

1. Using the equations of motion for the wavefunction, show that the density and current defined by

$$\rho(\vec{r}, t) = |\psi(\vec{r}, t)|^2,$$

$$\vec{j}(\vec{r}, t) = \frac{-i\hbar}{2m}(\psi^*(\vec{r}, t)\nabla\psi(\vec{r}, t) - (\nabla\psi^*(\vec{r}, t))\psi(\vec{r}, t)) - \frac{e\vec{A}}{mc}|\psi(\vec{r}, t)|^2,$$

satisfies the continuity equation,

$$\partial_t \rho + \nabla \cdot \vec{j} = 0.$$

$$\begin{aligned} \frac{d\rho}{dt} &= \frac{-i}{\hbar}(H\psi^*)\psi + \frac{i}{\hbar}\psi^*H\psi \\ &= \frac{i}{\hbar} \left\{ \frac{\hbar^2}{2m}(\nabla^2\psi^*)\psi - \frac{\hbar^2}{2m}\psi^*(\nabla^2\psi) \right. \\ &\quad \left. - \frac{\hbar e}{m}(\vec{A} \cdot \vec{\nabla}\psi^*)\psi - \frac{\hbar e}{m}\psi^*(\vec{A} \cdot \vec{\nabla}\psi) \right. \\ &\quad \left. - \cancel{\frac{e\vec{A}}{2m}\psi^*\psi} + \cancel{\left(\frac{e\vec{A}}{2m}\psi^*\psi\right)} \right\} \\ &= \vec{\nabla} \cdot \left\{ \frac{(\nabla\psi^*)\psi - \psi^*(\nabla\psi)}{2m} - \frac{e\vec{A}}{m}\psi^*\psi \right\} \\ &= \vec{\nabla} \cdot \vec{j} \end{aligned}$$

2. Consider a particle of charge e traveling in the electromagnetic potentials

$$\mathbf{A}(\mathbf{r}, t) = -\nabla\Lambda(\mathbf{r}, t), \quad \Phi(\mathbf{r}, t) = \frac{1}{c} \frac{\partial\Lambda(\mathbf{r}, t)}{\partial t}$$

where $\Lambda(\mathbf{r}, t)$ is an arbitrary scalar function.

- (a) What are the electromagnetic fields described by these potentials?
 (b) Show that the wave function of the particle is given by

$$\psi(\mathbf{r}, t) = \exp\left[-\frac{ie}{\hbar c}\Lambda(\mathbf{r}, t)\right] \psi^0(\mathbf{r}, t),$$

where ψ^0 solves the Schrödinger equation with $\Lambda = 0$

- (c) Let $V(\mathbf{r}, t) = e\Phi(t)$ be a spatially uniform time varying potential. Show that

$$\psi(\mathbf{r}, t) = \exp\left[-\frac{ie}{\hbar} \int_{-\infty}^t \Phi(t') dt'\right] \psi_0(\mathbf{r}, t)$$

is a solution if ψ_0 is a solution with $\Phi = 0$.

(a) $\vec{E} = \frac{1}{c} (\nabla \partial_t \Lambda - \nabla \partial_t \Lambda) = \underline{0}$
 $\vec{B} = \nabla \times (\nabla \Lambda) = \underline{0}$, $B_i = \epsilon_{ijk} \partial_j \partial_k \Lambda = 0$

(b) $i\hbar \partial_t \psi = \frac{e}{c} (\partial_t \Lambda) \psi + e^{-\frac{ie}{\hbar c} \Lambda} H_0 \psi_0$
 $H \psi = \frac{(-i\hbar \nabla - e\vec{A}/c)^2}{2m} e^{-\frac{ie}{\hbar c} \Lambda} \psi_0 + \frac{e}{c} \partial_t \Lambda \psi$
 $\vec{A} = -\nabla \Lambda$
 $(-i\hbar \nabla - e\vec{A}/c) e^{-\frac{ie}{\hbar c} \Lambda} \psi_0 = e^{-\frac{ie}{\hbar c} \Lambda} (-i\hbar \nabla \psi_0 - \cancel{\frac{e\hbar}{c}} - \cancel{\frac{\nabla \Lambda}{c}}) \psi_0$
 $H \psi = e^{-\frac{ie}{\hbar c} \Lambda} \left(-\frac{\hbar^2}{2m} \nabla^2 \psi_0 + \frac{e}{c} \partial_t \Lambda \psi_0 \right)$

Comparing

$$i\hbar \partial_t \psi = H \psi \quad \checkmark$$

(c) Let $\Lambda(t) = \int_{-\infty}^t dt' \frac{\Phi(t')}{c}$, from above,
 $\psi = e^{-\frac{ie}{\hbar c} \Lambda} \psi_0 = e^{-\frac{ie}{\hbar} \int_{-\infty}^t dt' \Phi(t')} \psi_0 \quad \checkmark$

3. For a gauge transformation, described in Eq. (3.7), including the associated phase change to the wave function ψ , described in Eq. (3.8),

(a) Show that the charge density $e\psi^*\psi$ and the current is unchanged by the gauge transformation

(b) Show that the current

$$\vec{j} = \frac{e}{2m} \left[\psi^* (-i\hbar \nabla - \frac{e}{c} \vec{A}) \psi + (i\hbar \nabla \psi^*) \psi \right]$$

is unchanged.

