{4} \left( AA + A^{\dagger}A^{\dagger} \right), \qquad (2)$$

where  $[A, A^{\dagger}] = 1$ . Find and solve the Heisenberg equations of motion for  $A_H(t)$  and  $A_H^{\dagger}(t)$ . Use these solutions to compute the expectation values of X and P, as well as the variances  $\Delta X$ , and  $\Delta P$ , as functions of time, for the case where the initial state satisfies  $A|\psi_S(0)\rangle = \alpha|\psi_S(0)\rangle$ , where  $\alpha$  is an arbitrary complex number. For the case  $\alpha = 0$ , show that  $\langle X \rangle_t = \langle P \rangle_t = 0$ , but  $\Delta X$  and  $\Delta P$  grow rapidly in time.

Now re-express the Hamiltonian in terms of X and P. Do your previous answers make sense given this viewpoint?

3. The Hamiltonian for a hydrogen atom is

$$H = \frac{P_r^2}{2\mu} + \frac{L^2}{2\mu R^2} - \frac{e^2}{4\pi\epsilon_0 R}.$$
(3)

First, use the properties

$$\langle \vec{r} | P_r^2 | \psi \rangle = -\hbar^2 \frac{1}{r^2} \frac{d}{dr} r^2 \frac{d}{dr} \langle \vec{r} | \psi \rangle, \tag{4}$$

and

$$\langle \vec{r} | P_r | \psi \rangle = -i\hbar \frac{d}{dr} \langle \vec{r} | \psi \rangle, \tag{5}$$

to compute the commutators  $[P_r^2, R]$  and  $[R^{-s}, P_r]$ , then use these commutators to derive the Heisenberg equations of motion for R and  $P_r$ .

4. Compute the density of states, n(E), for a massive particle in a cubic volume of side length L. Then compute the density of states for a two-dimensional massive particle confined to a square area of side length L, and also for a one-dimensional massive particle in an infinite square-well of width L. Then do the same for a photon, whose energy is related to its wavevector by  $E(\vec{k}) = \hbar c |k|$ .

To determine the density of states at energy E, first determine N(E), which is the number of quantized k-values inside a sphere of radius k(E). Do this by determining the volume in k-space occupied by a single state, and then divide the volume of the energy-sphere by the single-mode volume. Then compute the density of states via  $n(E) = \frac{d}{dE}N(E)$ .

- 5. Estimate the spontaneous photon-emission rate of an excited atom via Fermi's golden rule, use the density of states for a photon in a cube of volume V. To estimate the coupling matrix element, use the dipole energy operator  $V = -\vec{d} \cdot \vec{\mathcal{E}}$ , where  $d = ea_0$  is the atomic dipole moment, and  $\mathcal{E}$  is the electric field of a single photon in a volume V. To get the value of  $\mathcal{E}$ , take the photons energy to be  $\hbar\omega$ , and use the standard energy density of an electromagnetic field  $u = \frac{\epsilon_0}{2} (\mathcal{E}^2 + c^2 \mathcal{B}^2)$ . Relate  $\mathcal{B}$ , the magnetic field of the photon, to its electric field via  $\nabla \times \vec{\mathcal{E}} = -\frac{d}{dt}\vec{\mathcal{B}}$ , which from dimensional analysis gives  $k\mathcal{E} \approx \omega \mathcal{B}$ . For frequencies in the visible spectrum, what is your estimate of  $\Gamma$ ?
- 6. Consider a harmonic oscillator described by  $H^{(0)} = \hbar \Omega(A^{\dagger}A + 1/2)$ . Now consider two possible perturbations  $V_1 = bX^4 \cos(\omega t)$ , and  $V_2 = g\delta(X)e^{-\gamma t}$ . Assume that the system begins at t = 0, in the ground-state of  $H^{(0)}$ ,  $|0\rangle$ , and use time-dependent perturbation theory to compute the probability to second-order, for the system to be found in the the  $n^{th}$  eigenstate of  $H^{(0)}$  at time t. Consider both n = 0 and  $n \neq 0$  cases.