(c) Show that $\langle \chi | \mathbf{H} | \psi \rangle$ is unchanged in a gauge transformation where Λ is independent of time.

a) $\psi = e^{-\frac{ie}{\hbar c} \Lambda} \psi_0$, $\psi^* \psi = \psi_0^* \psi$
↑
pure phase

b) $\frac{e}{2m} e^{-\frac{ie}{\hbar c} \Lambda} \psi_0^* (-i\hbar \nabla - \frac{e}{c} \vec{A}) e^{-\frac{ie}{\hbar c} \Lambda} \psi_0$
 $+ \frac{e}{2m} (i\hbar \nabla (e^{-\frac{ie}{\hbar c} \Lambda} \psi_0^*)) \psi_0 e^{i\Lambda}$
 $= \frac{e}{2m} \psi_0^* (-i\hbar \nabla - \frac{e}{c} \vec{A}) \psi_0 + \frac{e}{2m} (i\hbar \nabla \psi_0^*) \psi_0$
 $+ \frac{e}{2m} e^{-\frac{ie}{\hbar c} \Lambda} \psi_0^* (i\hbar) \left(\frac{ie}{\hbar c} \nabla \Lambda \right) \psi_0 e^{-\frac{ie}{\hbar c} \Lambda}$
 $+ \frac{e}{2m} (i\hbar) \left(\frac{ie}{\hbar c} \right) (\nabla \Lambda) \psi_0^* \psi_0$

Let $\vec{A} = \vec{A}_0 - \nabla \Lambda$

→ $= \frac{e}{2m} (\psi_0^* (-i\hbar \nabla - \frac{e}{c} \vec{A}_0) \psi_0 + (i\hbar \nabla \psi_0^*) \psi_0)$
 $+ 2 \frac{e^2}{2mc} (\nabla \Lambda) \psi_0^* \psi_0 - \frac{e^2}{mc} (\nabla \Lambda) \psi_0^* \psi_0$
 $= \frac{e}{2m} (\psi_0^* (-i\hbar \nabla - \frac{e}{c} \vec{A}_0) \psi_0 + (i\hbar \nabla \psi_0^*) \psi_0)$

ⓐ $(-i\hbar \nabla - \frac{e}{c} \vec{A} + \frac{e}{c} \nabla \Lambda) e^{-\frac{ie}{\hbar c} \Lambda} \psi_0$
 $= e^{-\frac{ie}{\hbar c} \Lambda} (-i\hbar \nabla - \frac{e}{c} \vec{A} + \frac{e}{c} \nabla \Lambda - i\hbar \frac{(-ie}{\hbar c} \nabla \Lambda)) \psi_0 e^{-\frac{ie}{\hbar c} \Lambda}$
 $= (-i\hbar \nabla - \frac{e}{c} \vec{A}) \psi_0 e^{-\frac{ie}{\hbar c} \Lambda}$

4. Find the function $\Lambda(\vec{r}, t)$ that corresponds to the gauge transformation in Eq. (3.7) responsible for re-expressing the vector potential in Eq. (3.9) to the form of Eq. (3.10), and show that both forms give the same magnetic field.

$$\text{II. } A_y' = Bx, \quad A_z' = A_x' = 0, \quad \vec{\nabla} \times \vec{A}' = B \hat{z}$$

$$\text{I. } A_\varphi = \frac{1}{2} B \rho$$

$$A_x = -A_\varphi \sin \varphi, \quad A_y = A_\varphi \cos \varphi, \quad A_z = 0$$

$$A_x = -\frac{1}{2} B \rho \frac{y}{\rho} = -\frac{1}{2} B y$$

$$A_y = \frac{1}{2} B \rho \frac{x}{\rho} = \frac{1}{2} B x$$

$$\vec{A}' = \vec{A} + \frac{1}{2} B \times y, \quad \Lambda = -\frac{1}{2} B \times y$$

$$(\vec{\nabla} \times \vec{A}')_x = \frac{\partial A_y}{\partial z} - \frac{\partial A_z}{\partial y} = 0$$

$$(\vec{\nabla} \times \vec{A}')_y = \frac{\partial A_z}{\partial x} - \frac{\partial A_x}{\partial z} = 0$$

$$(\vec{\nabla} \times \vec{A}')_z = \frac{\partial A_x}{\partial y} - \frac{\partial A_y}{\partial x} = \frac{1}{2} B + \frac{1}{2} B = B \quad \checkmark$$

5. The expression for the \bar{v}_y in Eq. (3.20) is only valid for non-relativistic velocities, where $|E| \ll |B|$. For a uniform magnetic field $B\hat{z}$, with no electric field, consider the form for the vector potential in Eq. (3.10). Performing a relativistic boost (Lorentz transformation), but for non-relativistic velocities, in the y direction by a velocity v_y , what is the resulting zeroth component of the vector potential A_0 ? Equating this with the electric scalar potential, express the strength of the resulting electric field in terms of v_y and B .

In lab frame $A^{(0)} = \Phi = -Ex, A^{(y)} = Bx$

Boosted,

$$A^{(0)'} = \gamma \Phi + \gamma v A^{(y)}$$

$$= \gamma x (-E + v B)$$

Choose $v_y = E/B$ to make electric field disappear.